

# Exploring quantum spin-frustrated materials

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J. Phys.: Condens. Matter 25, 026003 (2013)

Chem. Phys. Lett. 570, 37 (2013)

New J. of Phys. 16, 093011 (2014)

Scientific Reports 5, 16419 (2015)

Phys. Rev. B 93, 060405 (R) (2016)

Phys. Rev. B 94, 024438 (2016)

Phys. Rev. Lett. 115, 167203 (2015)

Phys. Rev. Lett. 117, 097201 (2016)

Phys. Rev. Lett. 117, 267202 (2016)

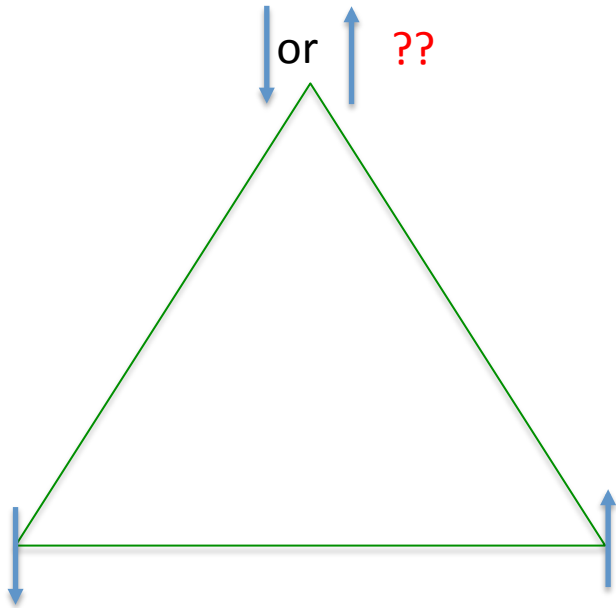
Phys. Rev. Lett. 118, 107202 (2017)

Nature 540, 559-562 (2016)

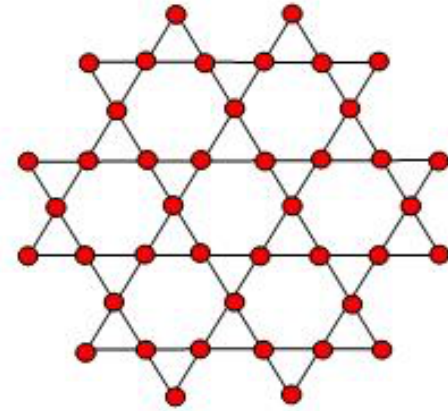
# Outline

- A short introduction
- $(\text{Zn,Co,Ni})\text{Cu}_3(\text{OH})_6\text{Cl}_2$ : two-dimensional kagome system
- $\text{ZnCu}_3(\text{OH})_6\text{SO}_4$ : anisotropic kagome system
- $\text{YbMgGaO}_4$ : structurally perfect triangular system
- Summary

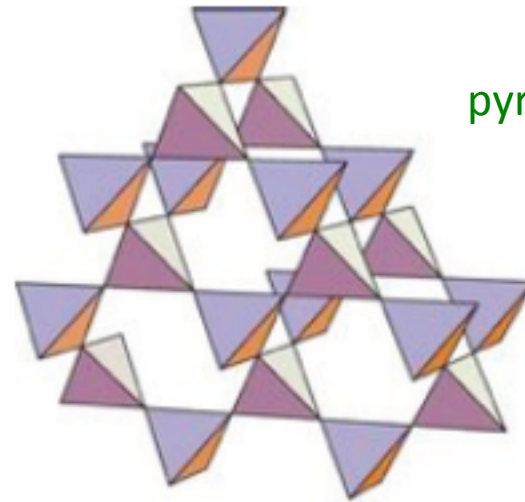
# Geometrical frustration



Building block for spin frustrated systems  
kagome, pyrochlore, spinel...



kagome



pyrochlore

# Quantum spin liquid—RVB state



P. W. Anderson, 1973

Resonating Valence Bond state

$$\Psi = \text{[Diagram 1]} + \text{[Diagram 2]} + \dots$$

The diagram shows two lattice configurations of blue ovals representing valence bonds. The first configuration shows a regular arrangement of ovals on the bonds of a triangular lattice. The second configuration shows a different arrangement, illustrating the concept of a resonating valence bond state.

$$\text{[Diagram 3]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

The diagram shows a single blue oval representing a singlet state.



X.-G. Wen

RVB state – an example of a topological phase

U(1) quantum spin liquid

Z<sub>2</sub> quantum spin liquid

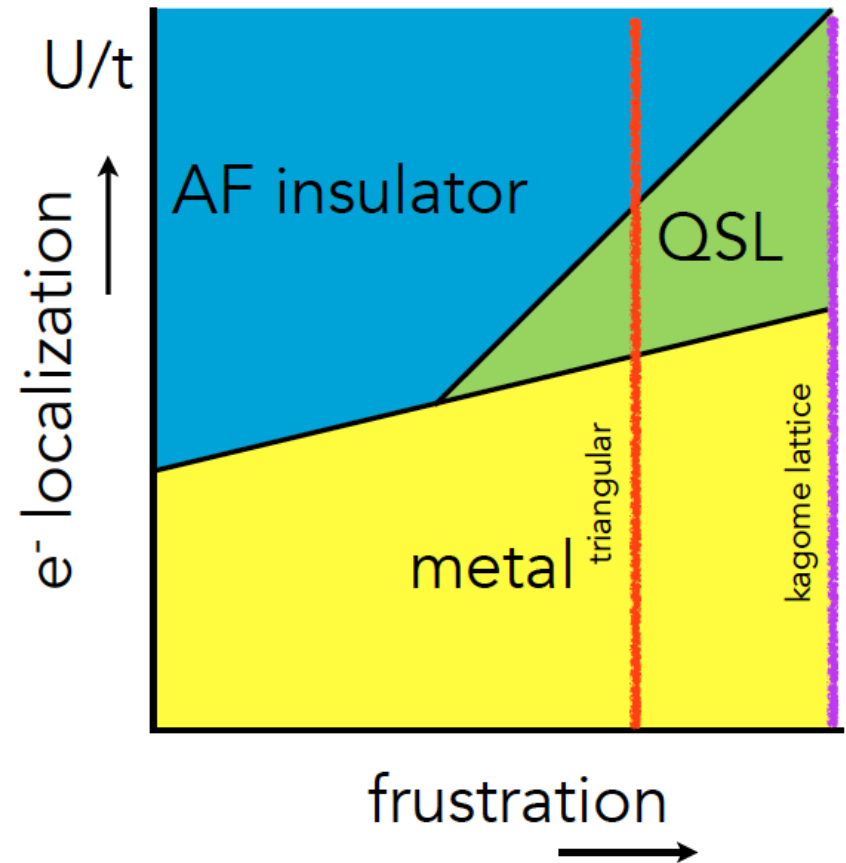
SU(2) quantum spin liquid

Buckley Prize winner 2017

# “Sheepskin Scroll”

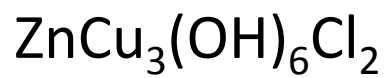
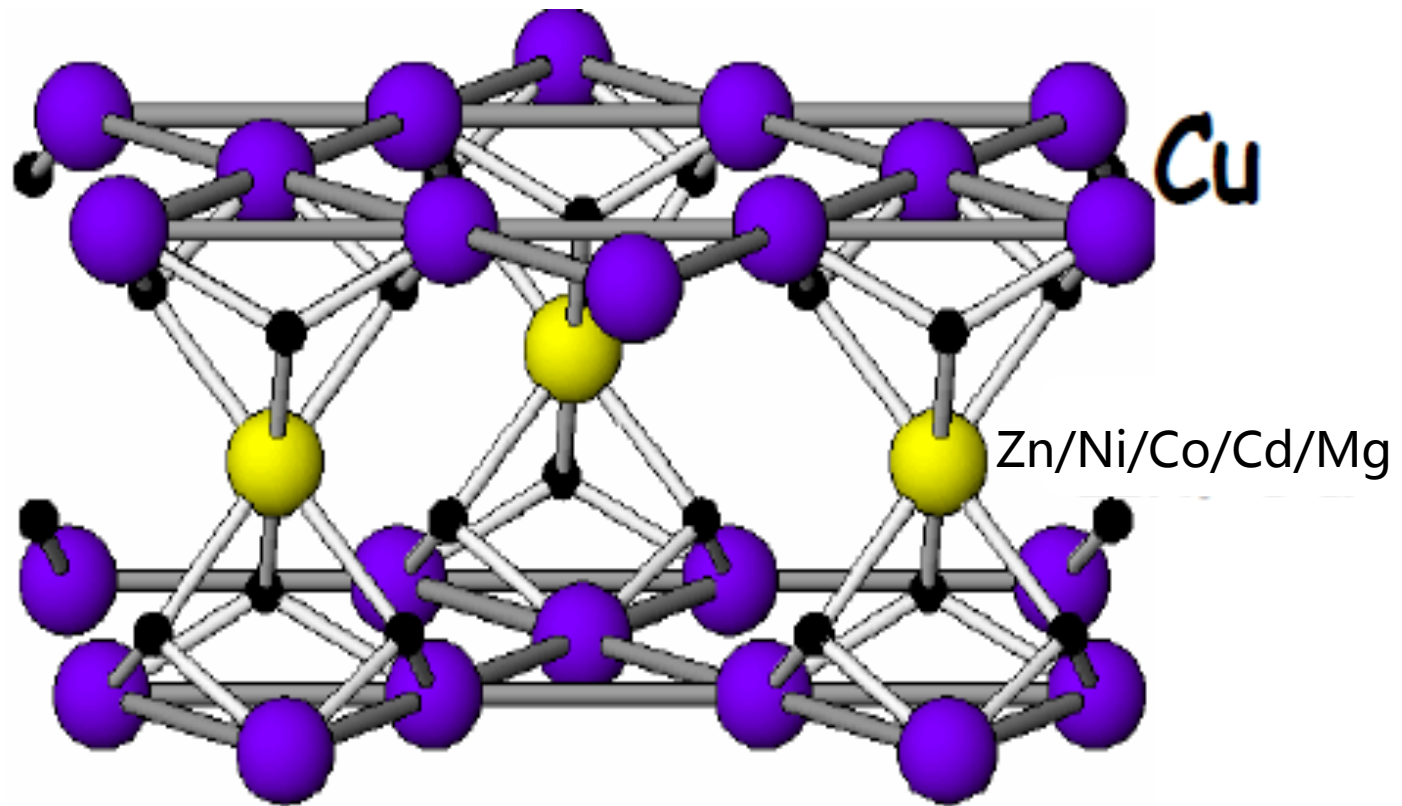
Where to find the mysterious “liquid” in the desert of condensed matter?

- $S=1/2$
- Antiferromagnetic coupling and frustration geomtry
- Large frustration factor  $f (= \Theta_W/T_N)$
- No long-range magnetic ordering
- No spontaneous symmetry breaking
- Strong charge fluctuation and/or exotic interactions
- Fractional excitation: spinon



From L. Balents

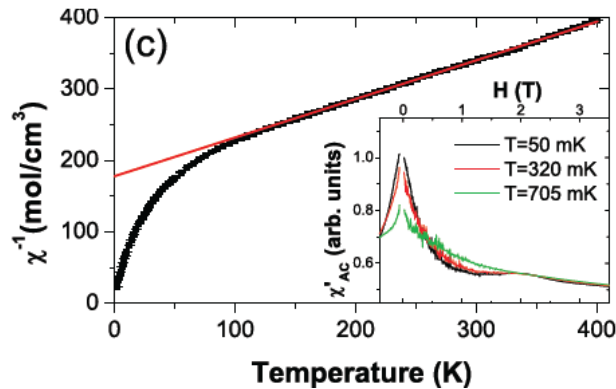
# Herbertsmithite



Antisite disorder/  
impurities

# Herbertsmithite

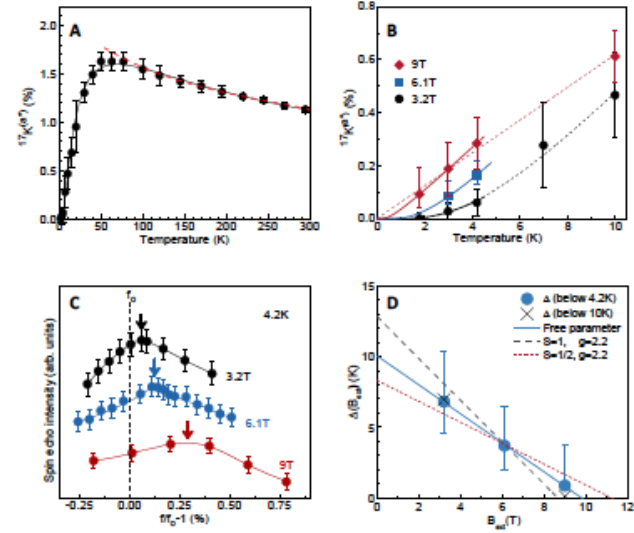
- AF coupling  $J \sim 200\text{K}$
- No magnetic ordering down to 50 mK
- Linear magnetic specific heat



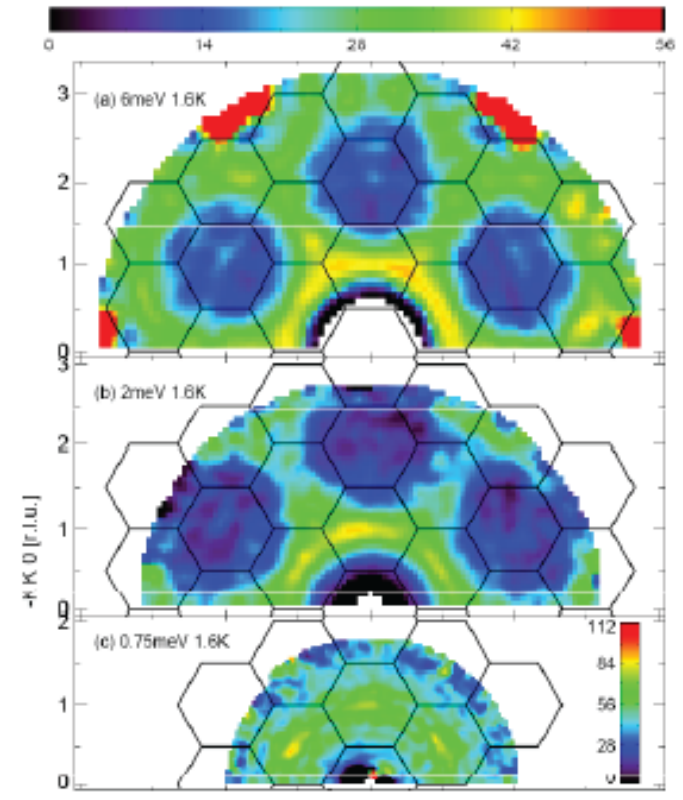
J. S. Helton et al., Phys. Rev. Lett. 98, 107204 (2007)

## Single crystal NMR

- Gapped spin liquid
- Gap  $\sim 10$  K



M. Fu et al., Science 350, 655 (2015)

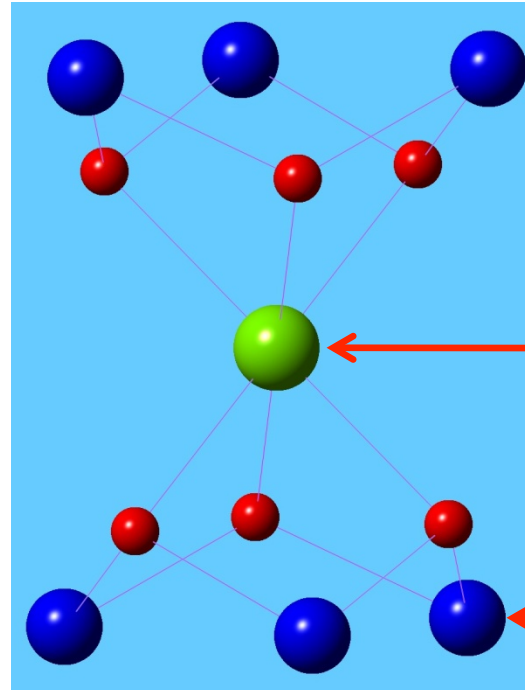
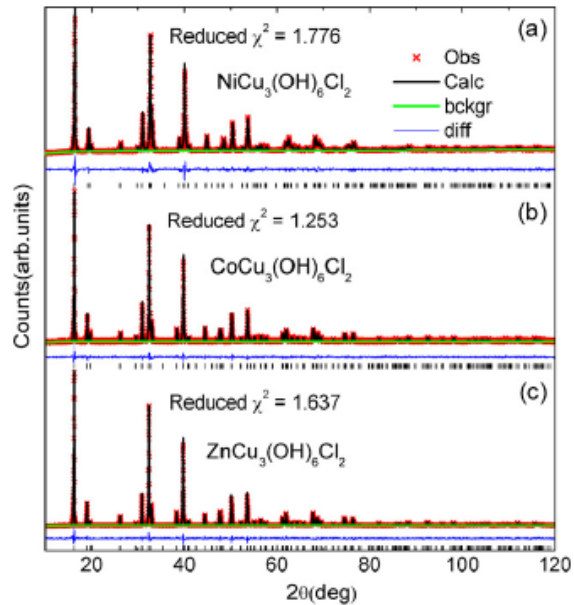


T. H. Han et al., Nature 492, 406 (2012)

## Continuum of spin excitations



# MCu<sub>3</sub>(OH)<sub>6</sub>Cl<sub>2</sub> (M=Zn, Co, Ni)



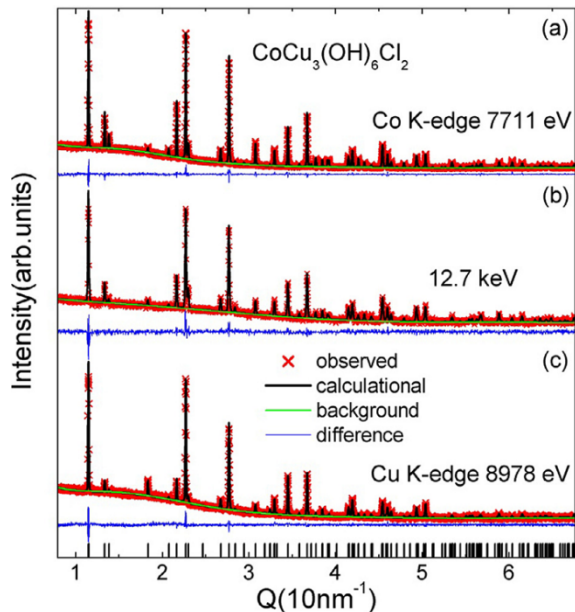
Interlayer/triangular position:

Co: 80(1)%  
Cu: 20(1)%

Intralayer/kagome position:

Cu: 93.6(1.0)%  
Co: 6.4(1.0)%

Significant site mixing

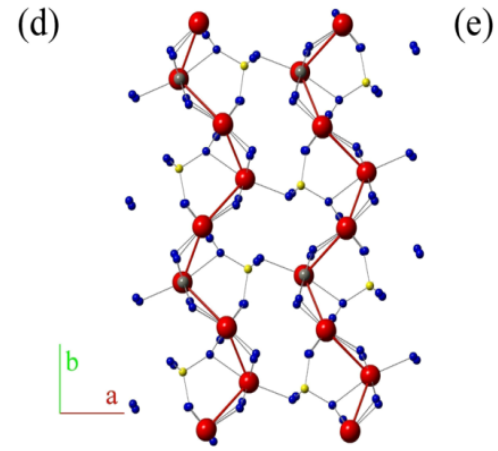
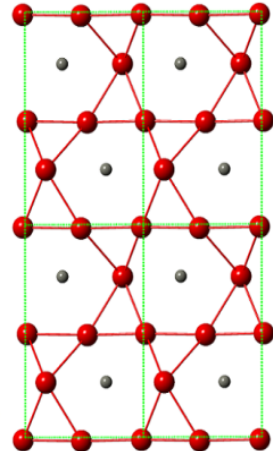
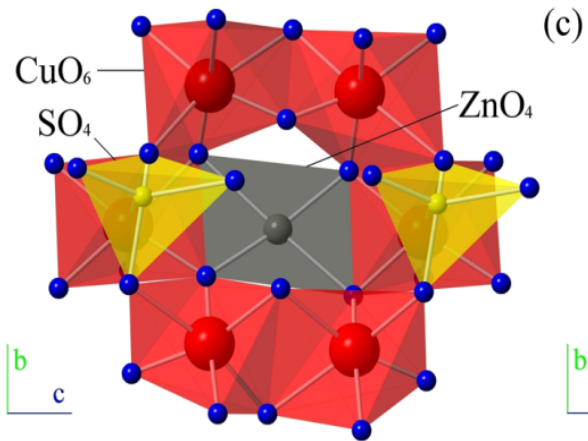
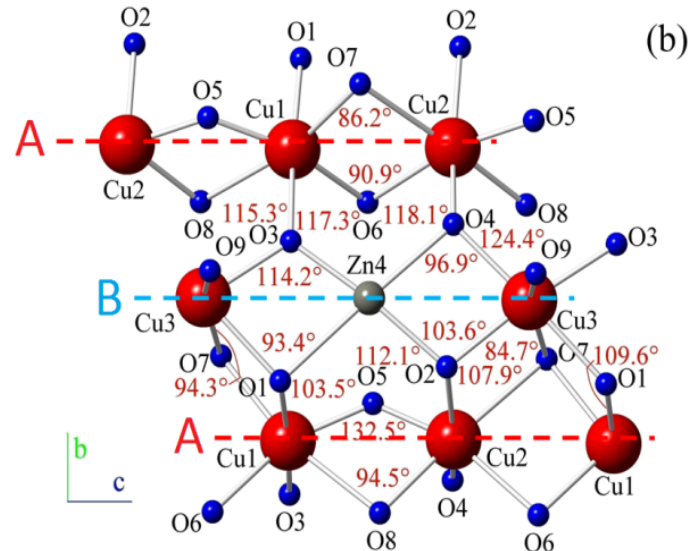
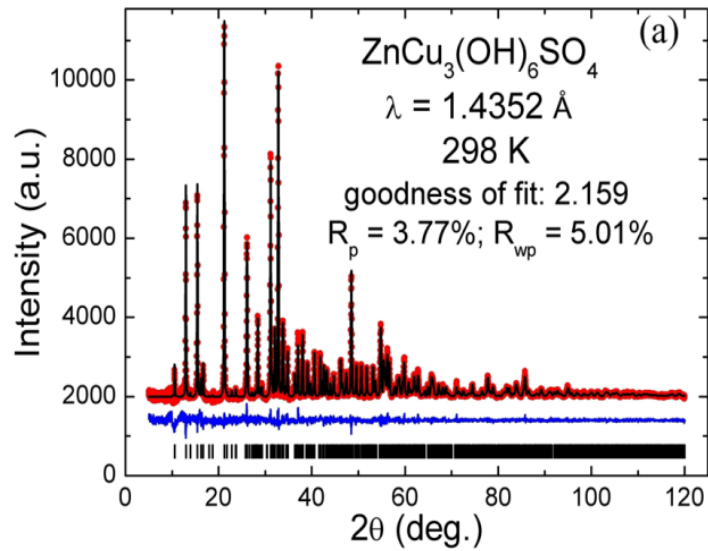


Combined structural refinements

Y. Li et al., J. Phys.: Condens. Matter 25, 026003 (2013)

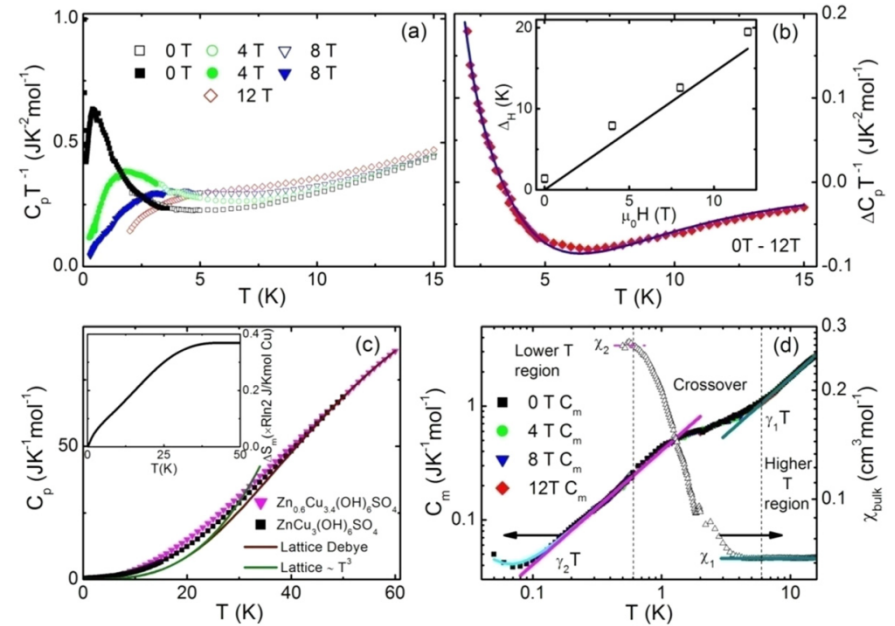
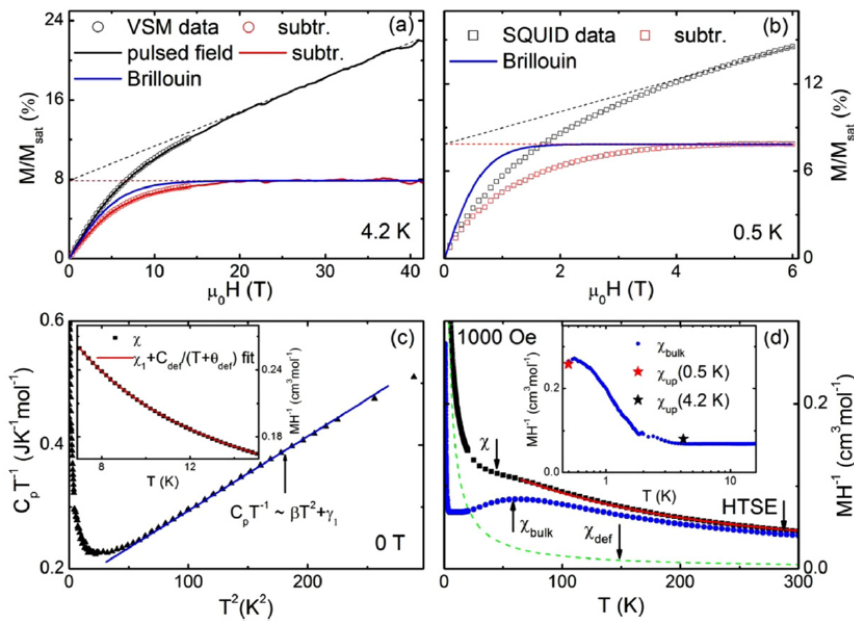
Y. Li et al., Chem. Phys. Lett. 570, 37 (2013)

# ZnCu<sub>3</sub>(OH)<sub>6</sub>SO<sub>4</sub>

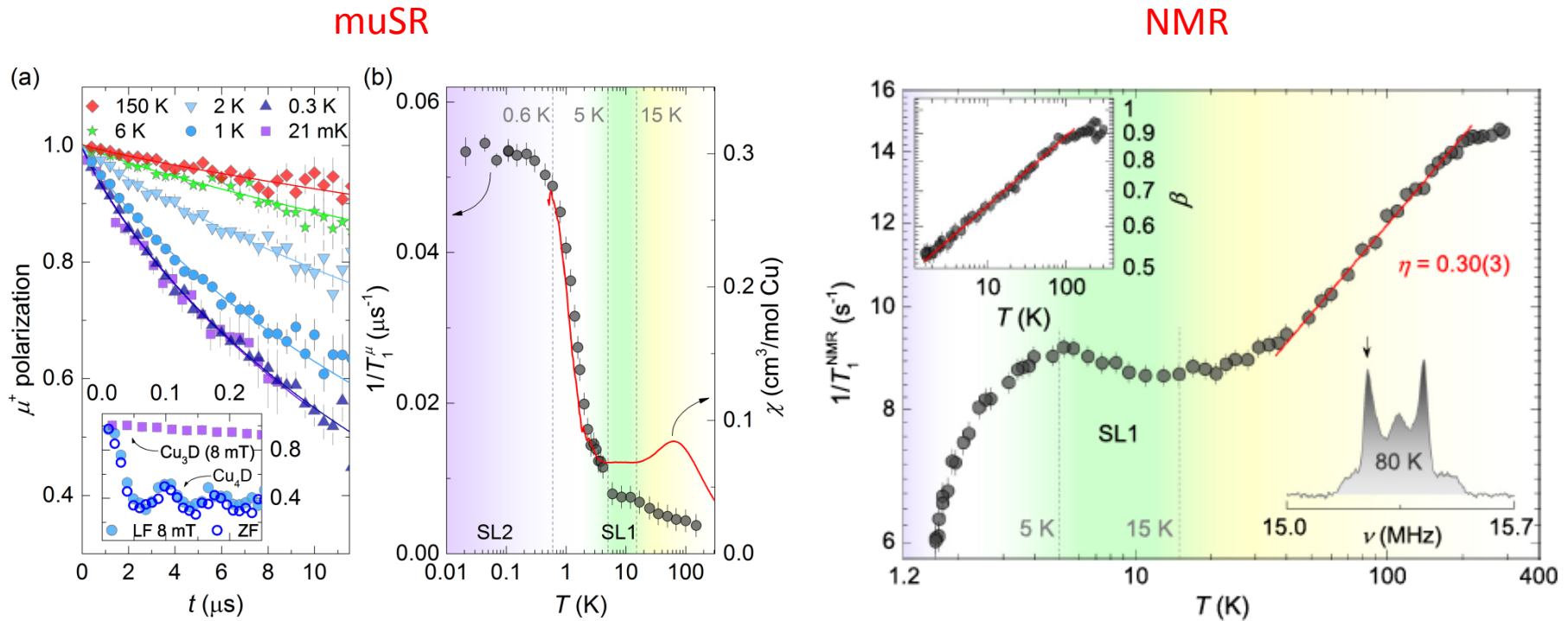


# Thermodynamics of $\text{ZnCu}_3(\text{OH})_6\text{SO}_4$

$\text{Zn}_x\text{Cu}_{4-x}$ $(\text{OH})_6\text{SO}_4$	Weiss tem- perature (K)	transition temperature (K)	frustration factor ( $l$ $\theta_w/T_c$ )	Curie constant ( $\text{Kcm}^3/$ $\text{mol}$ )	$\mu_{\text{eff}}$ ( $\mu_B/$ $\text{Cu}$ )	$g$
$x = 0$	-100	7.5	13	22.2	1.88	2.17
$x = 0.6$	-90	3.5	26	19.5	1.91	2.21
$x = 1$	-79	<0.05	>1580	17.1	1.90	2.20



# Spin dynamics of $\text{ZnCu}_3(\text{OH})_6\text{SO}_4$



- $\text{ZnCu}_3(\text{OD})_6\text{SO}_4$
- No oscillation down to 21 mK
- Two regions?

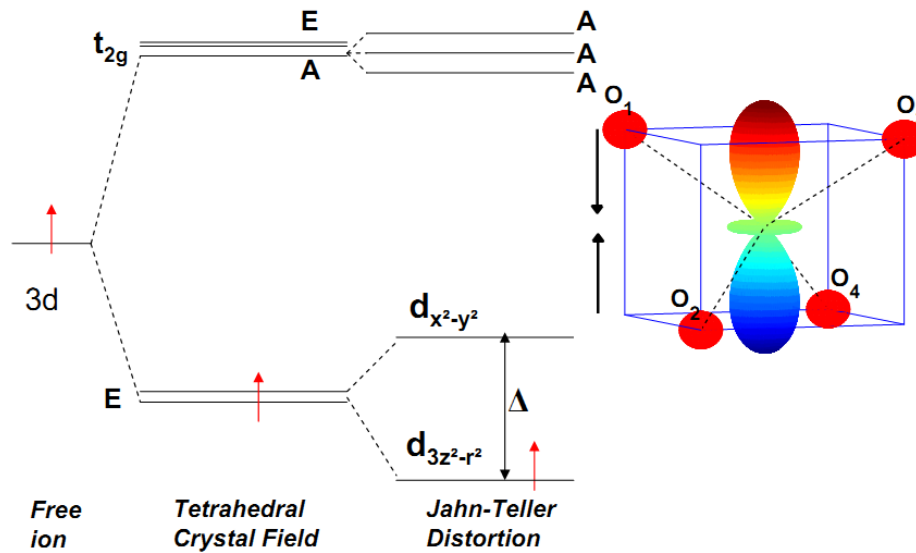
Sign of two regions

M. Gomilsek et al., Phys. Rev. B 93, 060405 (2016) (Rapid Commun.)

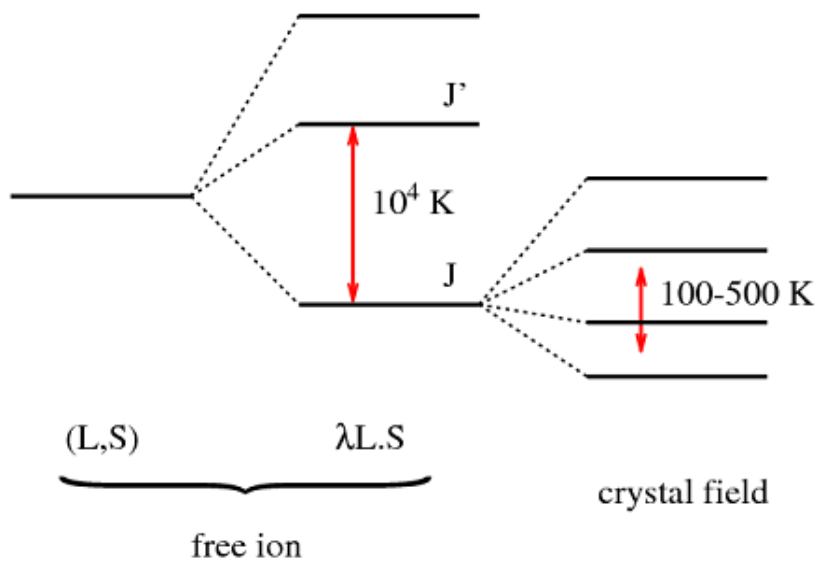
M. Gomilsek et al., Phys. Rev. B 94, 024438 (2016)

# 3d vs 4f

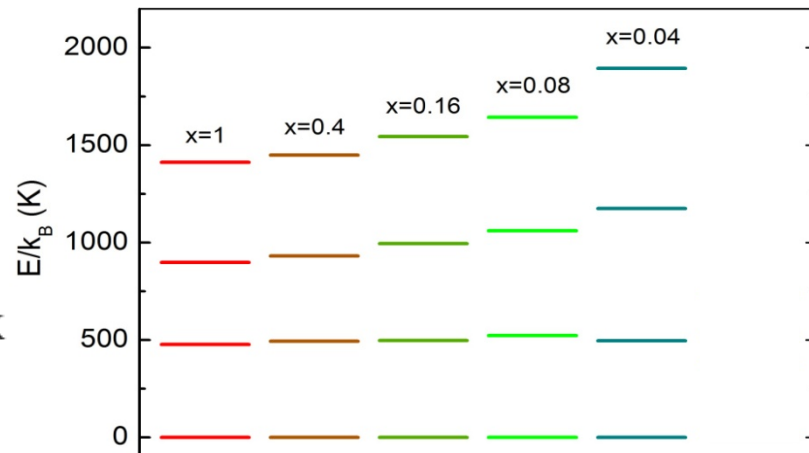
$\text{Cu}^{2+}$



$\text{Yb}^{3+}$



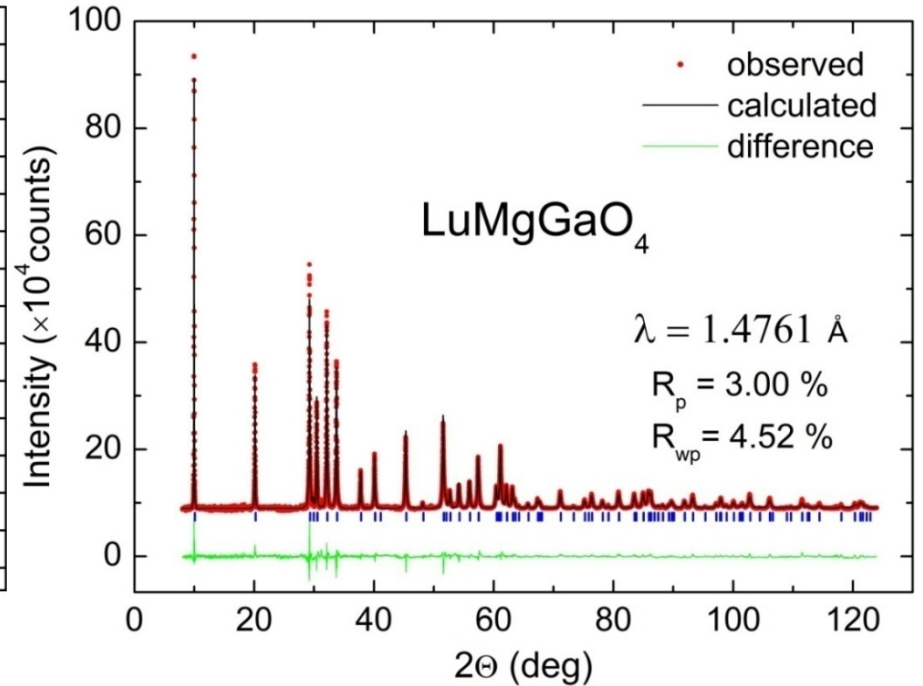
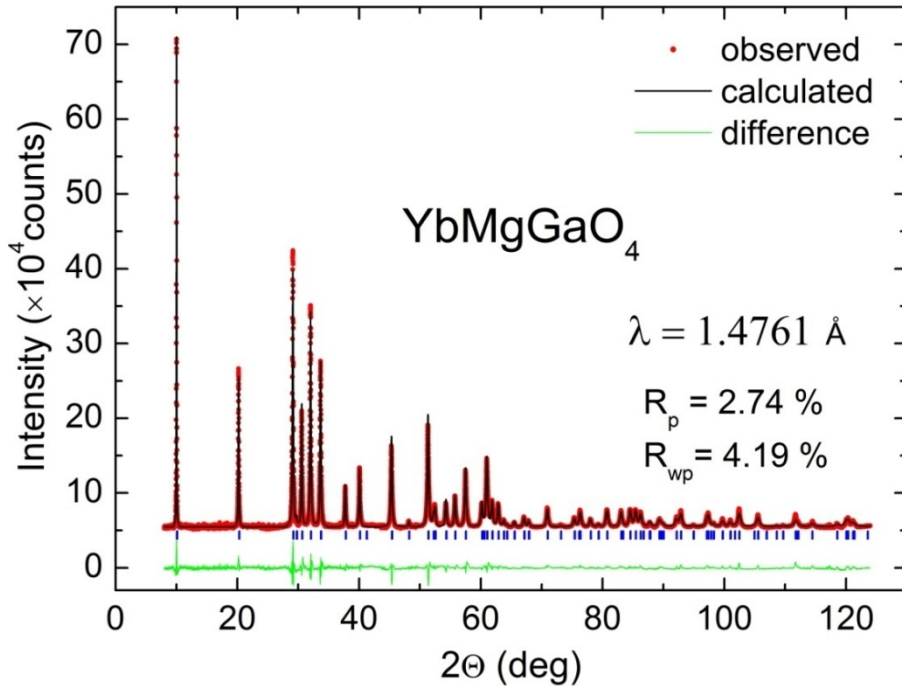
Kramers doublets



$4f^{13}$   $L=3, S=1/2, J=7/2$



# Powder diffraction

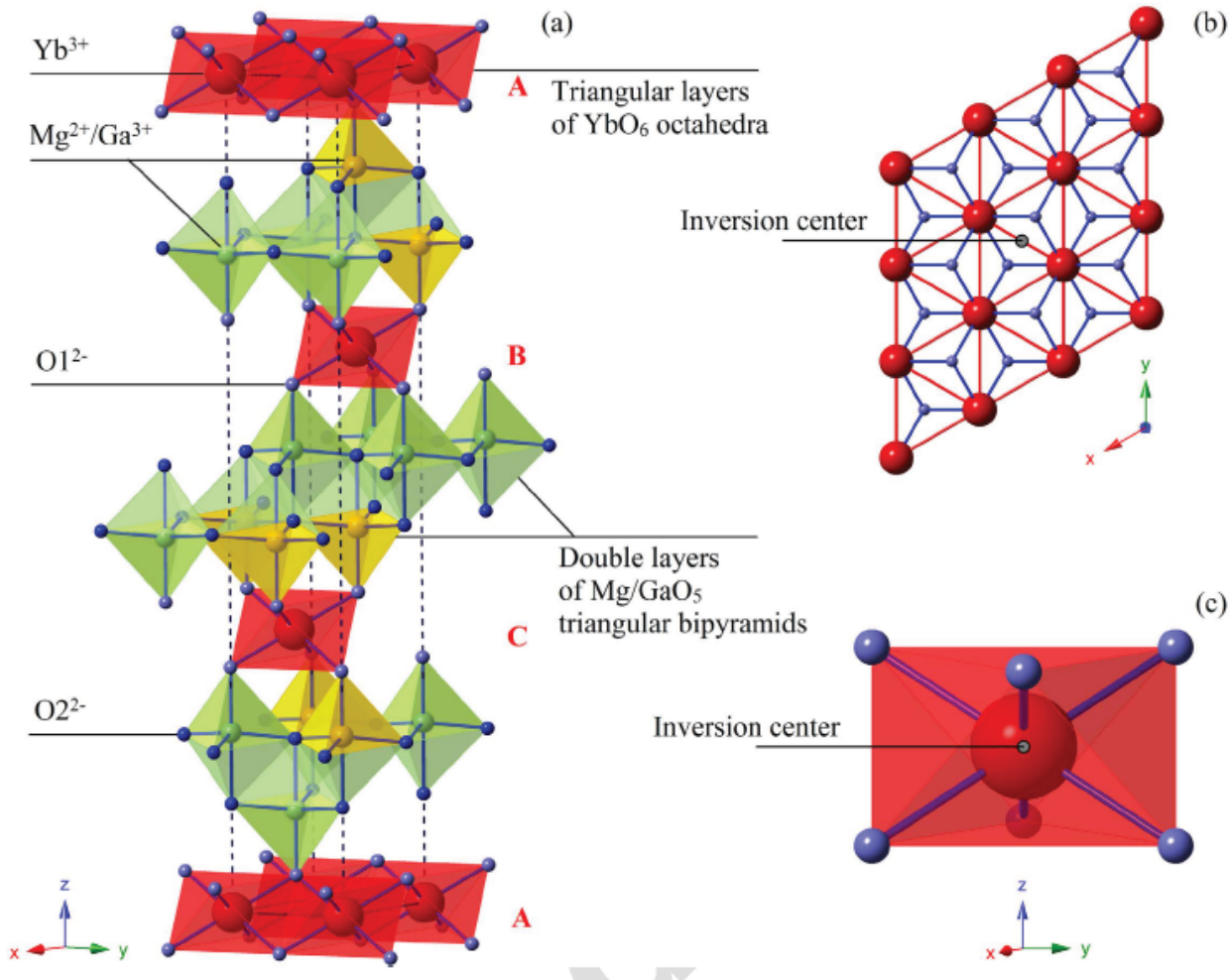


Yb<sup>3+</sup>: 4f<sup>13</sup>, S=1/2  
m=173

Lu<sup>3+</sup>: 4f<sup>14</sup>, S=0  
m=175

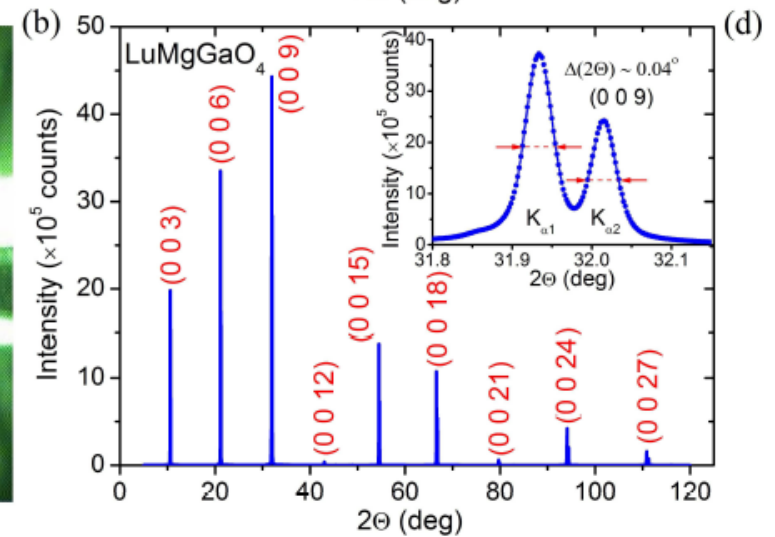
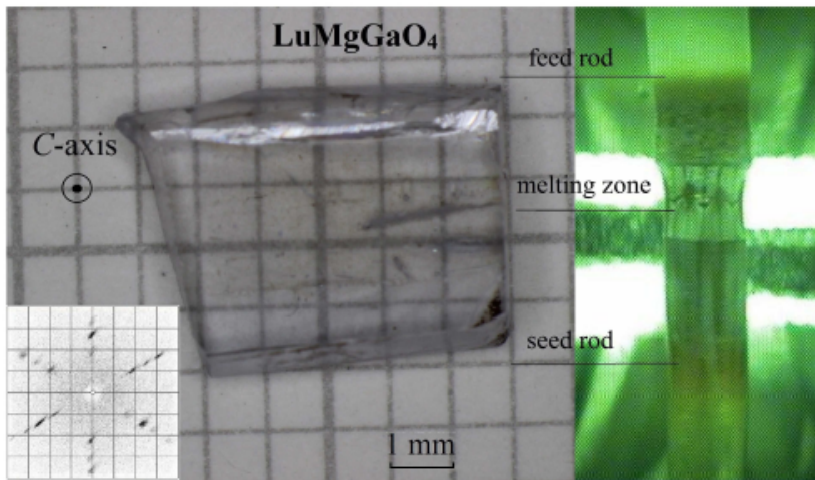
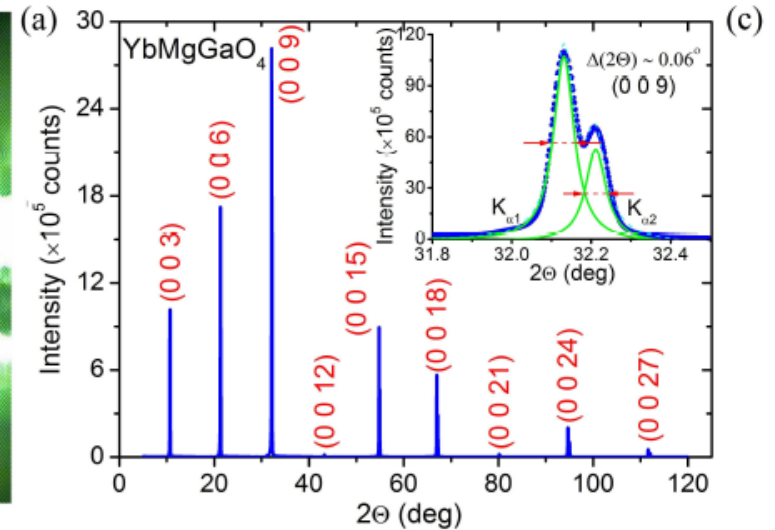
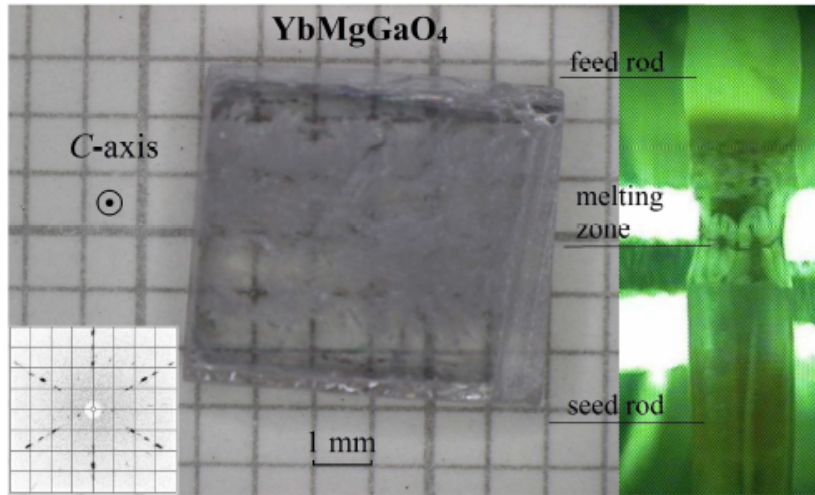
Lanthanides	6	58 <b>Ce</b> Cerium 140.115	59 <b>Pr</b> Praseodymium 140.9076	60 <b>Nd</b> Neodymium 144.24	61 <b>Pm</b> Promethium (145)	62 <b>Sm</b> Samarium 150.36	63 <b>Eu</b> Europium 151.965	64 <b>Gd</b> Gadolinium 157.25	65 <b>Tb</b> Terbium 158.9253	66 <b>Dy</b> Dysprosium 162.50	67 <b>Ho</b> Holmium 164.9303	68 <b>Er</b> Erbium 167.26	69 <b>Tm</b> Thulium 168.9342	70 <b>Yb</b> Ytterbium 173.04	71 <b>Lu</b> Lutetium 174.967	
Actinides	7	90 <b>Th</b> Thorium 232.0381	91 <b>Pa</b> Protactinium 231.0359	92 <b>U</b> Uranium 238.0289	93 <b>Np</b> Neptunium (237)	94 <b>Pu</b> Plutonium (244)	95 <b>Am</b> Americium (243)	96 <b>Cm</b> Curium (247)	97 <b>Bk</b> Berkelium (247)	98 <b>Cf</b> Californium (251)	99 <b>Es</b> Einsteinium (252)	100 <b>Fm</b> Fermium (257)	101 <b>Md</b> Mendelevium (258)	102 <b>No</b> Nobelium (259)	103 <b>Lr</b> Lawrencium (260)	7

# Structurally perfect $\text{YbMgGaO}_4$



- Strict R-3m symmetry
- Two-dimensional
- $S_{\text{eff}}=1/2$
- No antisite/impurity
- No DM interaction
- Perfect reference sample  $\text{LuMgGaO}_4$
- High-quality crystals available
- Small J

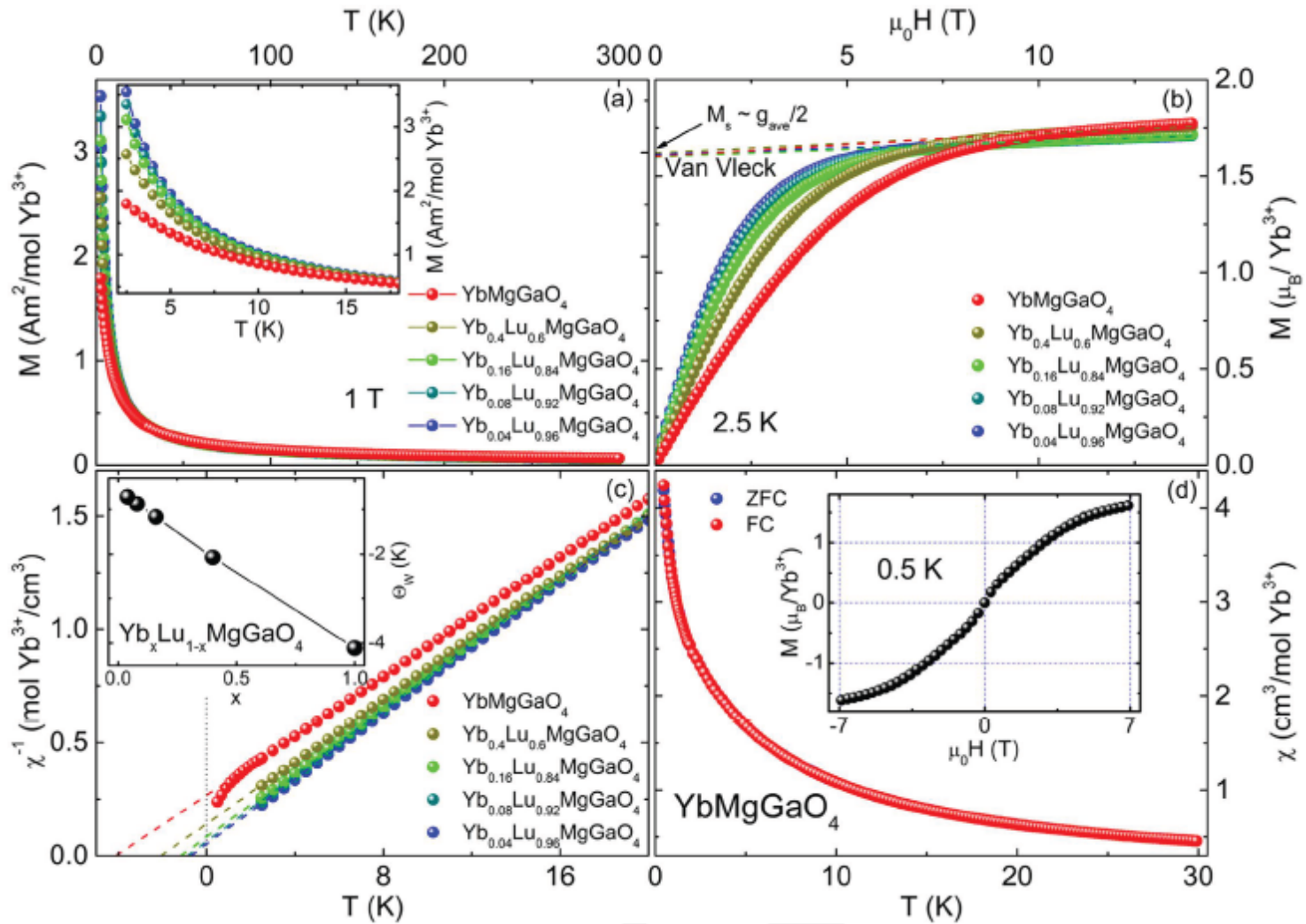
# Single Crystals



Charge Gap  $\sim 4$  eV



# Magnetization



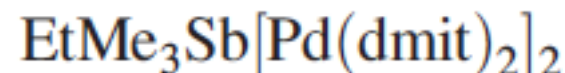
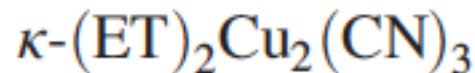
$$\Theta_w \sim -4 \text{ K}$$

# Spin Hamiltonian

Under R-3m symmetry

$$\mathcal{H} = \sum_{\langle ij \rangle} [J_{zz} S_i^z S_j^z + J_{\pm} (S_i^+ S_j^- + S_i^- S_j^+) + J_{\pm\pm} (\gamma_{ij} S_i^+ S_j^+ + \gamma_{ij}^* S_i^- S_j^-) - \frac{iJ_{z\pm}}{2} (\gamma_{ij}^* S_i^+ S_j^z - \gamma_{ij} S_i^- S_j^z + \langle i \leftrightarrow j \rangle)],$$

Anisotropic terms

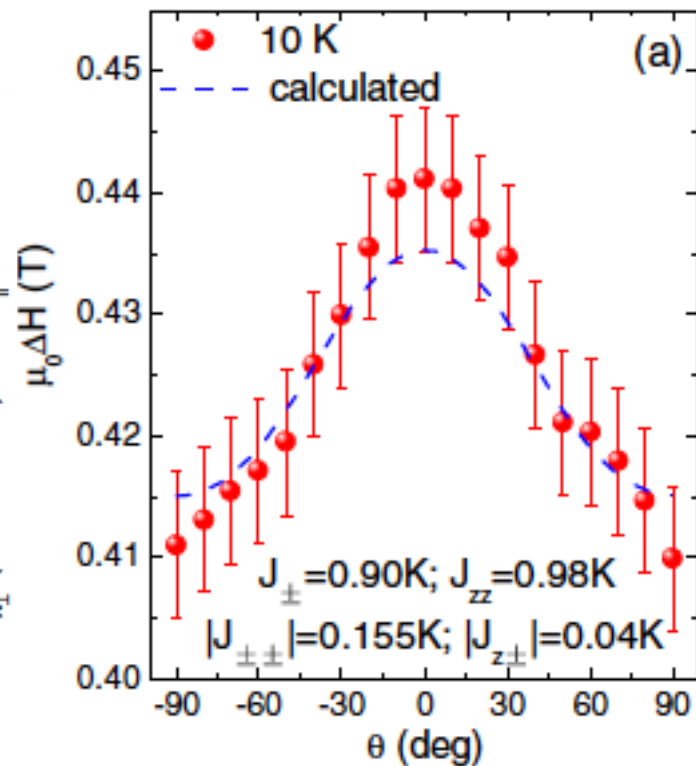
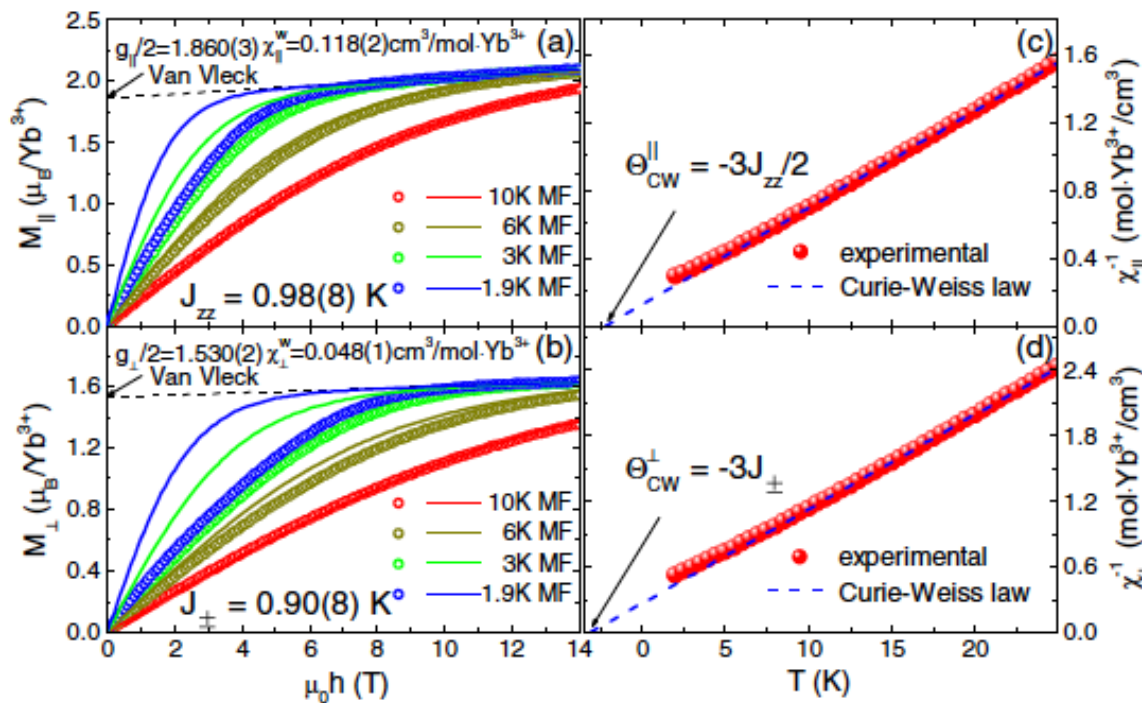


Strong charge fluctuations



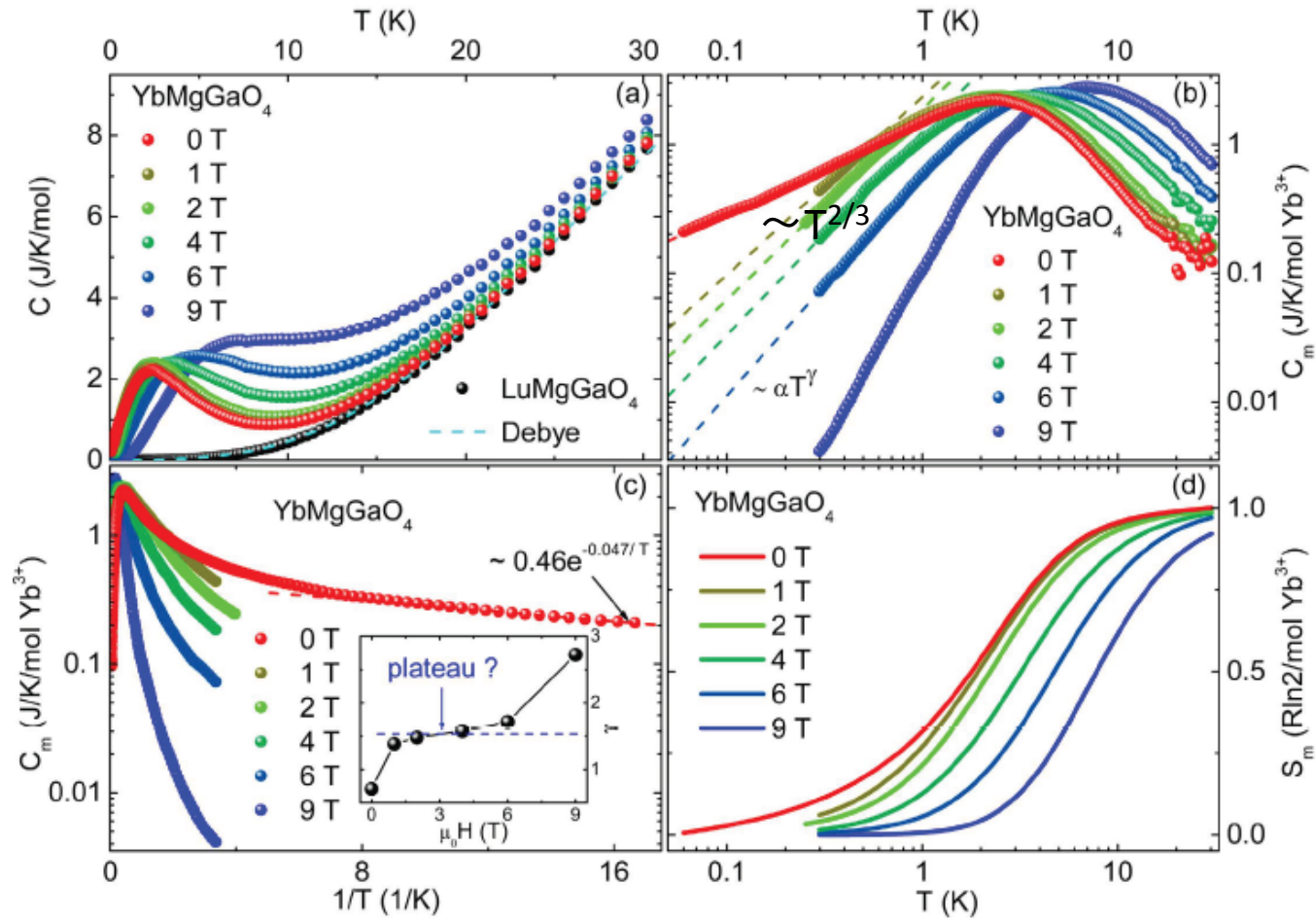
Ring exchange

# Exchange parameters



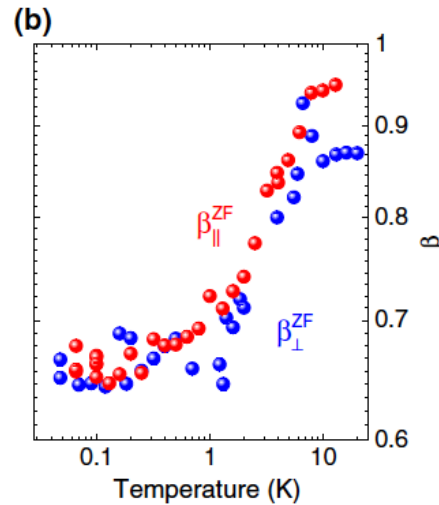
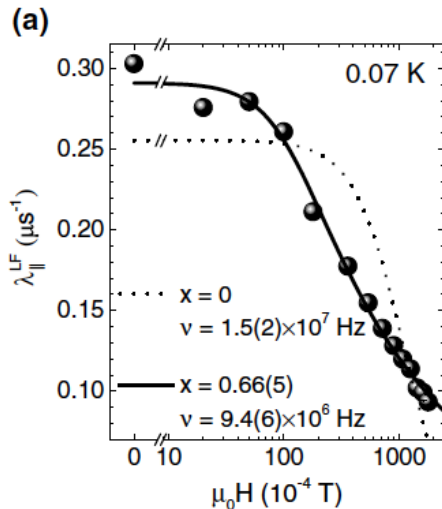
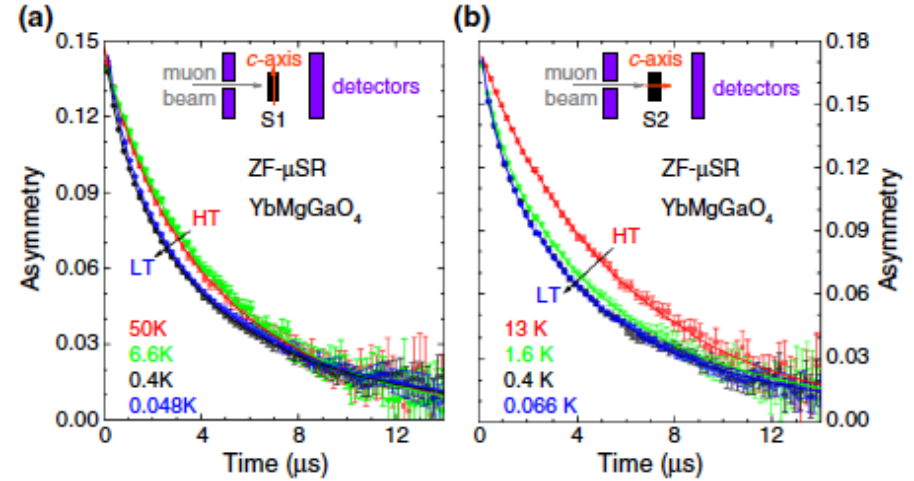
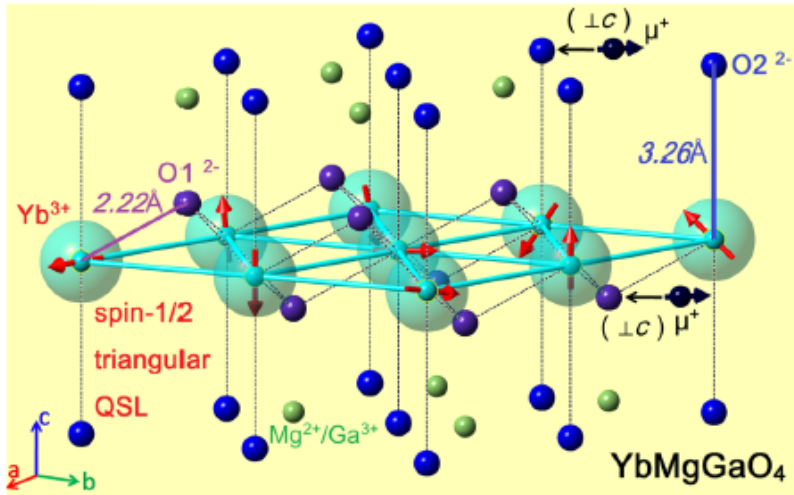
# Specific Heat

$m_{\text{Yb}}=173$   
 $m_{\text{Lu}}=175$



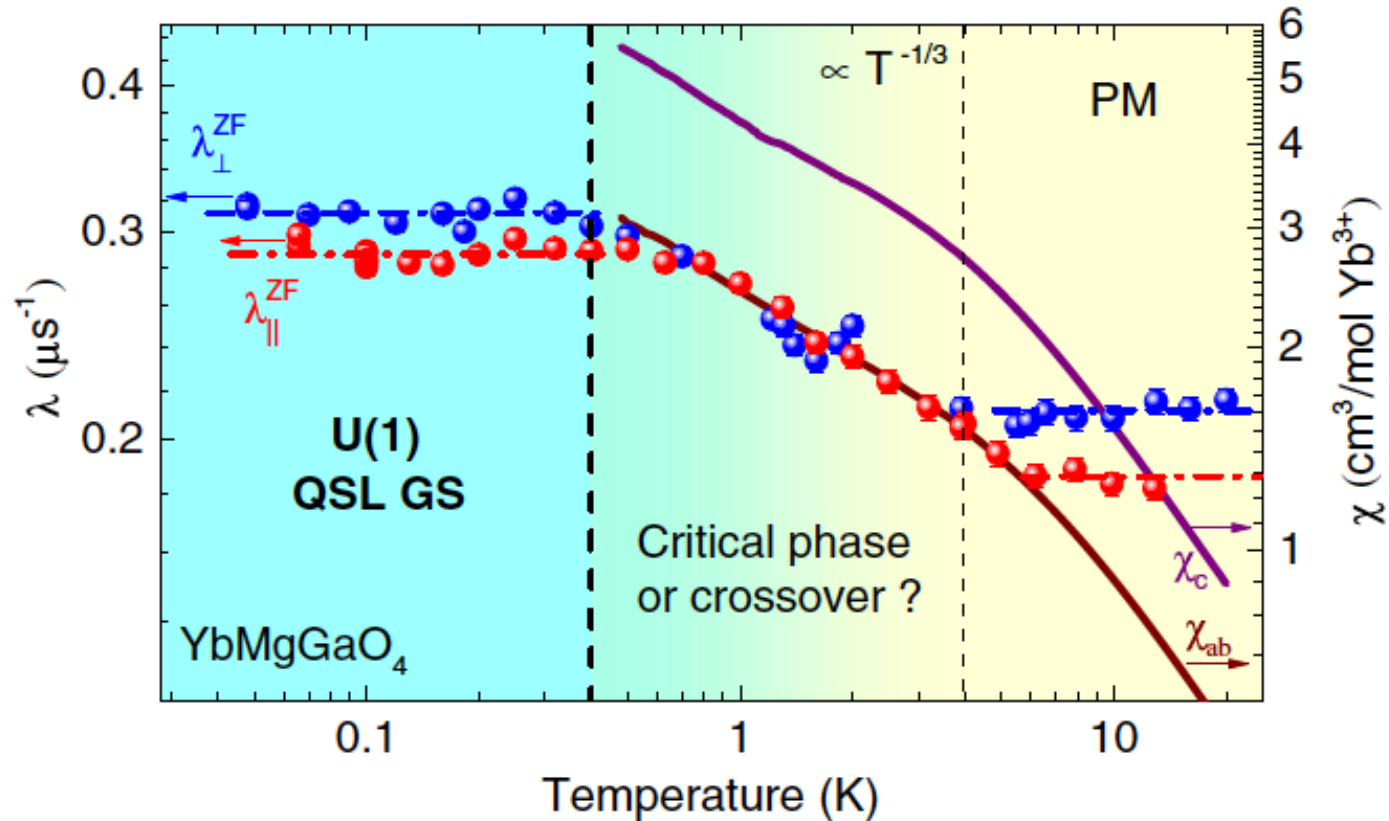
- No magnetic ordering down to 50 mK
- Accurate spin entropy
- Zero-entropy spin ground state (residual spin entropy below 50 mK  $< 0.6\%$ )

# U(1) quantum spin liquid ground state

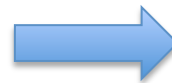


Rule out long-range magnetic ordering and spin freezing

# U(1) quantum spin liquid ground state

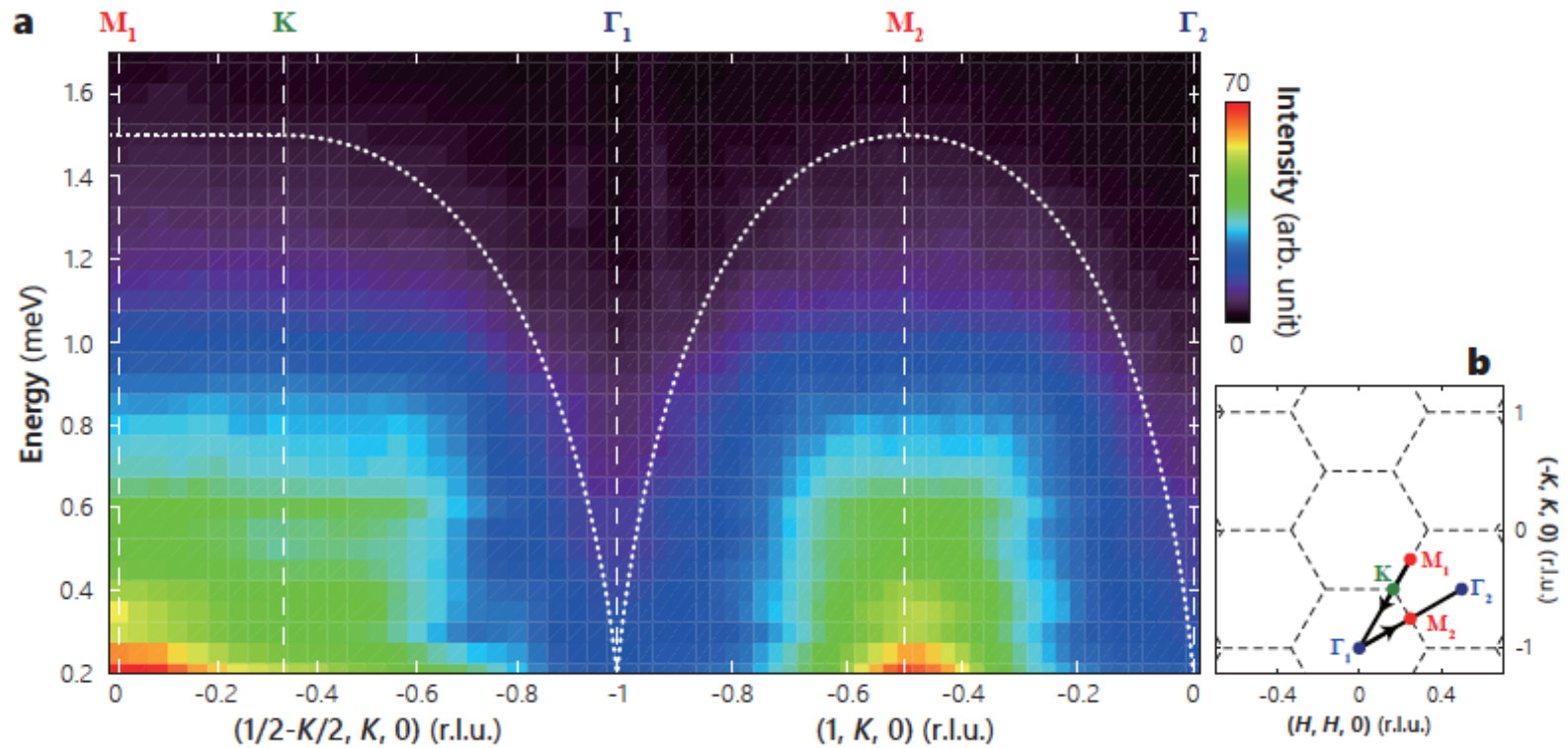


Zero spin entropy: spin ground state  
 $C_V \sim T^{2/3}$   
 Constant Muon spin relaxation rate



U(1) gapless quantum spin liquid

# Spin excitations



- Diffusive spin excitations
- Spinon Fermi surface
- Gapless U(1) quantum spin liquid

Y. Shen et al., Nature 540, 559-562 (2016)



# Summary

- ◇ A new two-dimensional triangular compound  $\text{YbMgGaO}_4$  with  $S_{\text{eff}}=1/2$  and high-quality single crystals available
- ◇ Anisotropic spin Hamiltonian
- ◇ Gapless  $U(1)$  quantum spin liquid ground state
- ◇ Strong spin-orbit coupling may play a key role in forming QSL ground state and our work may inspire a new route to search for QSL materials