

Symmetric Tensor Networks and Topological Phases

Congratulations to KITS!

Two types of symmetric topological phases

SPT --- symmetry protected topological phases: (Pollmann, Berg, Turner, Oshikawa, Chen, Liu, Gu, Wen...)

Example: Haldane spin chain, integer quantum Hall states, topological insulators...

Features:

- no topological order
- anomalous edge states protected by symmetry

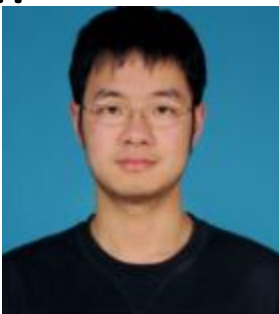
SET --- symmetry enriched topological phases: (Wen, Essin, Hermele, Mesaros, YR, Barkeshli....)

Example: toric code, gapped quantum spin liquids, fractional quantum Hall states...

Features:

- topological order (anyon excitations in 2d)
- symmetry can be fractionalized (e.g. $e/3$ quasiparticle in Laughlin's state).

Acknowledgement:



- **Collaborators:**

- (Boston College) Shenghan Jiang Xu Yang
- Panjin Kim, Hyungyong Lee, Jung Hoon Han (Sungkyunkwan University)
- Brayden Ware, Chao-Ming Jian, Michael Zaletel (StationQ)

- **References:**

- arXiv: 1505.03171, S. Jiang, Y. Ran
- arXiv: 1509.04358, P. Kim, H. Lee, S. Jiang, B. Ware, C. Jian, M. Zaletel, J. Han, Y. Ran
- arXiv: 1610.02024, S. Jiang, P. Kim, J. Han, Y. Ran
- arXiv: 1611.07652, S. Jiang, Y. Ran

Motivations

- We focus on **bosonic** topological (SET or SPT) phases, which require strong interactions to realize.

(1) Conceptual issues:

-- Classification problems

(SPT phases with spatial symmetries.)

(2) “Practical” issues: How to realize them?

-- Physical intuitions/guiding principles?

(Are there criteria like the band-inversion picture in topological insulators?)

-- Numerical methods suitable for searching for these topological phases in models?

(How to write down generic variational wavefunctions?)

Main result

Based on tensor-network formulation, we develop **a machinery** to:

- (1) systematically (but partially) classify topological phases
- (2) construct generic variational wavefunctions for these phases

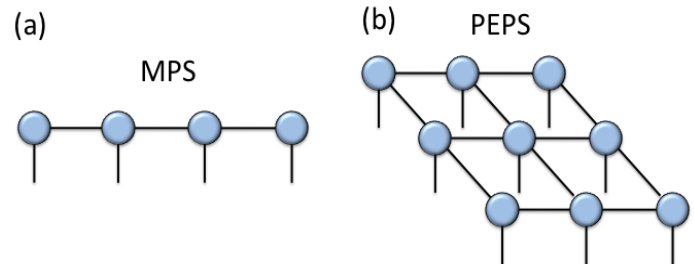
(onsite and spatial)
symmetries of the system



(Partially) classification of topological
phases and construction of generic
wavefunctions for each class

- This machinery answers:
How many classes of symmetric tensor-network wavefunctions that cannot be smoothly deformed into each other under certain assumptions?

1D-MPS, 2D-PEPS, and 3D generalizations

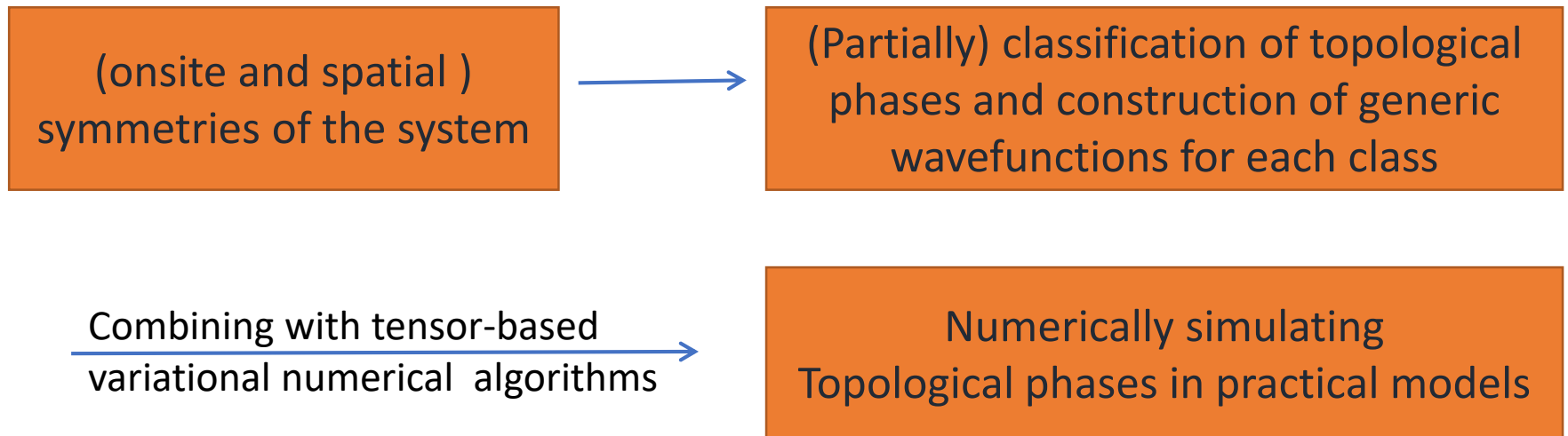


figures from R. Orus,
Annals Phys. (2014)

Main result

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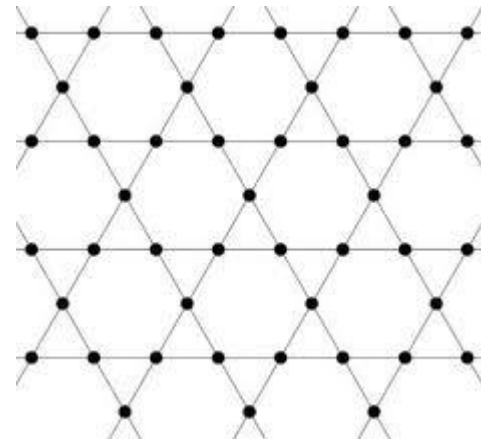
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Some applications of this machinery

(1) Classification and simulation of competing spin liquids in the spin-1/2 Heisenberg model on the kagome lattice.

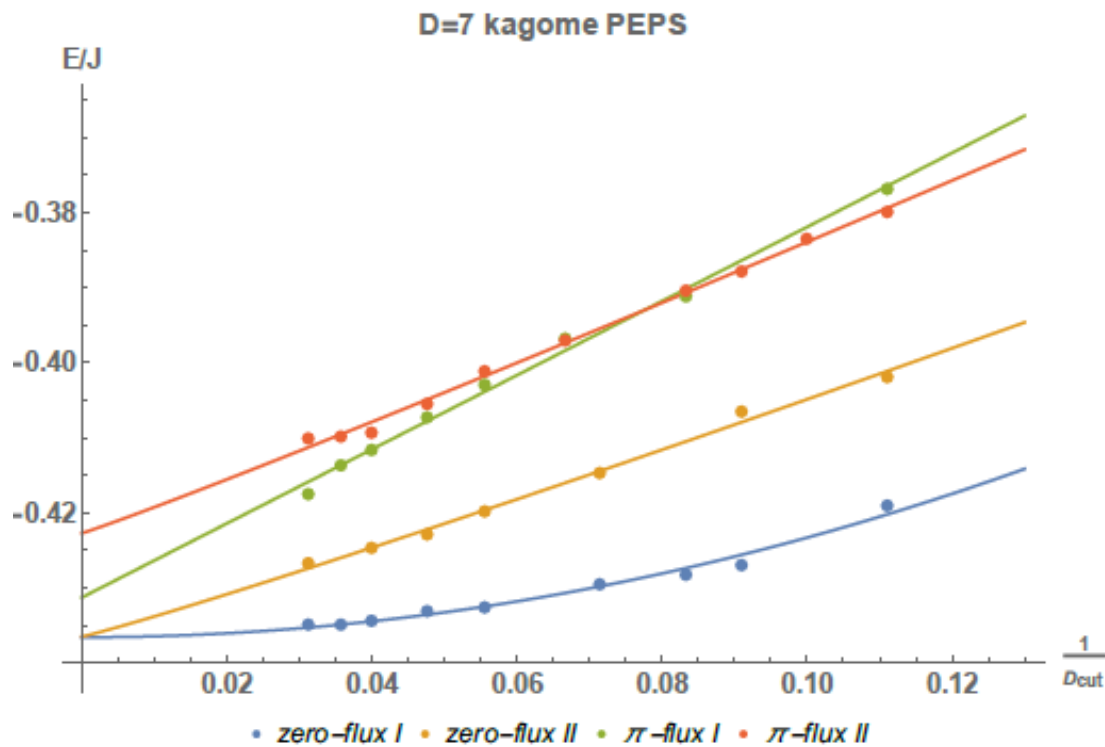
$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j$$



Sachdev, Marston, Senthil, Singh, Evenbly, Vidal, Ran, Hermele, Wen, Lee, Wang, Vishwanath, Weng, Xiang, Sheng, Iqbal, Becca, Sorella, Poilblanc, White, Huse, Depenbrock, McCulloch, Schollwöck, Jiang, Balents, Mei, He, Zaletel, Oshikawa, Pollmann... and many more

Some applications of this machinery

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$D_{cut} \sim$ virtual states kept
when performing tensor
contraction

Some applications of this machinery

(1) Classification and simulation of competing spin liquids in the spin-1/2 Heisenberg model on the kagome lattice.

(2) Classification of bosonic cohomological SPT: $H^{d+1}(SG, U(1))$

- SG : on-site and lattice symmetries (onsite (Chen, Liu, Gu, Wen...), lattice (Chen, Hermele, Fu, Qi, Furusaki, Cheng...))
- T and P (mirror) should be treated as “anti-unitary”
- Generic tensor wavefunctions for every class (if SG is discrete)

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(3) A by-product: a general connection between “conventional” fractionalized phases and SPT phases in 2D via anyon condensation.

Plan

- Anyon condensation mechanism:
“conventional” fractionalized phases \rightarrow SPT phases

A by-product: anyon condensation

An example:

Z_2 gauge theory $\xrightarrow{\text{Condense gauge fluxes}}$ Trivial confined phase

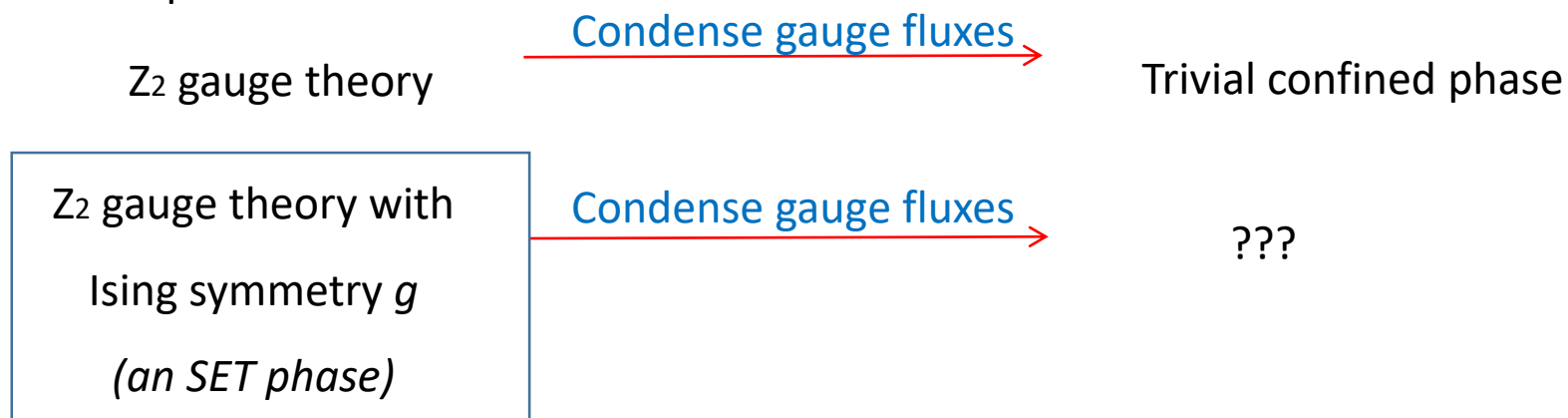
This is the well-studied deconfinement-confinement phase transition.

I will show:

implementing symmetry into such a transition can lead to SPT phases

A by-product: anyon condensation

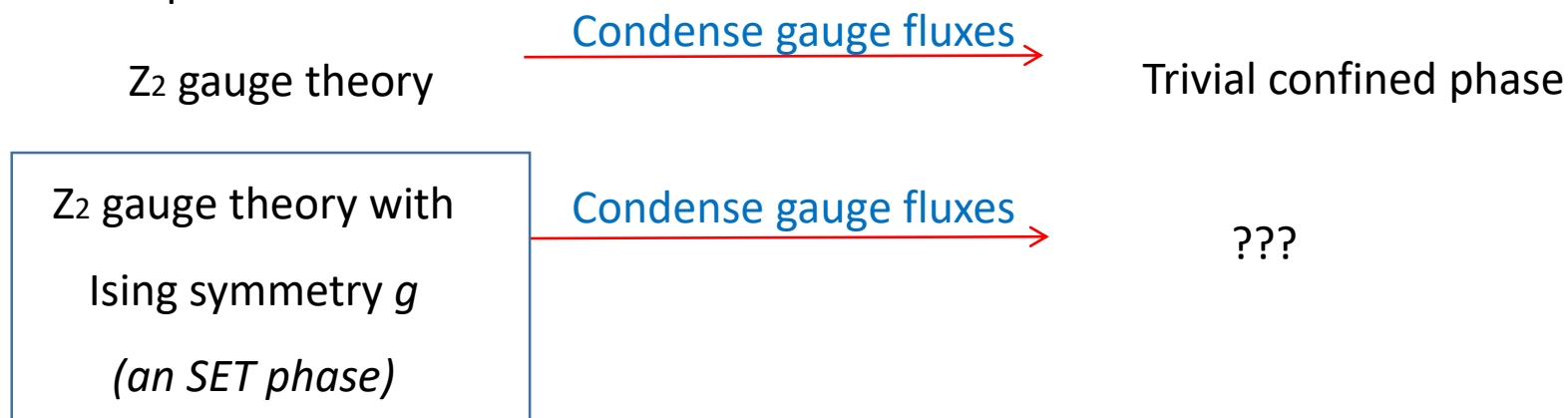
An example:



SET: symmetry-enriched topological phases

A by-product: anyon condensation

An example:



- Consider a particular SET phase:

(e: gauge charge, m: gauge flux)

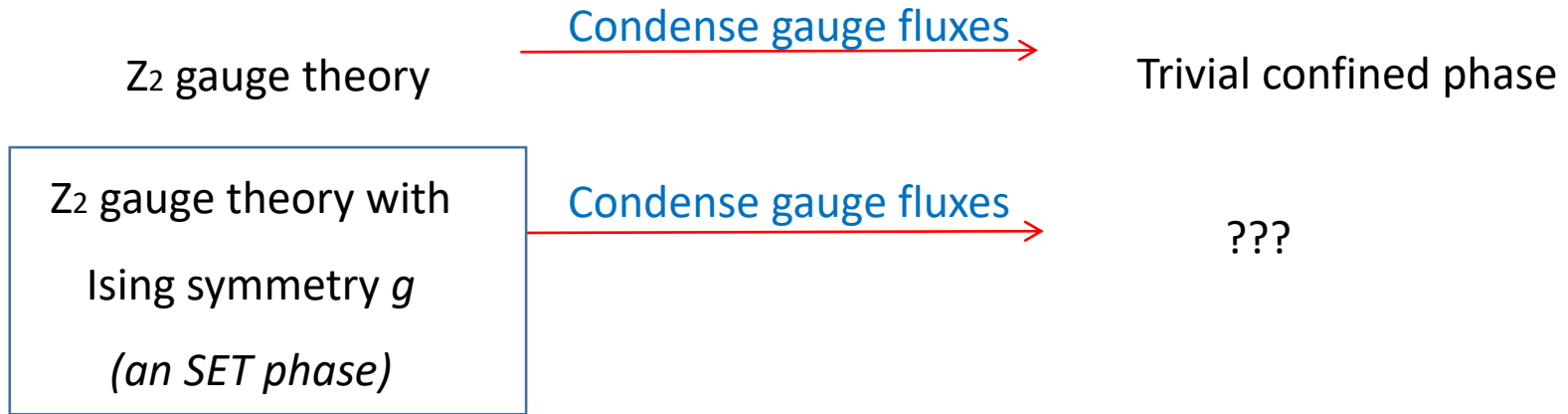
$$[g(e)]^2 = -1, [g(m)]^2 = 1$$

Namely: Ising symmetry is fractionalized on the e-particle

This is a rather conventional fractionalized phase without gapless edge states

A by-product: anyon condensation

An example:



- Consider a particular SET phase:

(e : gauge charge, m : gauge flux)

$$[g(e)]^2 = -1, [g(m)]^2 = 1$$

- Condense m with $g(m) = 1 \rightarrow$ trivial Ising paramagnet
- Condense m with $g(m) = -1 \rightarrow$ nontrivial Ising SPT



Schematic phase diagram

General criteria for anyon condensation

an SET phase $\xrightarrow{\text{Condense gauge fluxes}}$ an SPT phase

- Gauge group: $Z_{N_1} \times Z_{N_2} \times \dots$ & symmetry group: SG
- e-particles feature nontrivial symmetry fractionalization
- m-particles have trivial fractionalization, but can carry usual quantum numbers

$$\Omega_{g_1} \cdot \Omega_{g_2} = \lambda(\mathbf{g}_1, \mathbf{g}_2) \cdot \Omega_{g_1 g_2} \quad \Omega_g \sim \text{symmetry defect}, \quad \lambda \sim \text{certain } m \text{ particle}$$

- Condensing m-particles without breaking symmetry, which requires:
 1. Condensed m 's carry 1D symmetry irrep: $\chi_m(\mathbf{g})$
 2. $\chi_m(\mathbf{g}) \cdot \chi_{m'}(\mathbf{g}) = \chi_{mm'}(\mathbf{g})$
- After condensing those m 's, we get an SPT phase

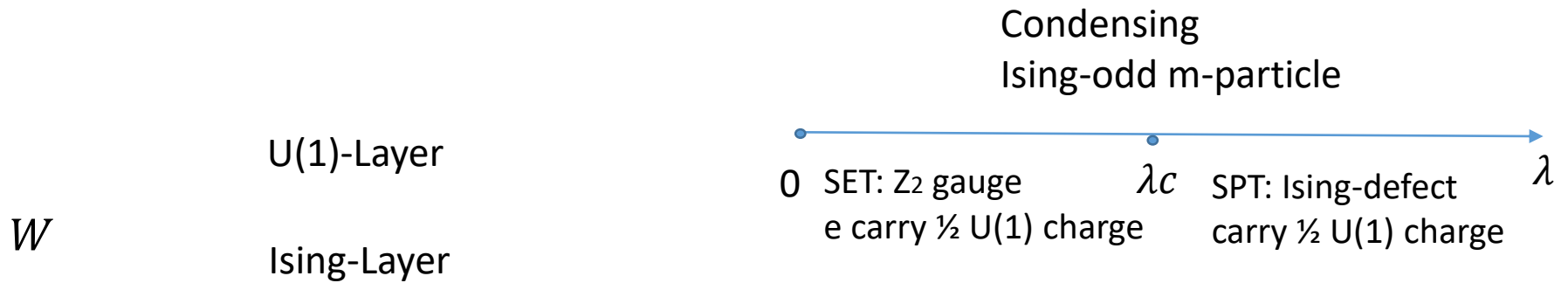
$$\omega(\mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3) \equiv \chi_{\lambda(\mathbf{g}_2, \mathbf{g}_3)}(\mathbf{g}_1), \quad [\omega] \in H^3(SG, U(1))$$

A somewhat simple model realizing SPT

- Following the anyon-condensation mechanism, we can design somewhat simple models realizing bosonic SPT phases. (need 3-spin interactions)
- The model looks like this:

Global symmetry: $U(1) \times \text{Ising}$

$$H = H_{U(1)} + H_{\text{Ising}} + \lambda \cdot W$$



$$H_{\text{Ising}} = h \cdot \sum \sigma^x$$

A somewhat simple model realizing SPT

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$$H = H_{U(1)} + H_{\text{Ising}} + \lambda \cdot W$$

U(1)-layer: Half-filled hard-core bosons on the **kagome lattice**

$$H_{U(1)} = -t \sum b_i^\dagger b_j + V_1 \sum n_i n_j + V_2 \sum n_i n_j + V_3 \sum n_i n_j$$

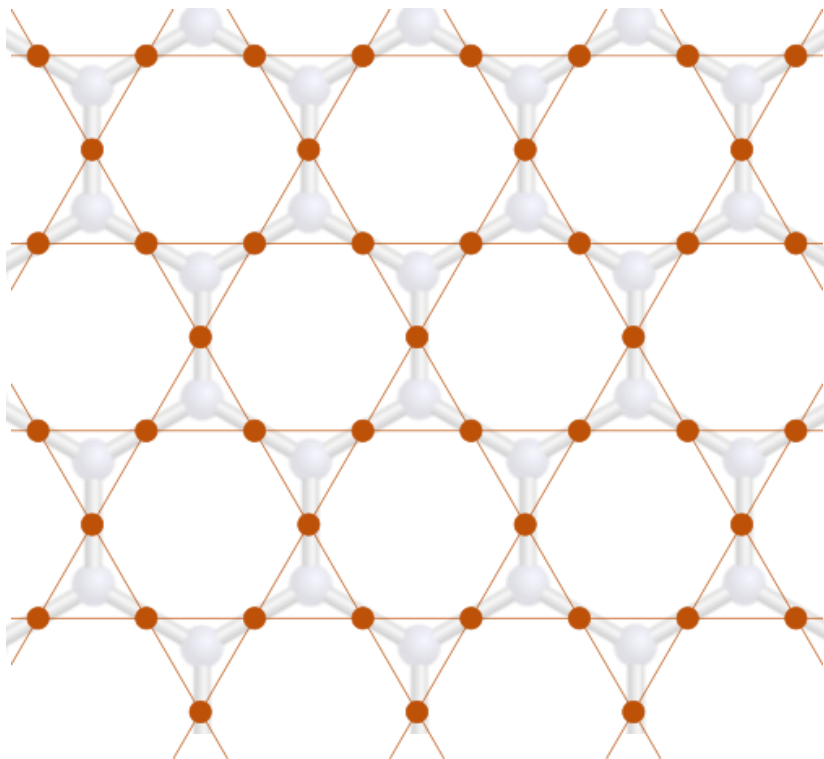
$$t \ll V_1 = V_2 = V_3 = V$$

In this regime, $H_{U(1)}$ is in a deconfined \mathbb{Z}_2 spin liquid phase:
e-particle carries $\frac{1}{2}$ U(1)-charge.
(Balents, Fisher, Girvin 2001)

0 **SET: \mathbb{Z}_2 gauge**
e carry $\frac{1}{2}$ U(1) charge

λc

λ



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$U(1)$ -layer: Half-filled hard-core bosons on the kagome lattice

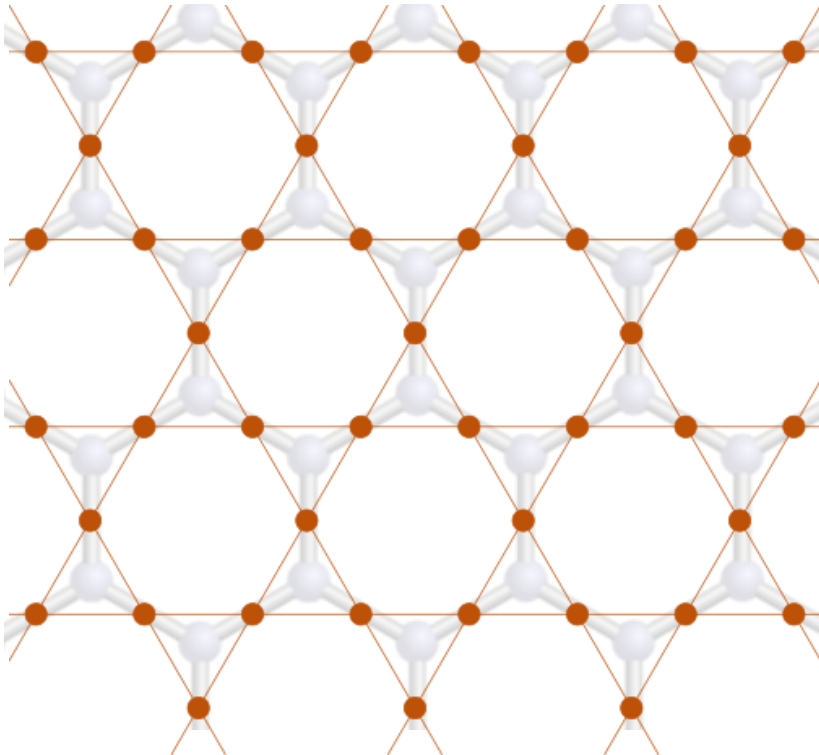
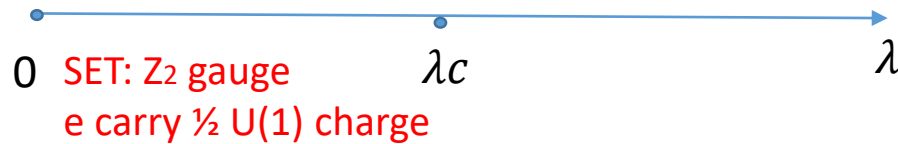
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Ising-layer: transverse field Ising spins on the **honeycomb lattice**

$$H_{\text{Ising}} = h \cdot \sum \sigma^x$$

$$h \ll V$$



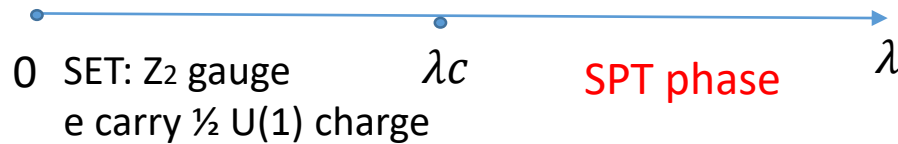
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Condensing
Ising-odd m-particle

$U(1)$ -layer: Half-filled hard-core bosons
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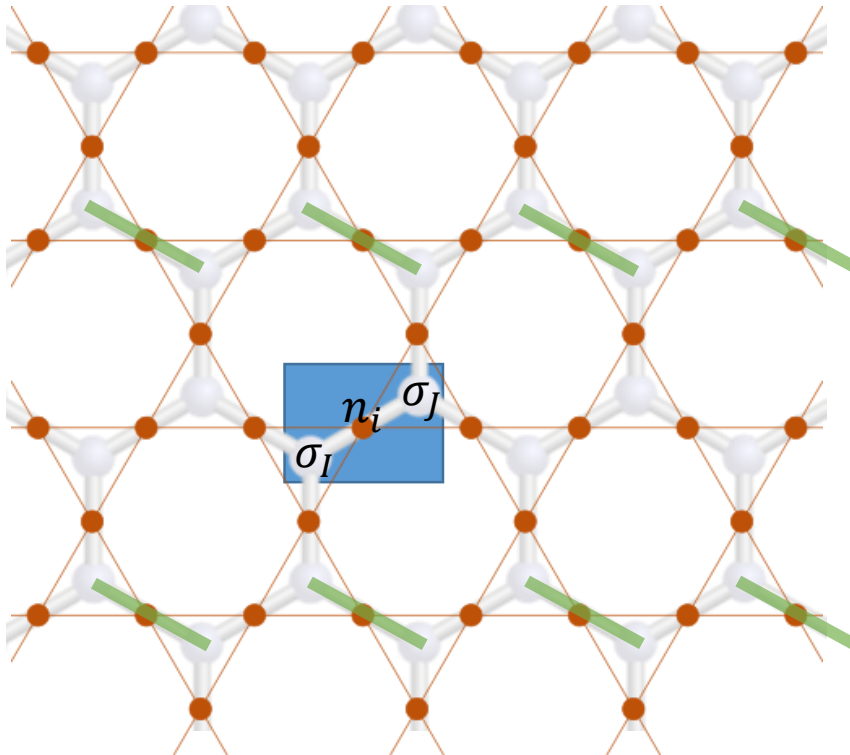
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$$h \ll V$$

$\lambda \cdot W$: 3-spin interaction coupling two layers

$$\lambda \cdot W = \lambda \cdot \sum (n_i - 1/2) \cdot (s_{IJ} \sigma_I^z \sigma_J^z)$$

$s_{IJ} = -1$ on green bonds, $s_{IJ} = +1$ otherwise



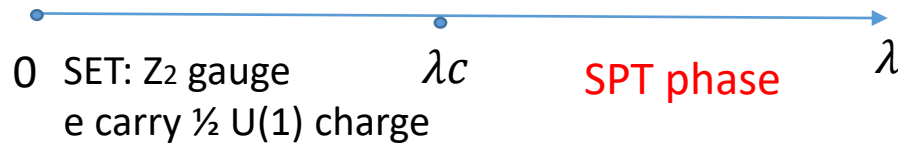
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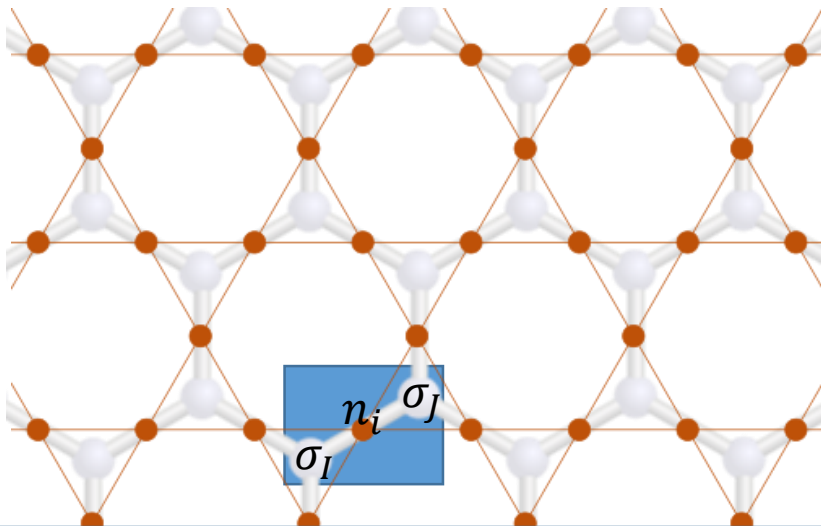
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One can analytically show:
SPT phase is realized when
 $t, h \ll \lambda \ll V$

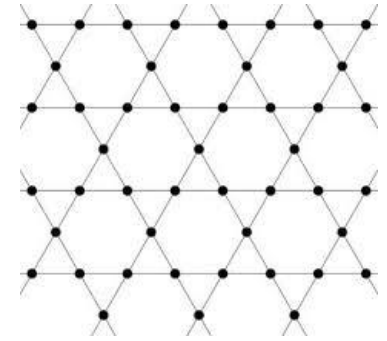
Summary

SPT: partial classification

$$H^{d+1}[SG, U(1)]$$

SG : on-site & spatial symmetries

T and P antiunitary



Numerical simulation

Input: symmetries of the model

Running Machinery

Output: SET/SPT classes and generic wavefunctions for every classes

Fractionalized phases

anyon condensation

SPT

Thank you!