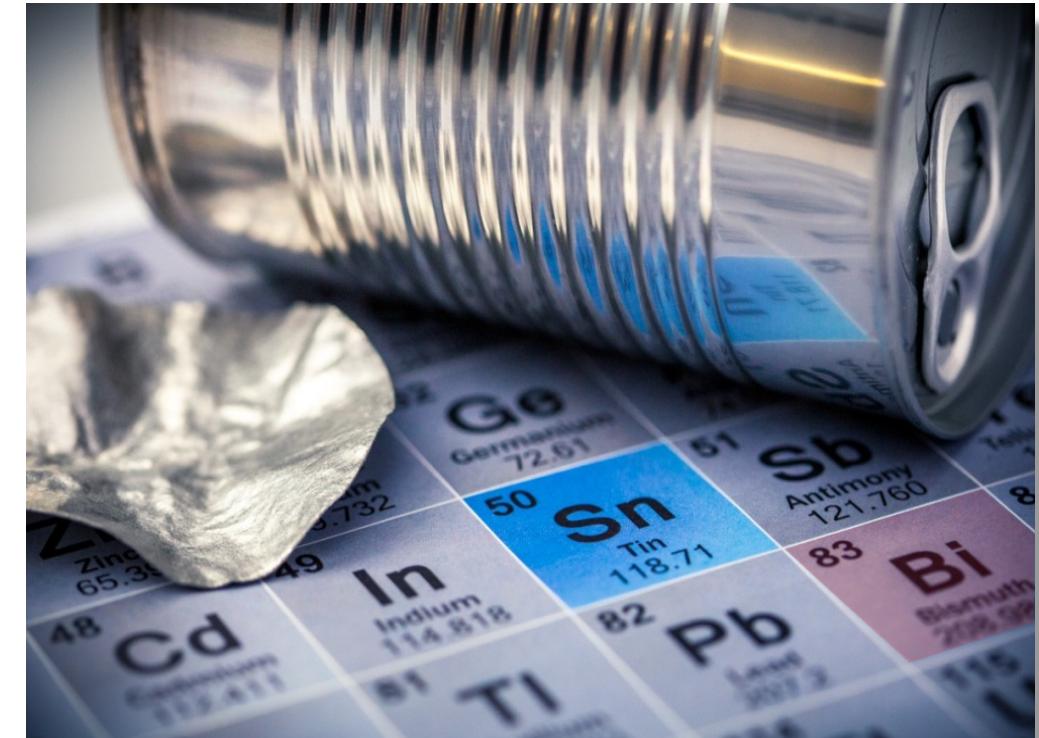


Elemental and elementary fun with topological insulators: α -Sn, stanene, bismuthene

Ralph Claessen

*Physikalisches Institut and
Röntgen Center for Complex Materials (RCCM),
Universität Würzburg, Germany*

- Sn as 3D and 2D TI
- bismuthene/SiC: a quantum spin Hall paradigm
- electronic correlations in 1D edge states



Jörg Schäfer

Felix Reis

Raul Stühler

Florian Adler

Arne Barfuss

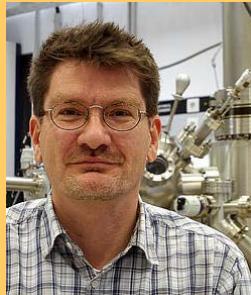
Max Bauernfeind

Lenart Dudy

Stefan Glass

Victor Rogalev

Markus Scholz



experiments

Gang Li (now Shanghai Tech)

Werner Hanke

Ronny Thomale

Domenico Di Sante

Andrzej Fleszar

Congjun Wu (UCSD)

Ewelina Hankiewicz

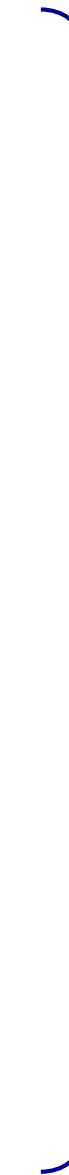
Fernando Dominguez

Dimitri Jungblut

Benedikt Scharf



theory



(almost) all @ U Würzburg



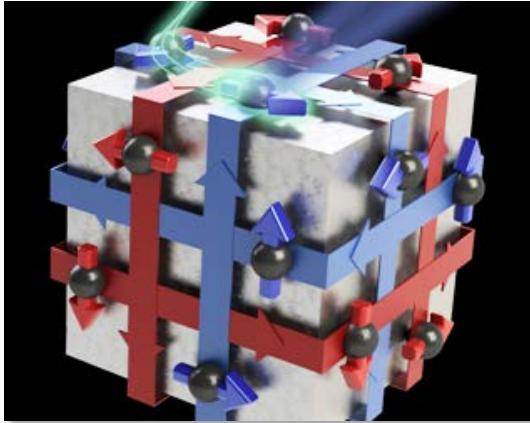
SFB1170
ToCoTronics*

*Topological and correlated electronics
at surfaces and interfaces

Helical surface/edge states in topological insulators

3D TIs

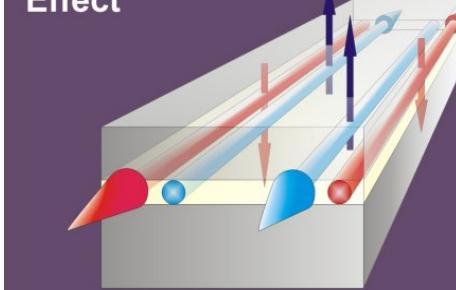
with 2D surface states



2D TIs

with 1D edge states

Quantum Spin Hall
Effect



"topologically protected" surface/edge states in the bulk band gap:

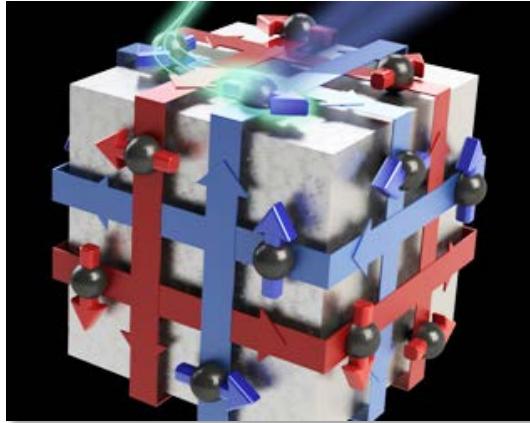
- metallic
- spin-polarized
- helical

→ search for Dirac materials & quantum spin Hall (QSH) insulators for
electronic/spintronic applications at room temperature

Helical surface/edge states in topological insulators

3D TIs

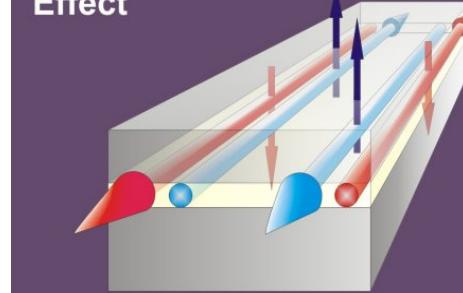
with 2D surface states



2D TIs

with 1D edge states

Quantum Spin Hall
Effect



"topologically protected" surface/edge states in the bulk band gap:

- metallic
- spin-polarized
- helical

→ search for Dirac materials & quantum spin Hall (QSH) insulators for
electronic/spintronic applications at room temperature

theoretical proposal:

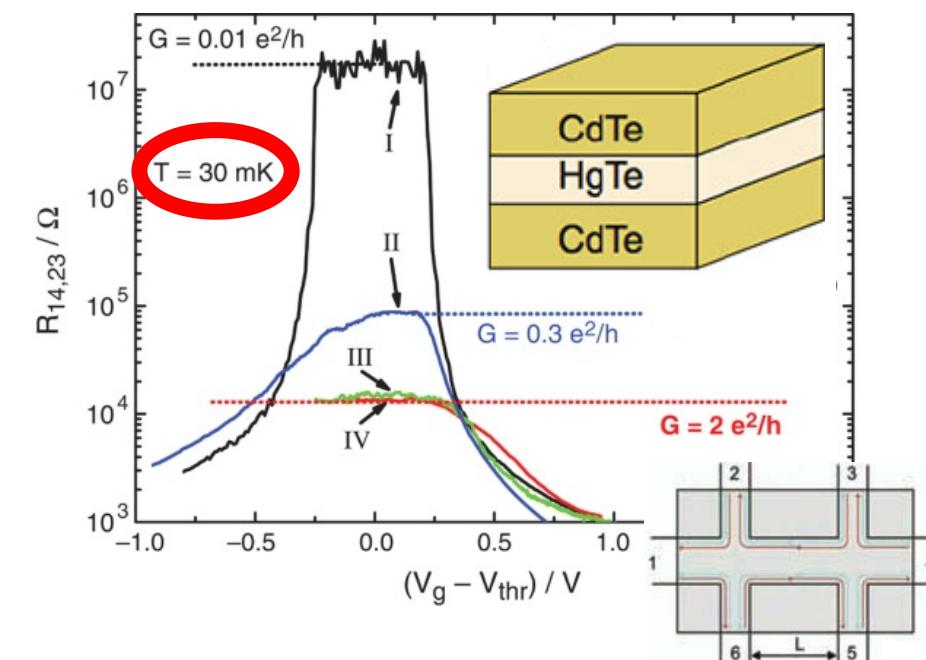
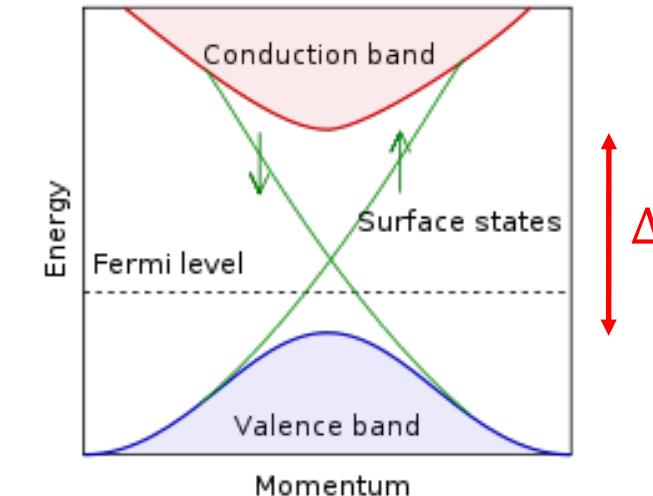
semiconductor quantum wells with inverted band gap

*Bernevig, Hughes & Zhang, Science (2006)***experimental realizations:**

- HgTe/CdTe quantum well structures

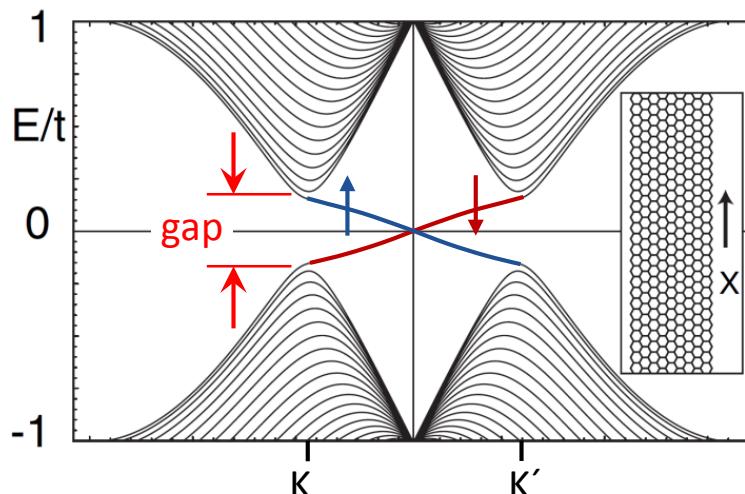
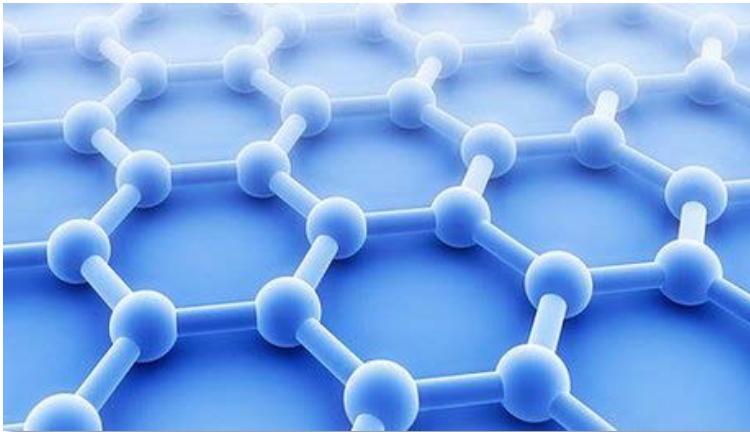
M. König et al., Science (2007)→ effective band gap: $\Delta < 40 \text{ meV}$

- InAs/GaSb QWs

Knez, Du & Sullivan, PRL (2011)→ $\Delta \approx 4 \text{ meV}$ 

Quantum spin Hall effect in honeycomb lattices

Kane & Mele, PRL (2005): QSHE in graphene



SOC-induced
energy gap:

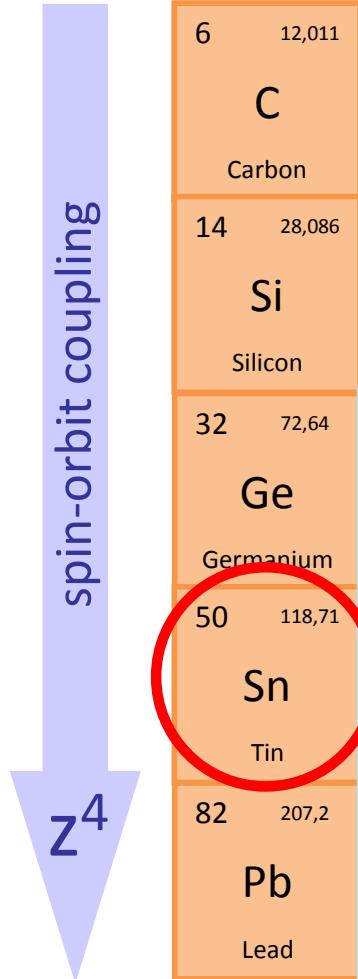
24 μ eV
graphene

2 meV

24 meV

100 meV
stanene

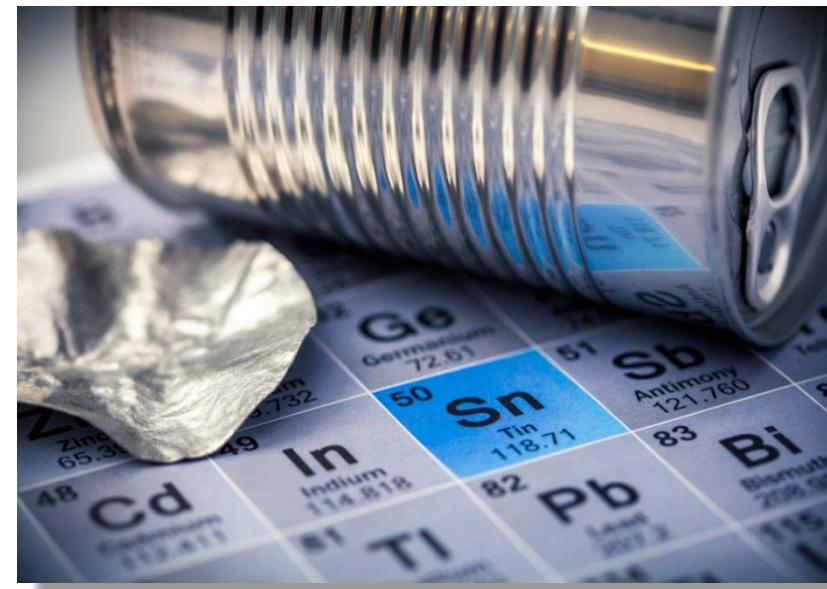
Y. Xu et al.,
PRL 2013



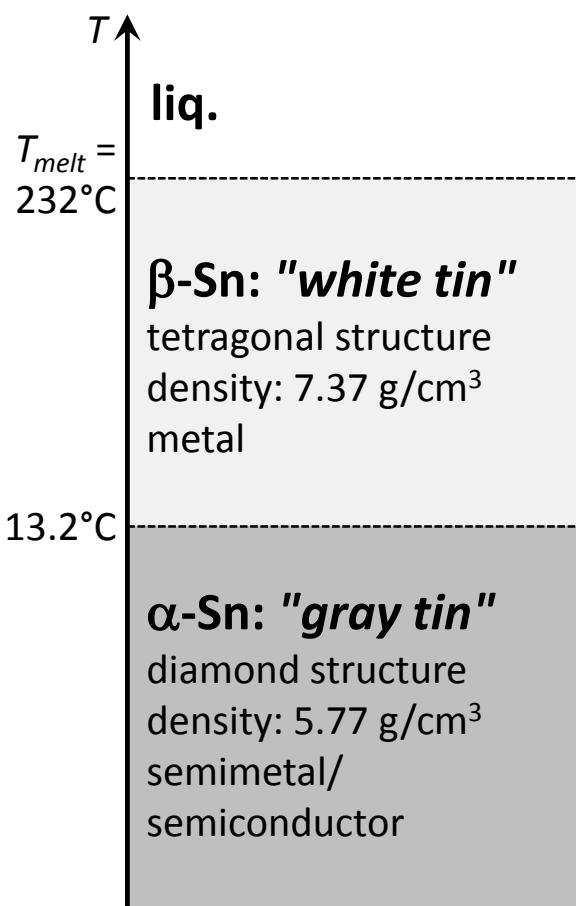
→ spin-polarized (helical) metallic edge states



Elemental Sn as 3D and 2D topological insulator

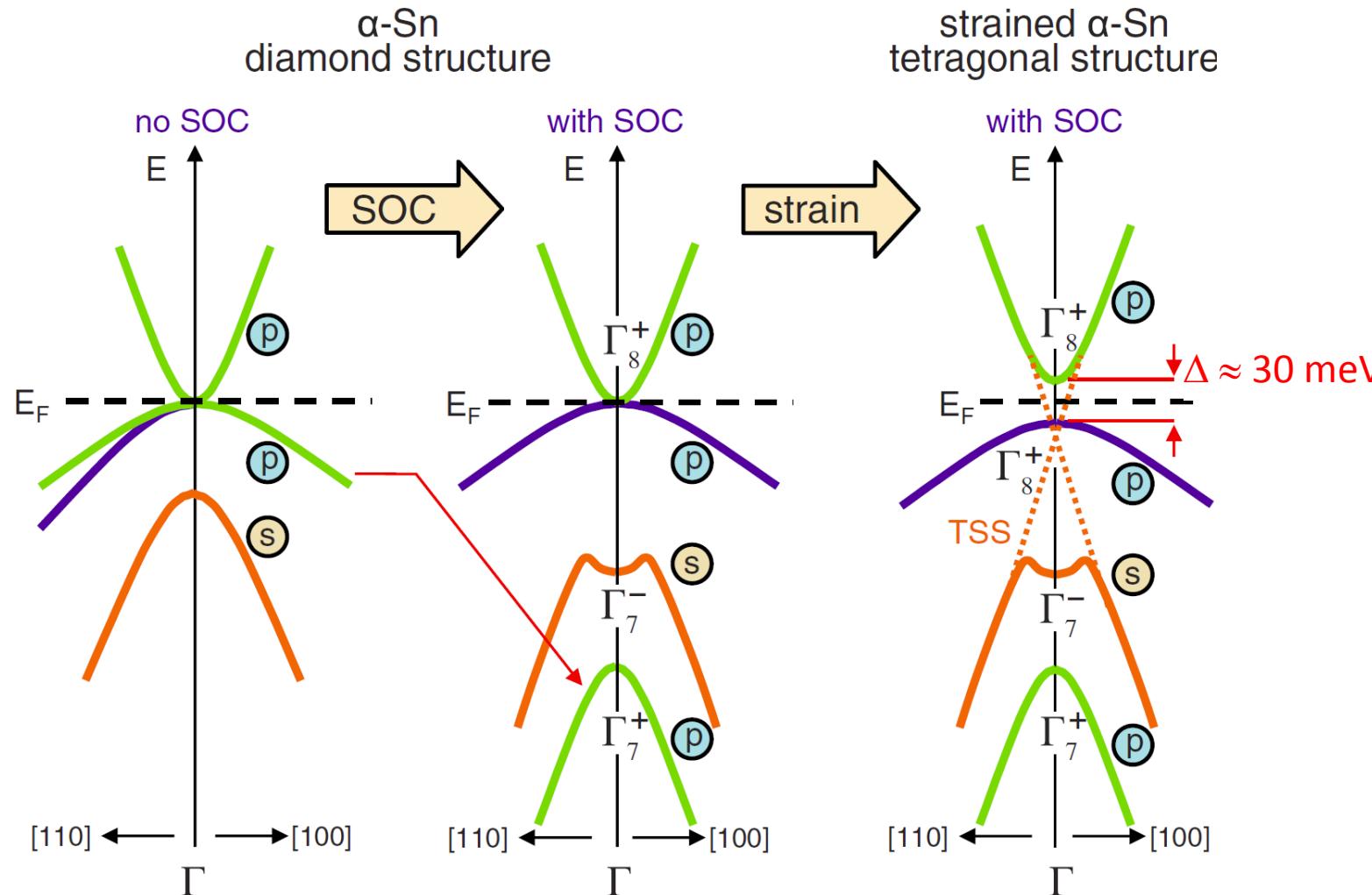


Sn as elemental TI

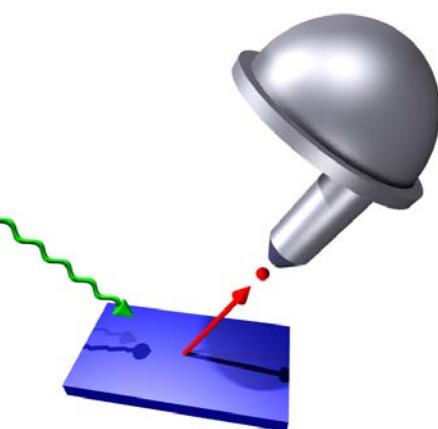
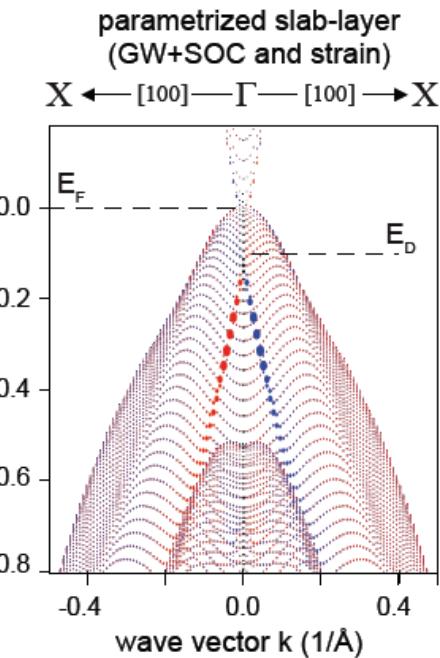
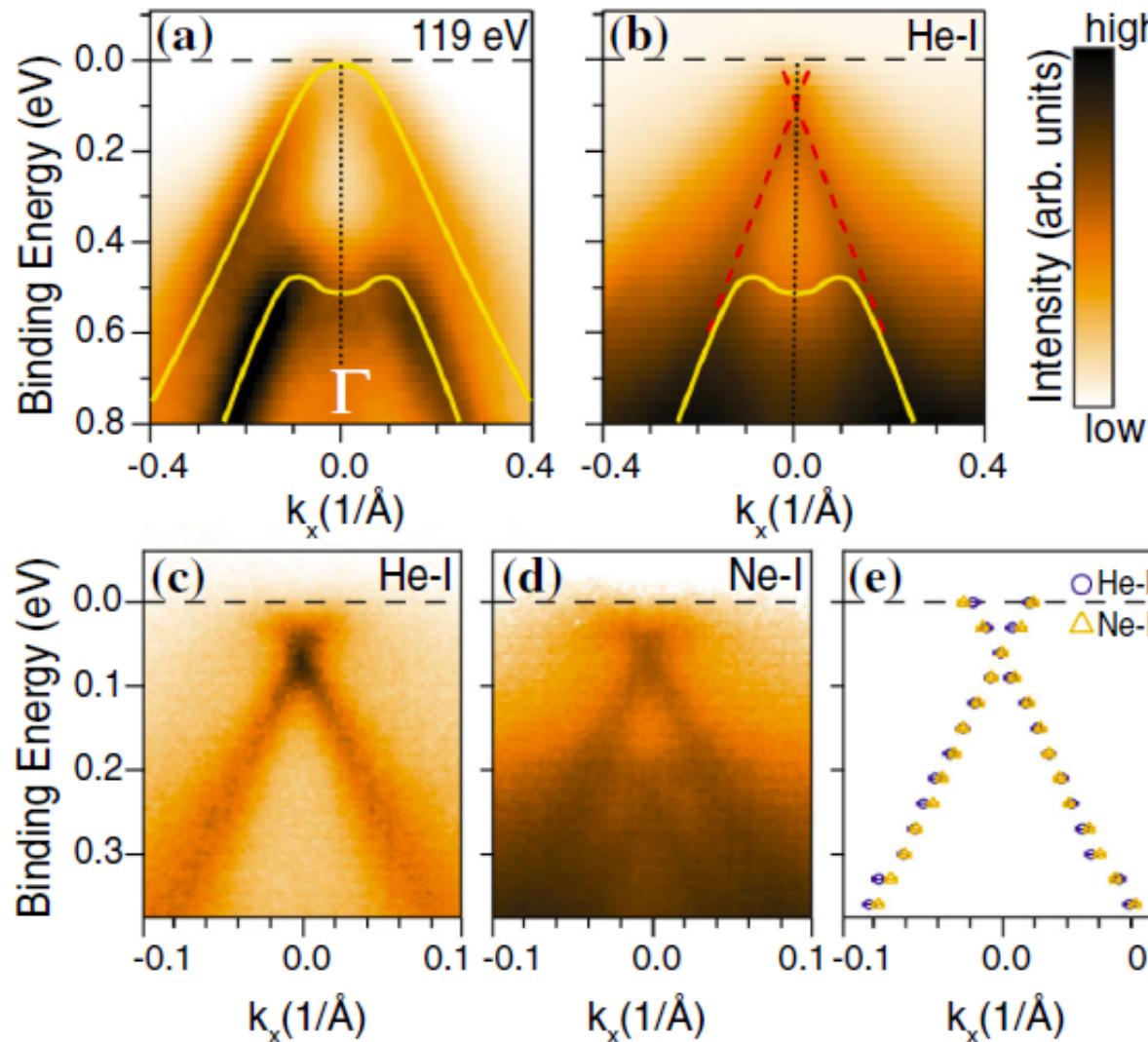


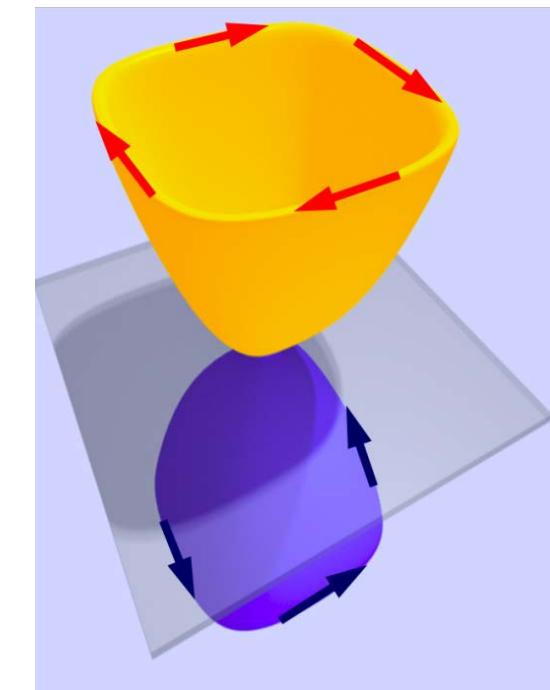
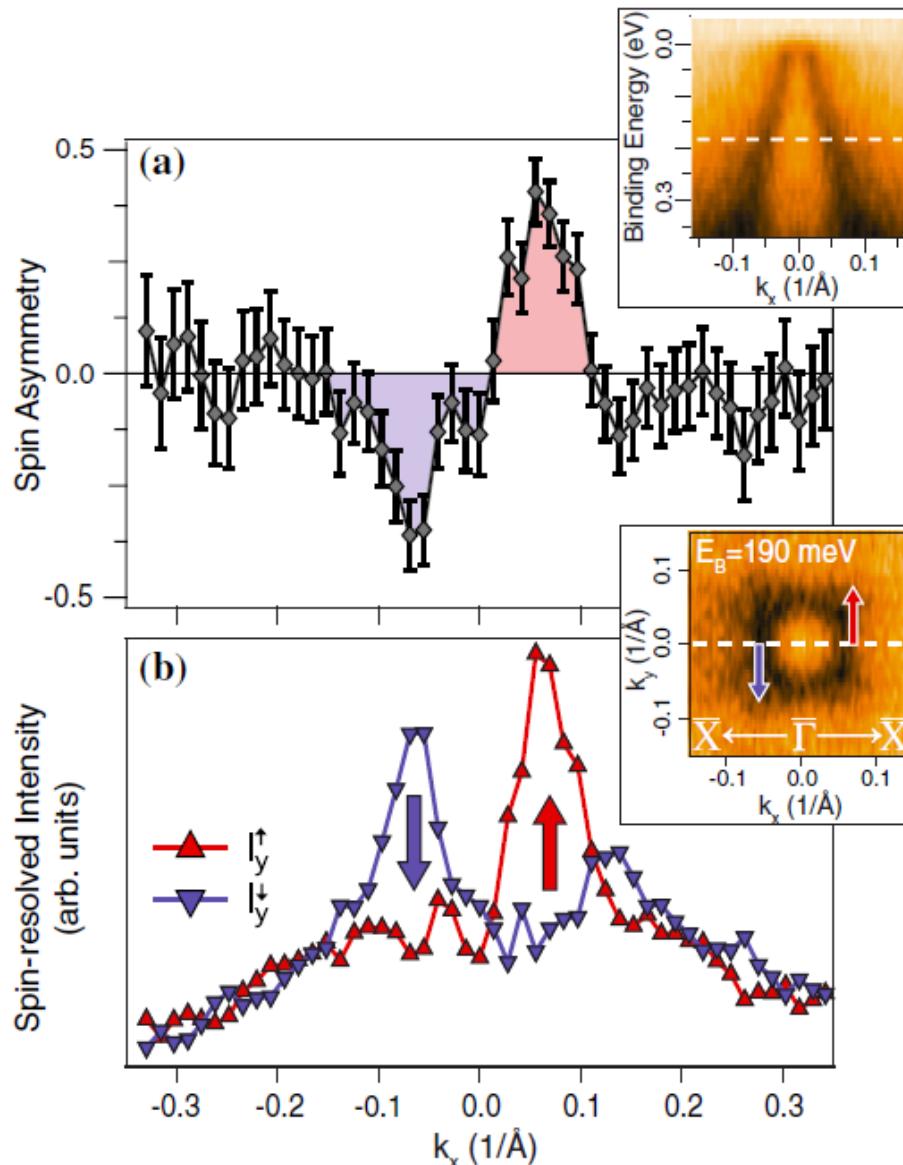
13	III A 3A	14	IV A 4A	15	V A 5A	16	VI A 6A	17	VII A 7A
5	Boron 10.811	6	Carbon 12.011	7	Nitrogen 14.007	8	Oxygen 15.999	9	Fluorine 18.998
11	IB 1B	12	IIB 2B	13	Aluminum 26.982	14	Silicon 28.086	15	Phosphorus 30.974
29	Copper 63.546	30	Zinc 65.39	31	Gallium 69.732	32	Germanium 72.61	33	Arsenic 74.922
47	Silver 107.868	48	Cadmium 112.411	49	Indium 114.818	50	Tin 118.71	51	Antimony 121.760
79	Gold 196.967	80	Mercury 200.59	81	Thallium 204.383	82	Pb Lead 207.2	83	Bismuth 208.980
								84	Po Polonium [208.982]
								85	At Astatine 209.987

- strained $\alpha\text{-Sn} - \text{diamond lattice}$
can be stabilized on suitable substrate
- strained HgTe - *zincblende structure*
→ 3dim topological insulator
(Brüne et al., PRL 2011)

Strained α -Sn/InSb(001): electronic band structure**experimental realization:**

- epitaxial growth of α -Sn on InSb(001)
- compressive strain: -0.14%
- thickness: ~ 10 nm

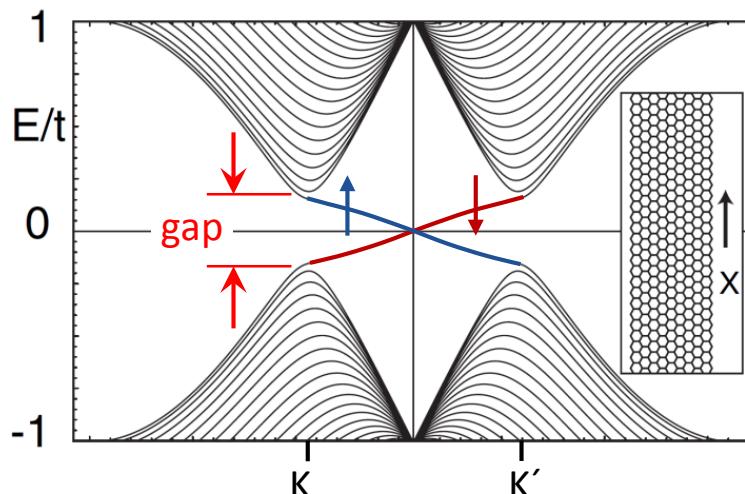
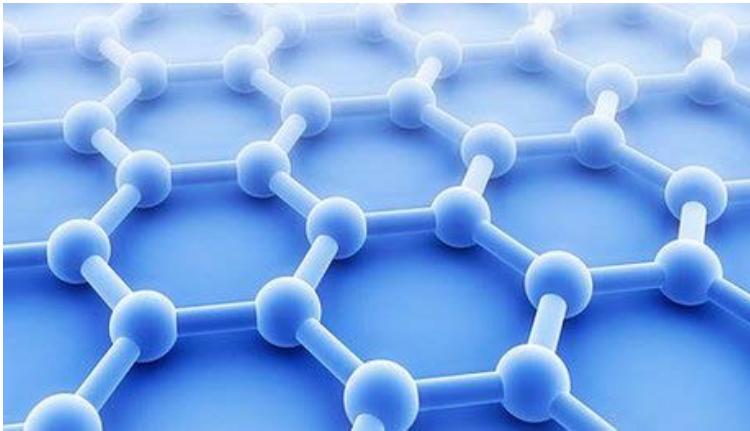




→ spin-momentum locking !

Quantum spin Hall effect in honeycomb lattices

Kane & Mele, PRL (2005): QSHE in graphene



SOC-induced
energy gap:

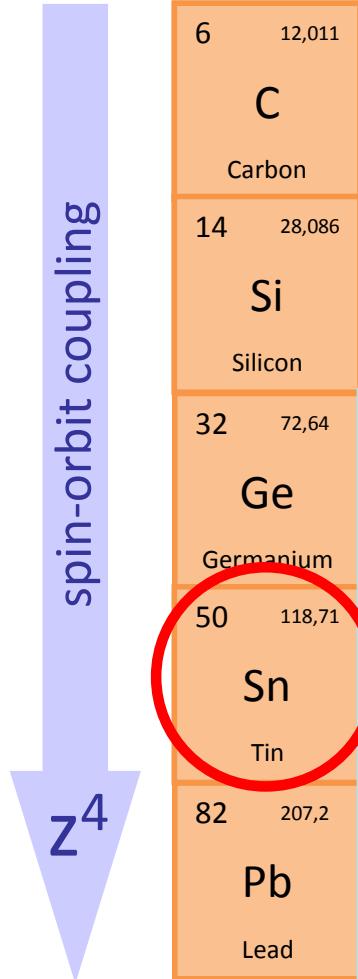
24 μ eV
graphene

2 meV

24 meV

100 meV
stanene

Y. Xu et al.,
PRL 2013



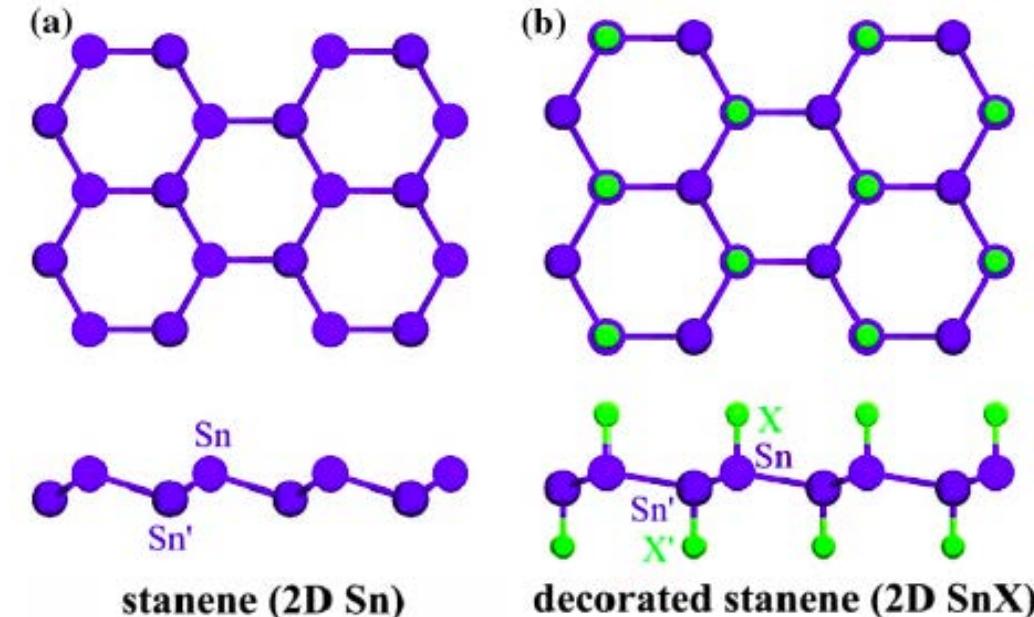
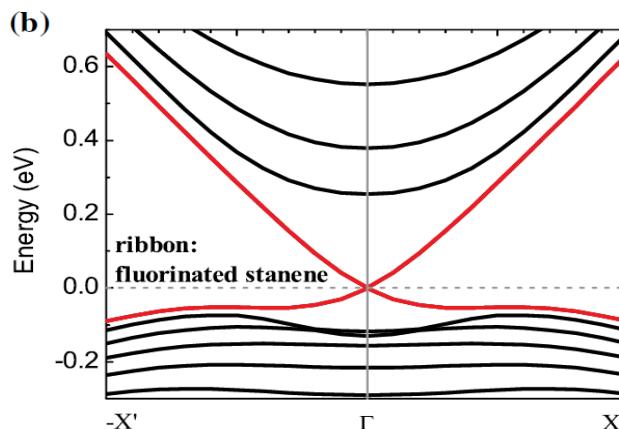
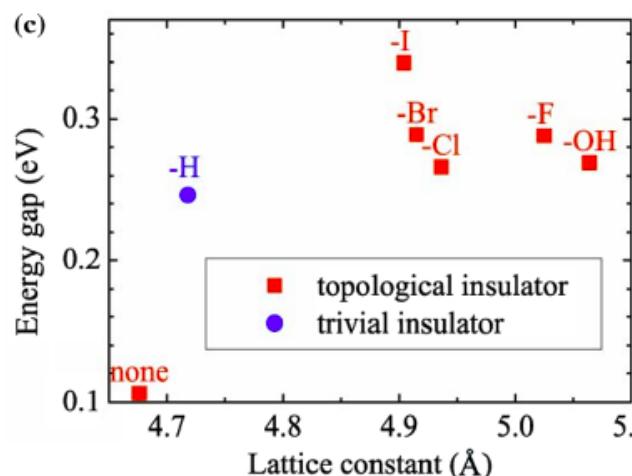
→ spin-polarized (helical) metallic edge states

Sn as 2D topological insulator: stanene

free-standing stanene

predicted to be a QSHE system/2D TI

Y. Xu *et al.*, PRL 2013



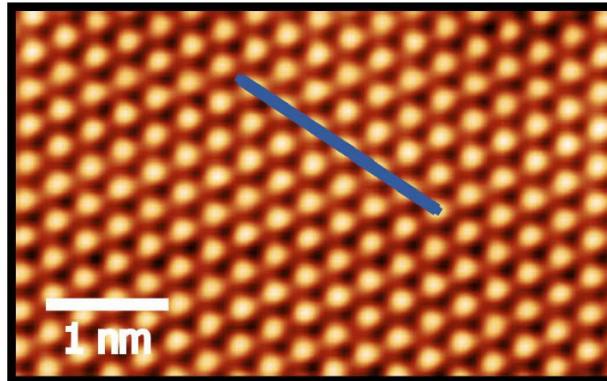
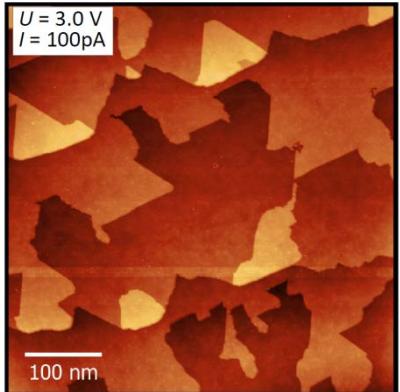
control parameters

- strain
- functionalization

real stanene requires suitable substrate
→ SiC(0001)

Epitaxial Sn sub-monolayers on SiC(0001)

substrate: 4H-SiC(0001)



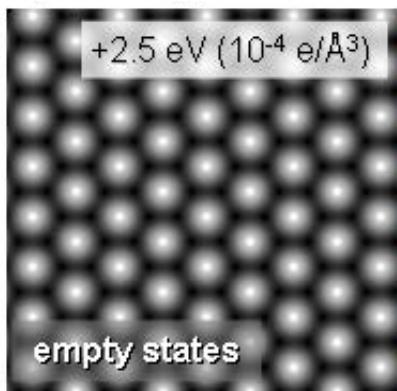
large, well-ordered terraces
prepared by hydrogen etching
(1 bar, 1200 °C)

J. Phys. Chem. C **120**, 10361 (2016)

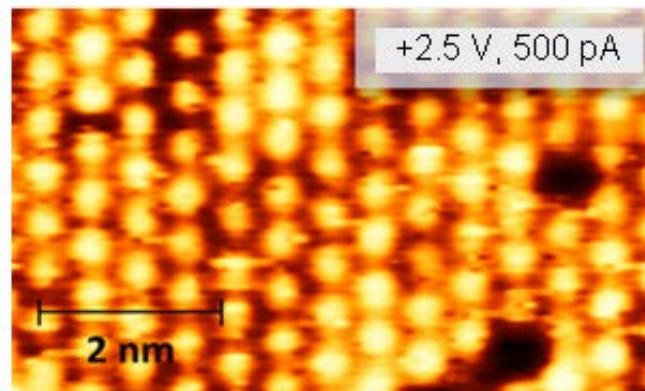
Sn epitaxy on SiC(0001)

$\sqrt{3} \times \sqrt{3}$ reconstruction (dilute phase: 1/3 ML)

theory



experiment (STM)



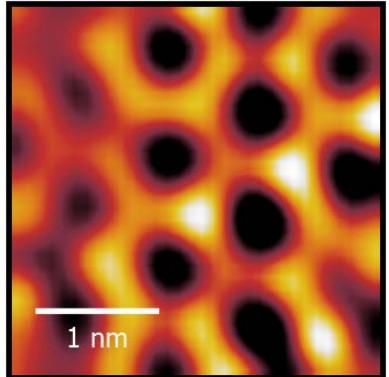
Mott insulator on a
triangular lattice

PRL **114**, 247602 (2015)
Nat. Comm. **4**, 1620 (2013)

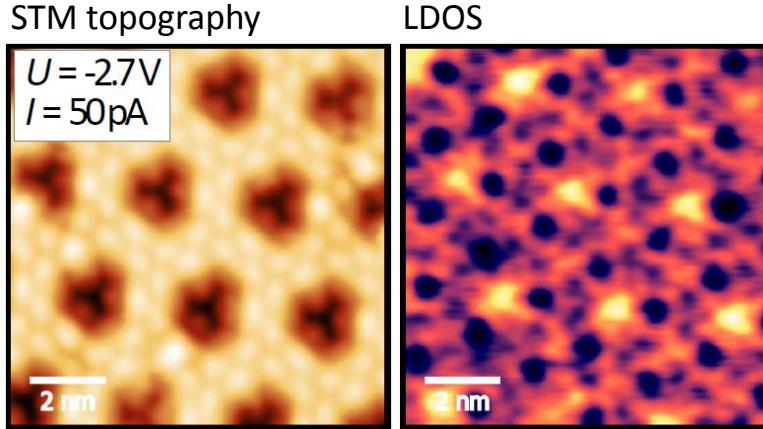
Epitaxial Sn sub-monolayers on SiC(0001)

More Sn on SiC(0001): "honeycomb" reconstructions

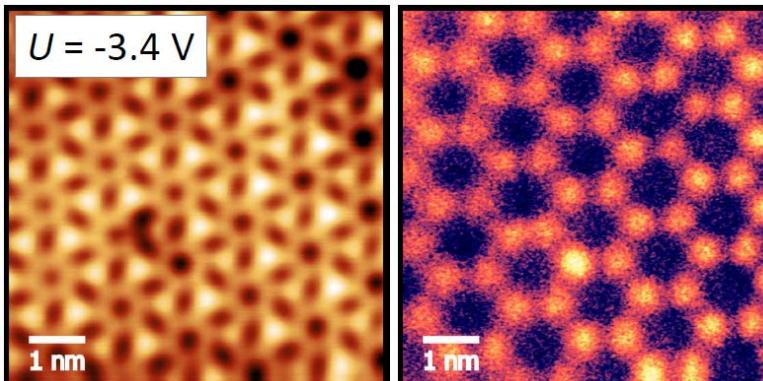
3×3



$6\sqrt{3} \times 6\sqrt{3}$



$\sqrt{21} \times \sqrt{21}$



multitude of phases for increasing Sn coverage/
decreasing substrate temp.:

"super-stanene" -
lattice constants too large
for real stanene

→ real topological stanene
yet to be realized !

Stanene on Bi_2Te_3 ?

ARTICLES

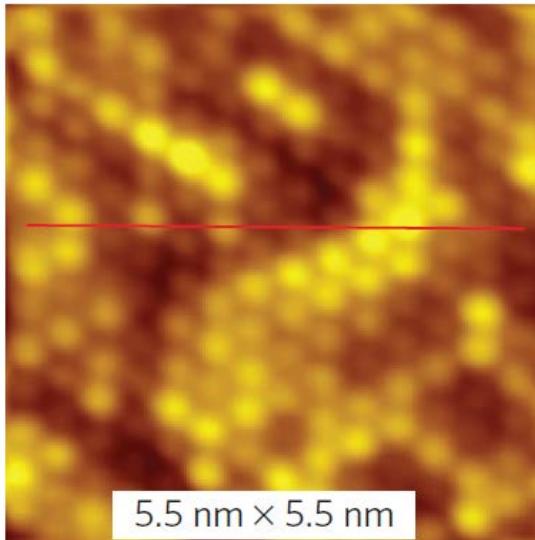
Nat. Mat. **14**, 1020 (2015)nature
materials

PUBLISHED ONLINE: 3 AUGUST 2015 | DOI: 10.1038/NMAT4384

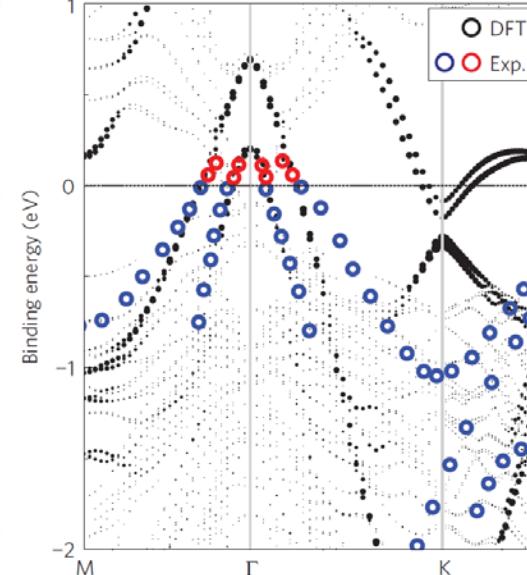
Epitaxial growth of two-dimensional stanene

Feng-feng Zhu^{1†}, Wei-jiong Chen^{1†}, Yong Xu^{2,3,4†}, Chun-lei Gao^{1,5}, Dan-dan Guan^{1,5}, Can-hua Liu^{1,5}, Dong Qian^{1,5*}, Shou-Cheng Zhang^{2,3,4} and Jin-feng Jia^{1,5*}

STM

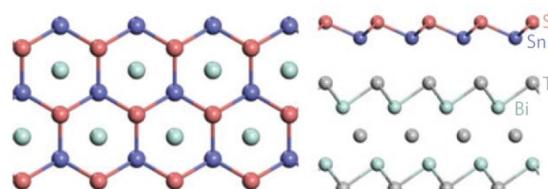


ARPES vs. DFT



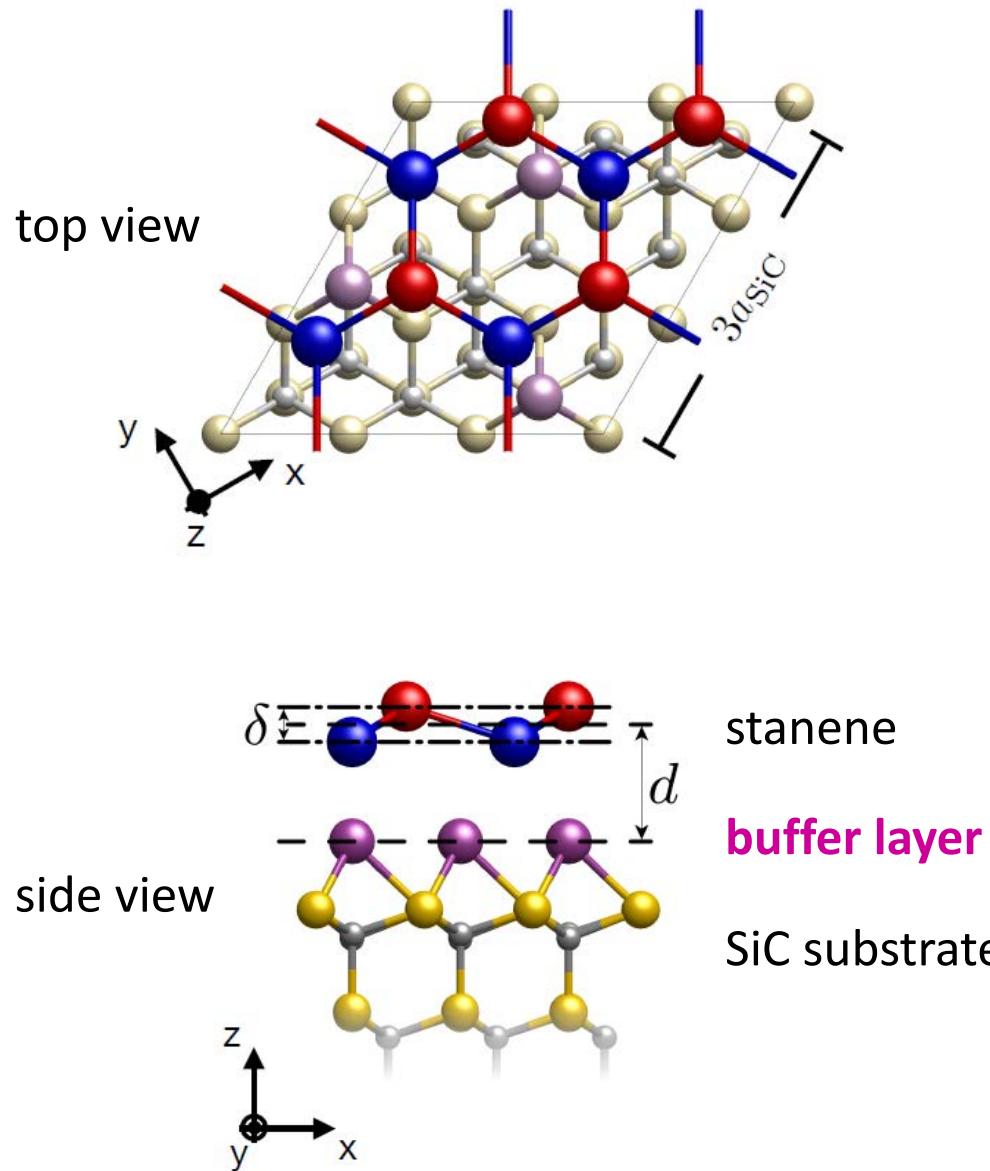
- STM shows hexagonal (triangular) lattice, **not honeycomb**
- ARPES finds **metal, not 2D TI**

→ **real topological stanene
yet to be realized !**



?

Quasi-freestanding stanene via substrate engineering: a proposal

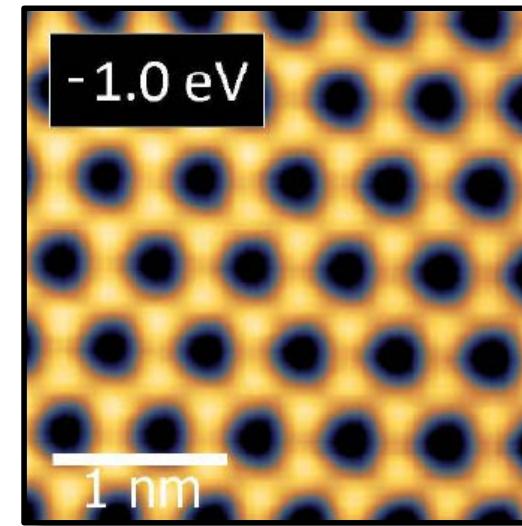


	P	As	Sb	Bi
Si				
ΔE_K (meV)	68	43	40	61
d_{Buf} (Å)	1.8	1.9	2.2	2.3
d (Å)	3.2	3.5	3.6	3.6
δ (Å)	0.46	0.45	0.46	0.46
\mathbb{Z}_2	0	0	0	0
C				
ΔE_K (meV)	69	55	42	11
d_{Buf} (Å)	1.3	1.4	1.6	1.8
d (Å)	3.9	3.8	3.9	3.7
δ (Å)	0.44	0.44	0.44	0.45
\mathbb{Z}_2	1	1	1	0

DFT calculations by D. Di Sante, P. Eck,
R. Thomale, and G. Sangiovanni

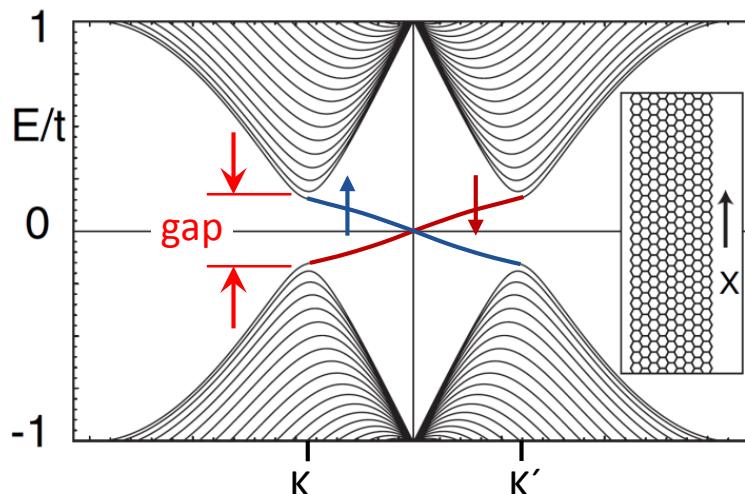
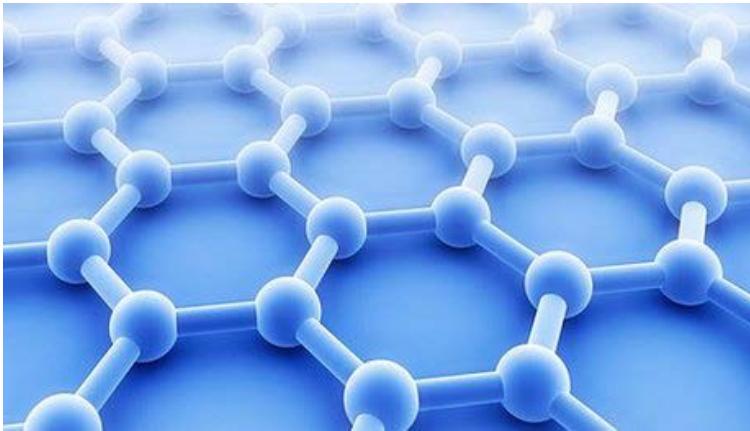


A new quantum spin Hall paradigm: bismuthene/SiC(0001)



Quantum spin Hall effect in honeycomb lattices

Kane & Mele, PRL (2005): QSHE in graphene



SOC-induced
energy gap:

24 μ eV
graphene

2 meV

24 meV

100 meV
stanene
Y. Xu et al.,
PRL 2013

spin-orbit coupling
 Z^4

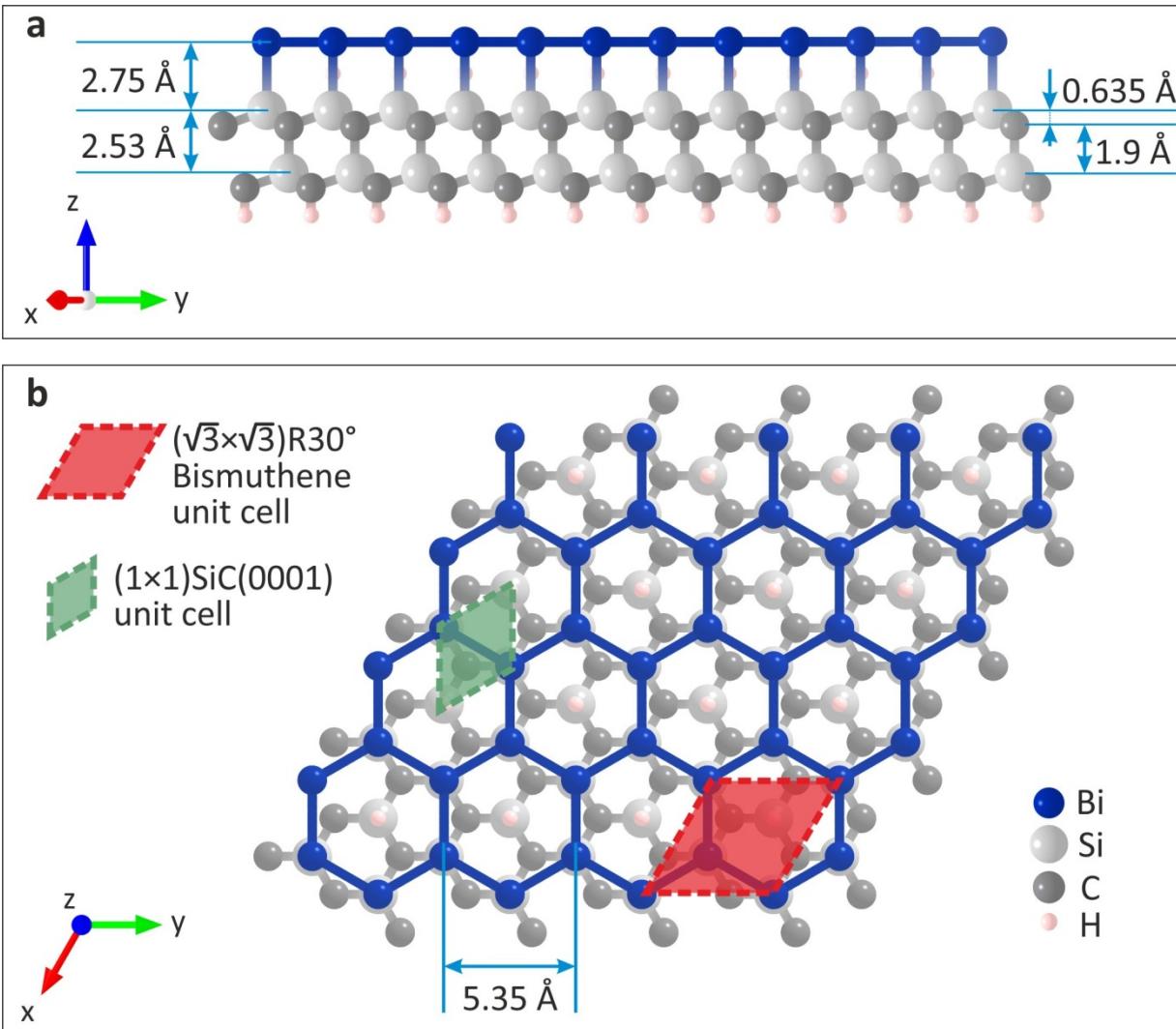
IV		V	
6	12,011	7	14,007
C	Carbon	N	Nitrogen
14	28,086	15	30,974
Si	Silicon	P	Phosphorous
32	72,64	33	74,922
Ge	Germanium	As	Arsenic
50	118,71	51	121,76
Sn	Tin	Sb	Antimony
82	207,2	83	208,98
Pb	Lead	Bi	Bismuth

bismuthene:
Bi/SiC(0001)

Hsu et al.,
NJP 2015

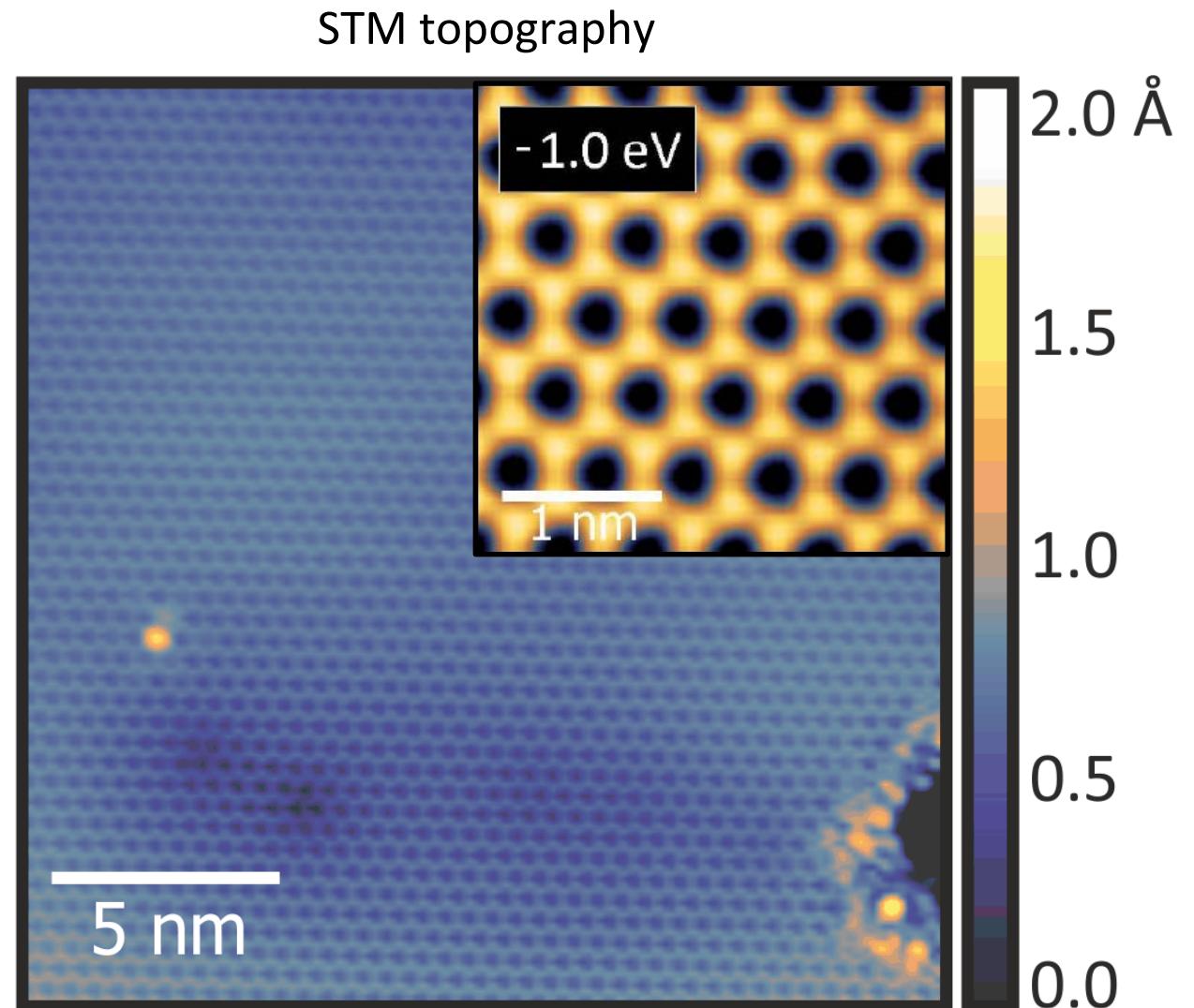
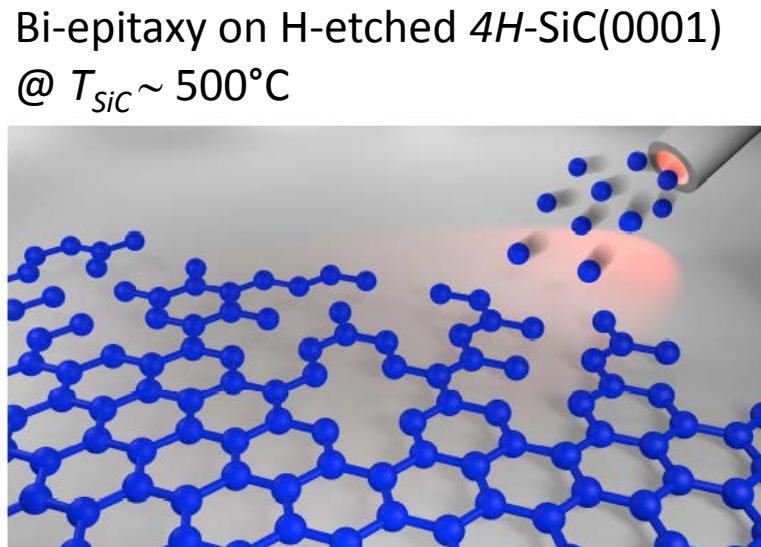
→ spin-polarized (helical) metallic edge states

Bismuthene/SiC(0001): structure

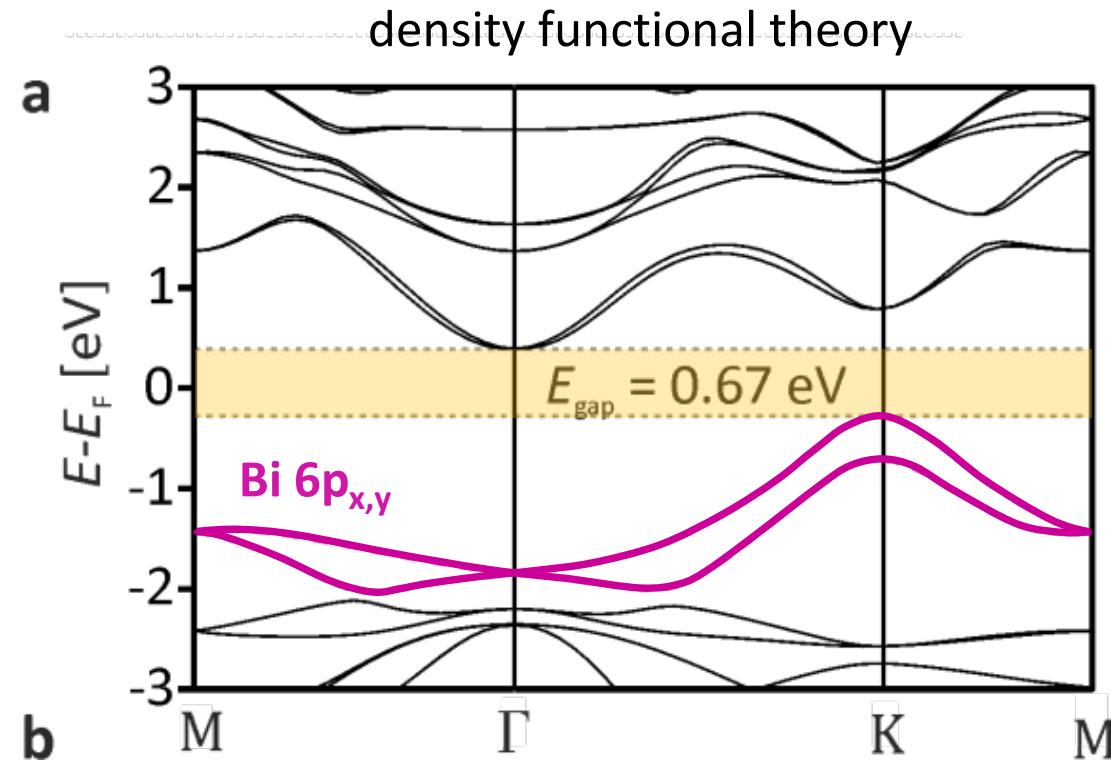


18% tensile strain
→ planar honeycombs

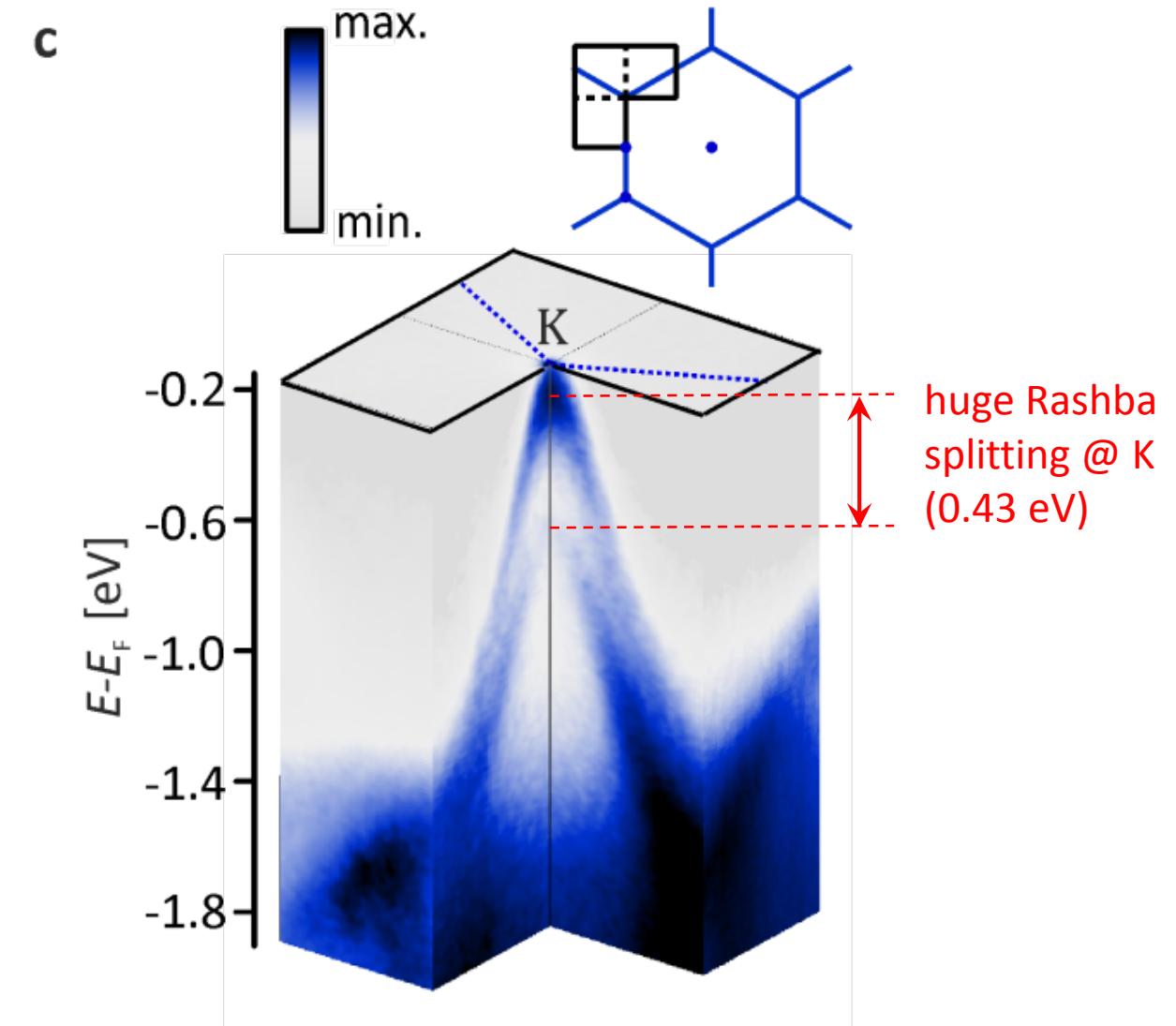
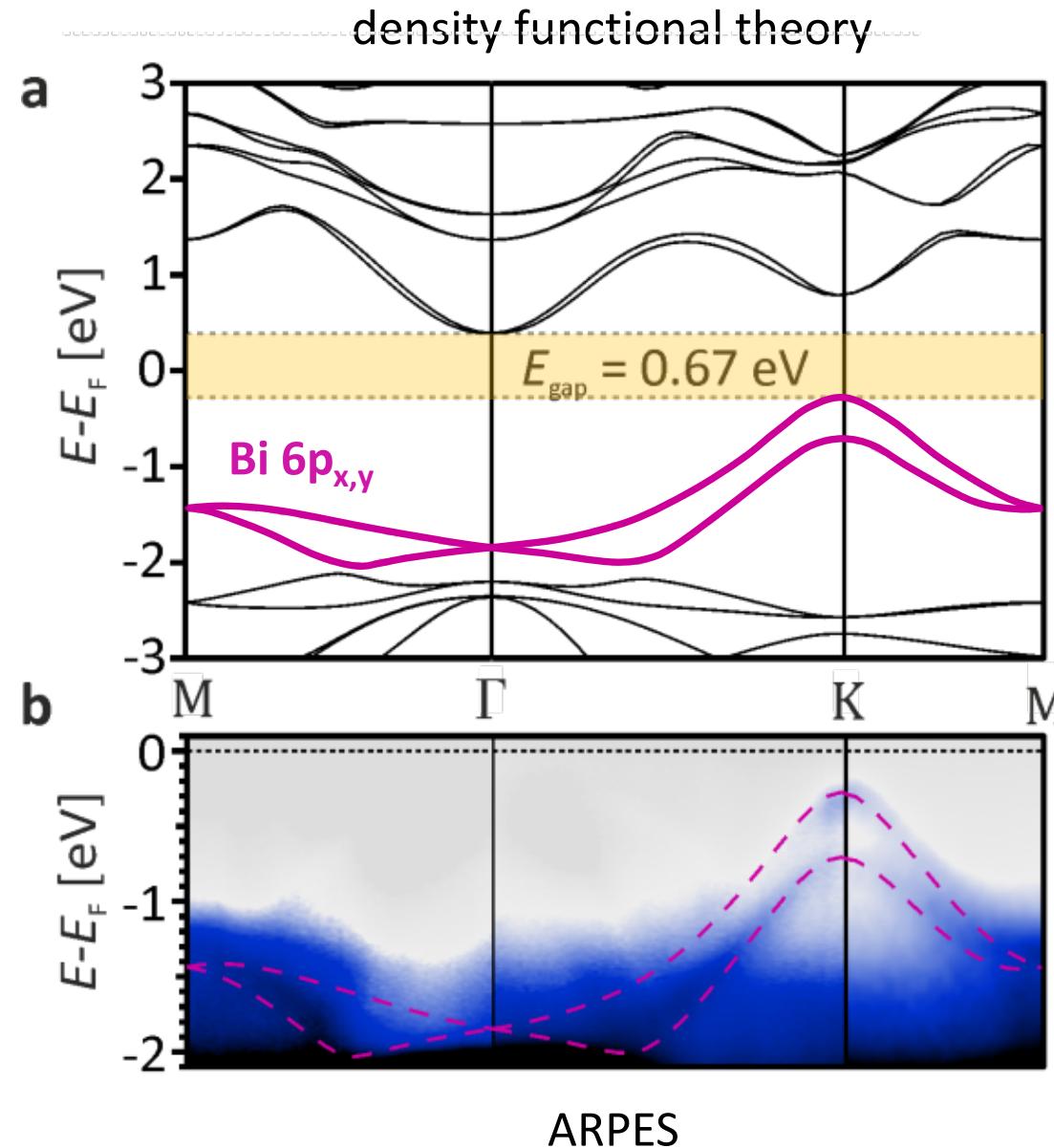
($\sqrt{3} \times \sqrt{3}$)R30°
reconstruction



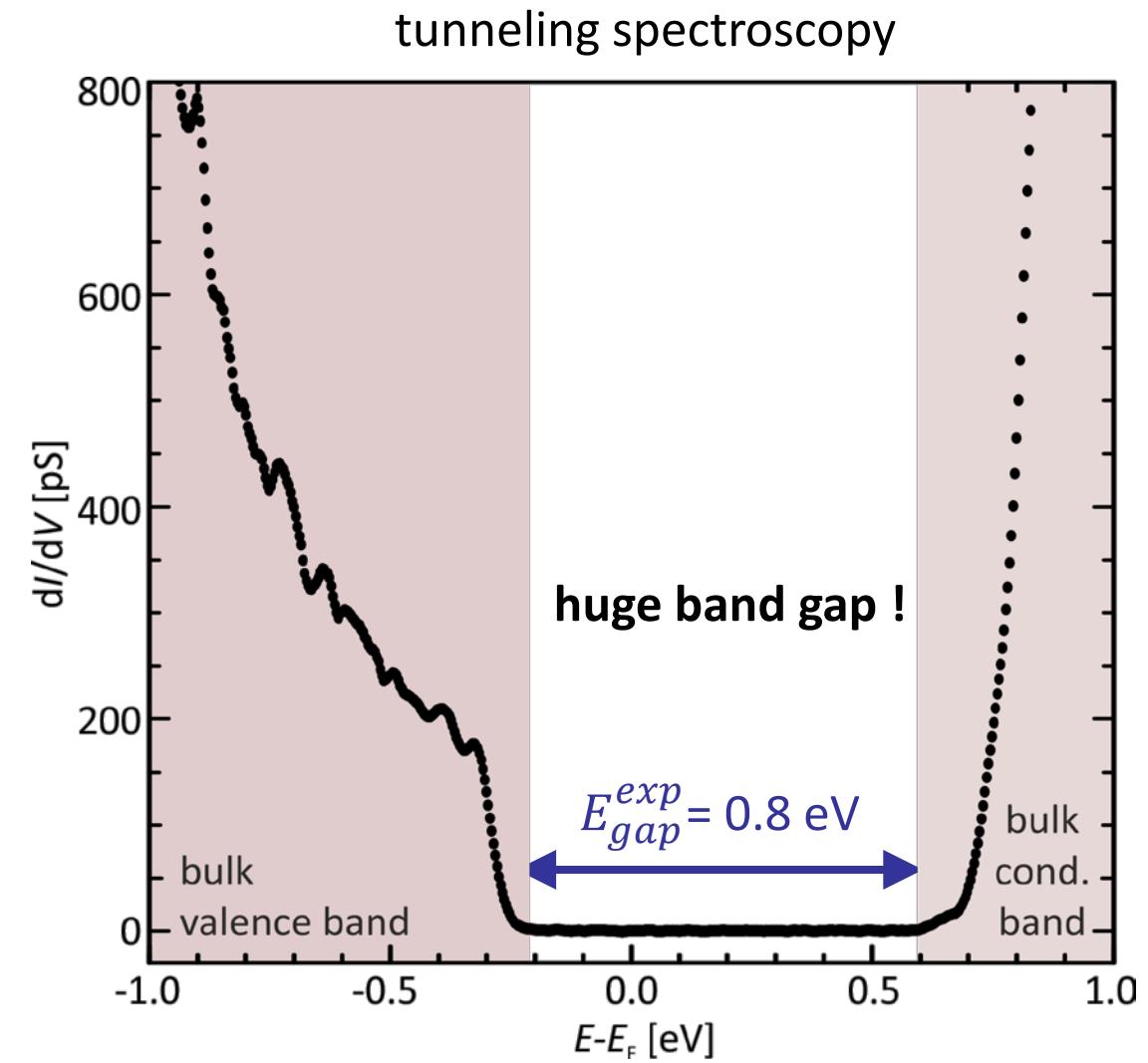
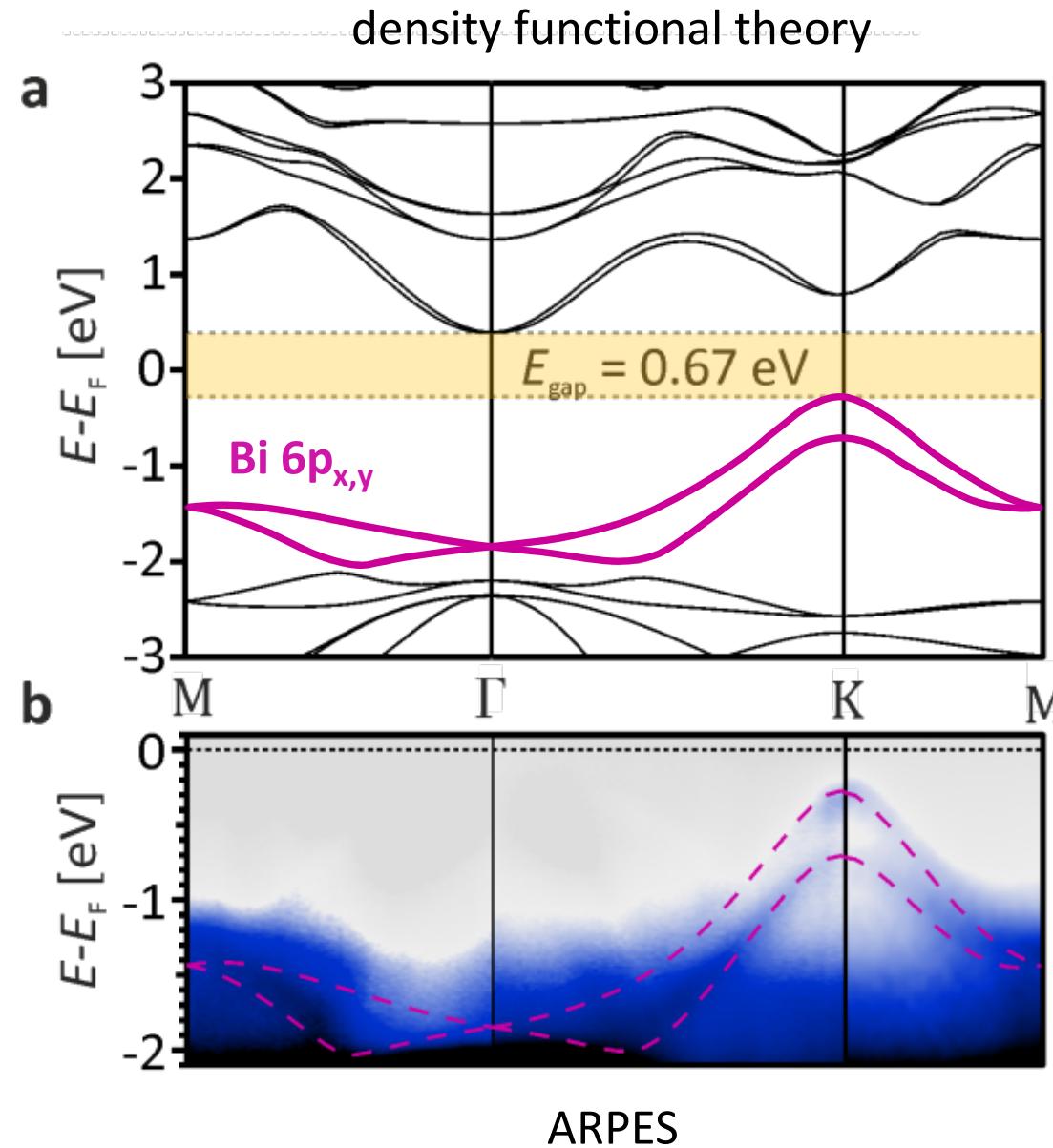
Bismuthene/SiC(0001): band structure and energy gap



Bismuthene/SiC(0001): band structure and energy gap

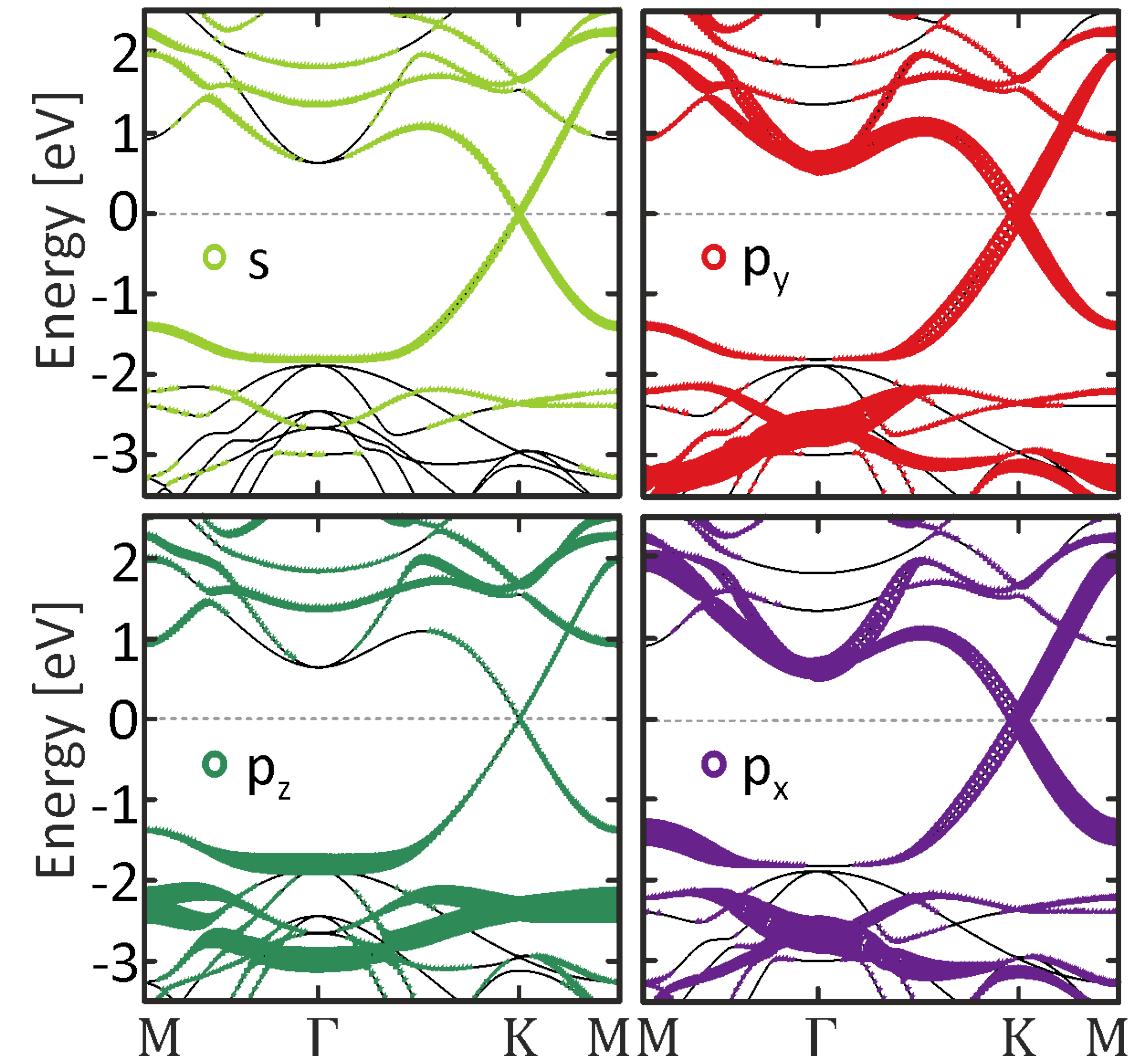


Bismuthene/SiC(0001): band structure and energy gap



band structure analysis w/o
spin-orbit coupling:

- sp^2 -type σ -bonding
 - p_x and p_y orbitals dominate at low energies
 - **important role of substrate:**
removes p_z from Fermi level
→ "orbital filtering"
- exploitation of *local* ($L_z S_z$) SOC
→ distinctly different from graphene
& Kane-Mele model !



basis states: $|p_{x\uparrow}^A\rangle, |p_{y\uparrow}^A\rangle, |p_{x\uparrow}^B\rangle, |p_{y\uparrow}^B\rangle; \quad |p_{x\downarrow}^A\rangle, |p_{y\downarrow}^A\rangle, |p_{x\downarrow}^B\rangle, |p_{y\downarrow}^B\rangle$

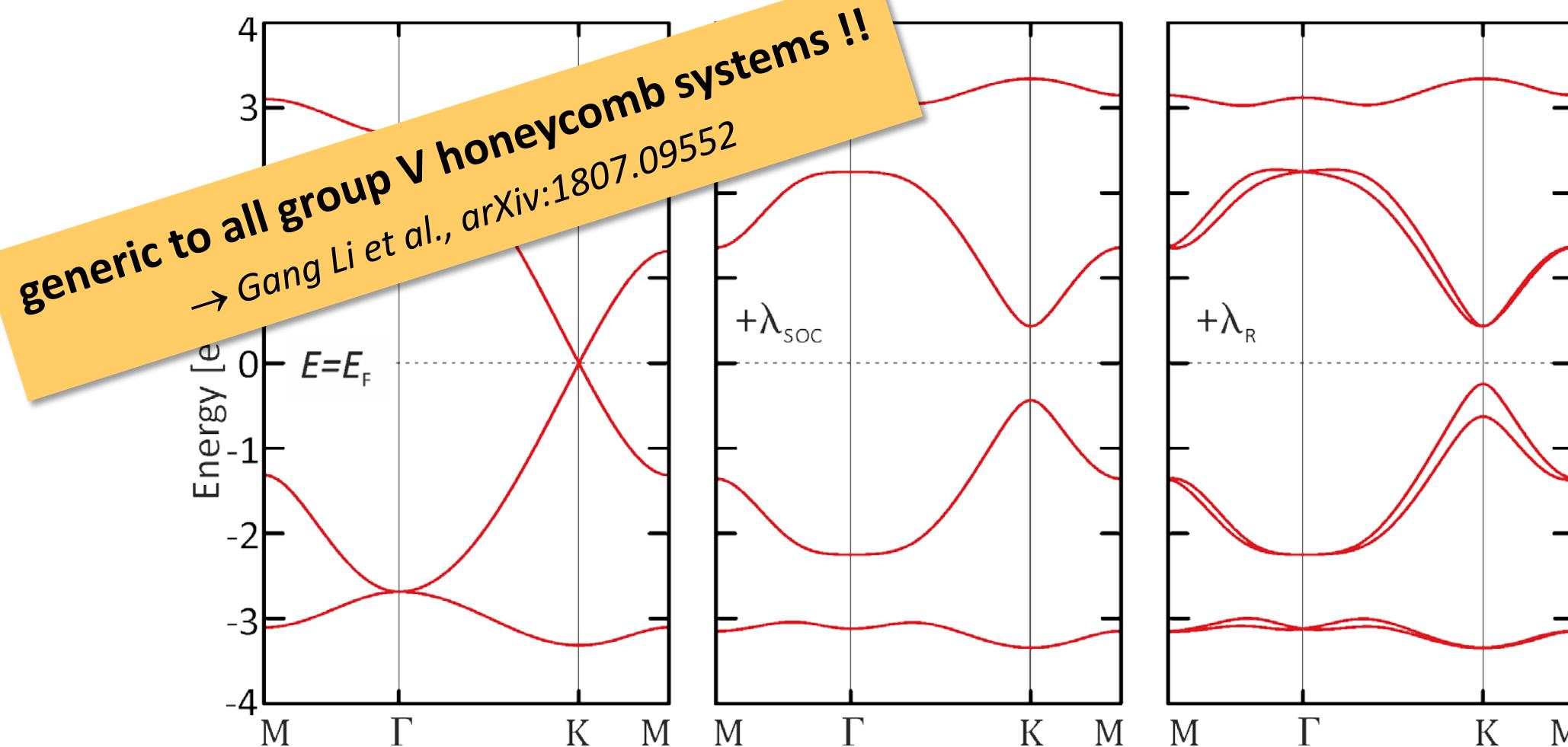
tight-binding Hamiltonian: $H_{eff}^{\sigma\sigma} = \begin{pmatrix} H_{\uparrow\uparrow}^{\sigma\sigma} & H_{\uparrow\downarrow}^{\sigma\sigma} \\ H_{\downarrow\uparrow}^{\sigma\sigma} & H_{\downarrow\downarrow}^{\sigma\sigma} \end{pmatrix}$ with

$$H_{\uparrow\uparrow/\downarrow\downarrow}^{\sigma\sigma} = H_{0,\uparrow\uparrow/\downarrow\downarrow}^{\sigma\sigma} \pm \lambda_{SOC} \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{pmatrix}$$

**off-diagonal orbital term
mixes in local (atomic) SOC**

$$H_{\uparrow\downarrow}^{\sigma\sigma} = (H_{\downarrow\uparrow}^{\sigma\sigma})^+ = \lambda_R \begin{pmatrix} 0 & 0 & m_1 & m_2 \\ 0 & 0 & m_2 & m_3 \\ m_4 & m_5 & 0 & 0 \\ m_5 & m_6 & 0 & 0 \end{pmatrix}$$

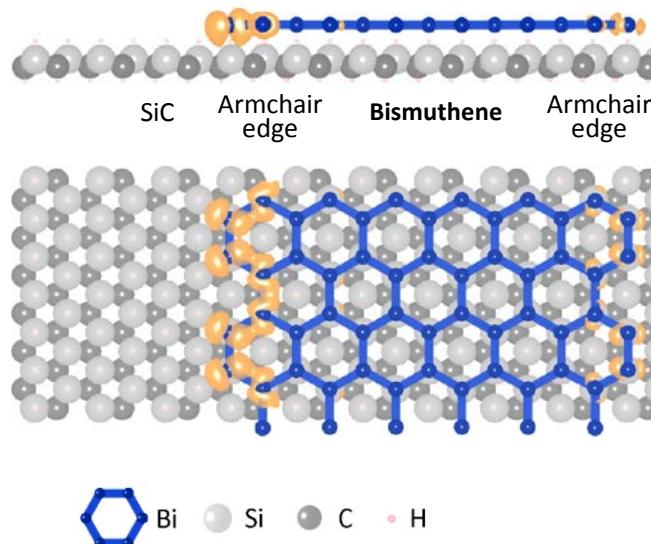
Rashba

Bismuthene/SiC(0001): effective model *including SOC*

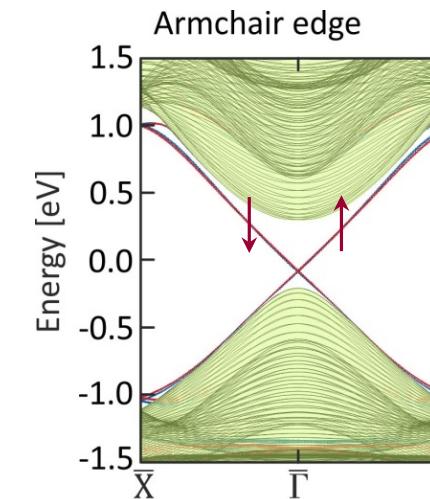
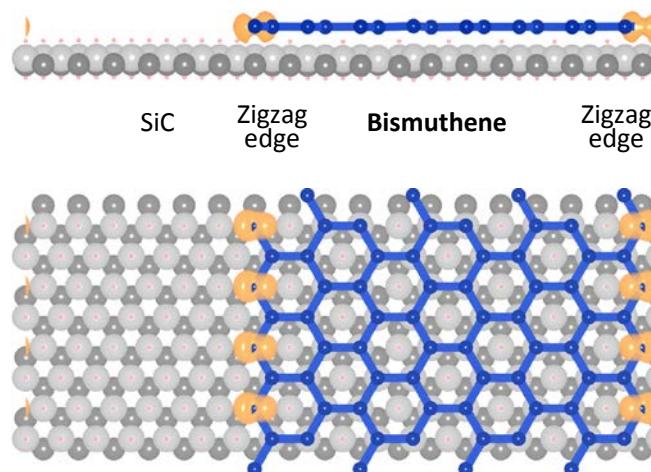
$$H_{eff}^{\sigma\sigma} = H_0^{\sigma\sigma} + \lambda_{SOC} H_{SOC}^{\sigma\sigma} + \lambda_R H_R^{\sigma\sigma}$$

Bismuthene/SiC(0001): topological edge states

Armchair edge ribbon



Zigzag edge ribbon



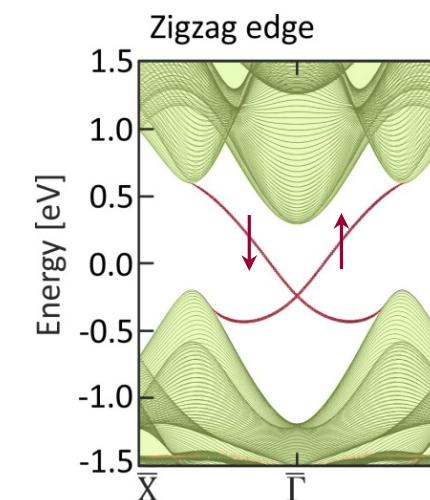
TB calculations for nanoribbon geometries:

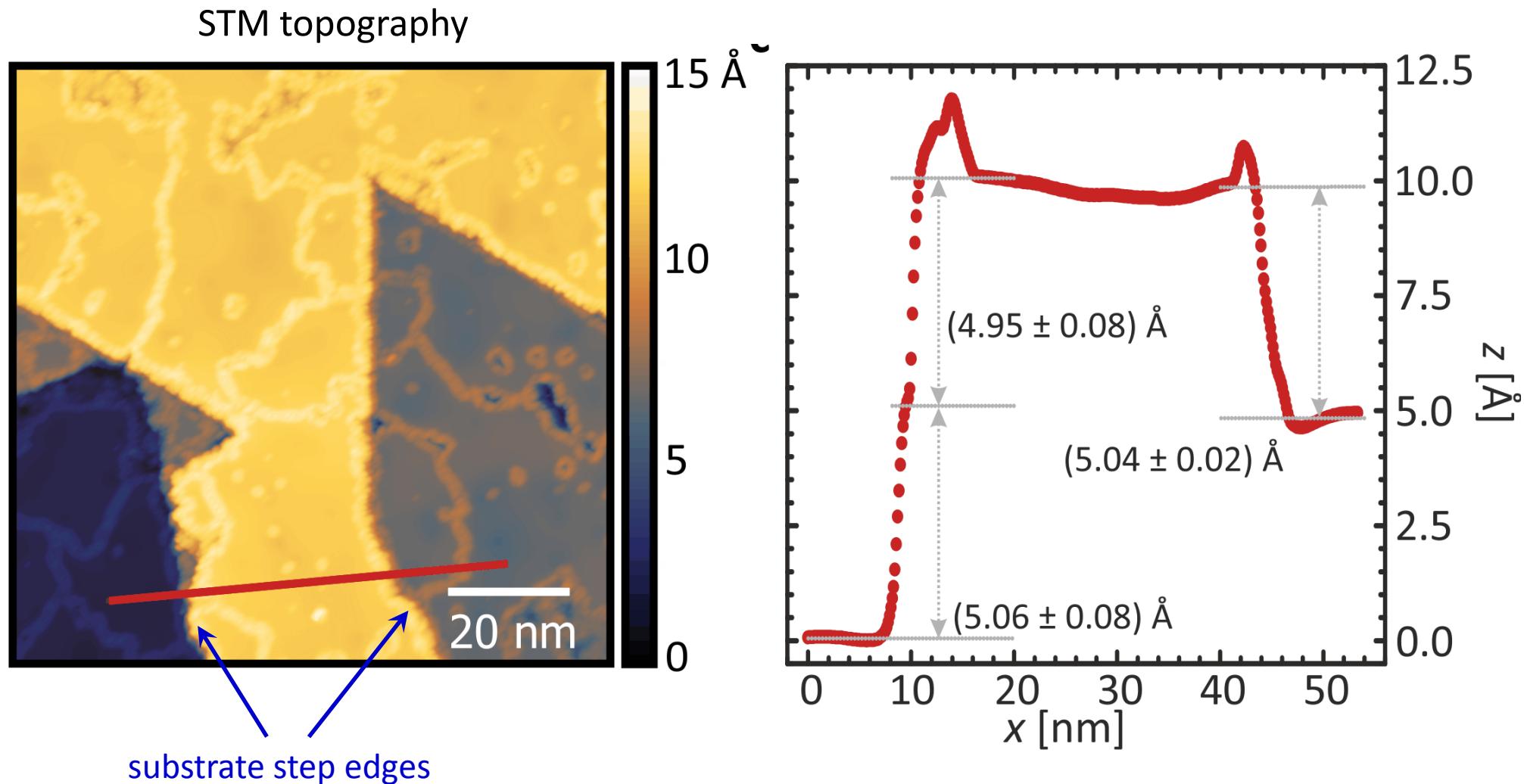
→ helical metallic edge states

→ topological invariant $Z_2 = 1$

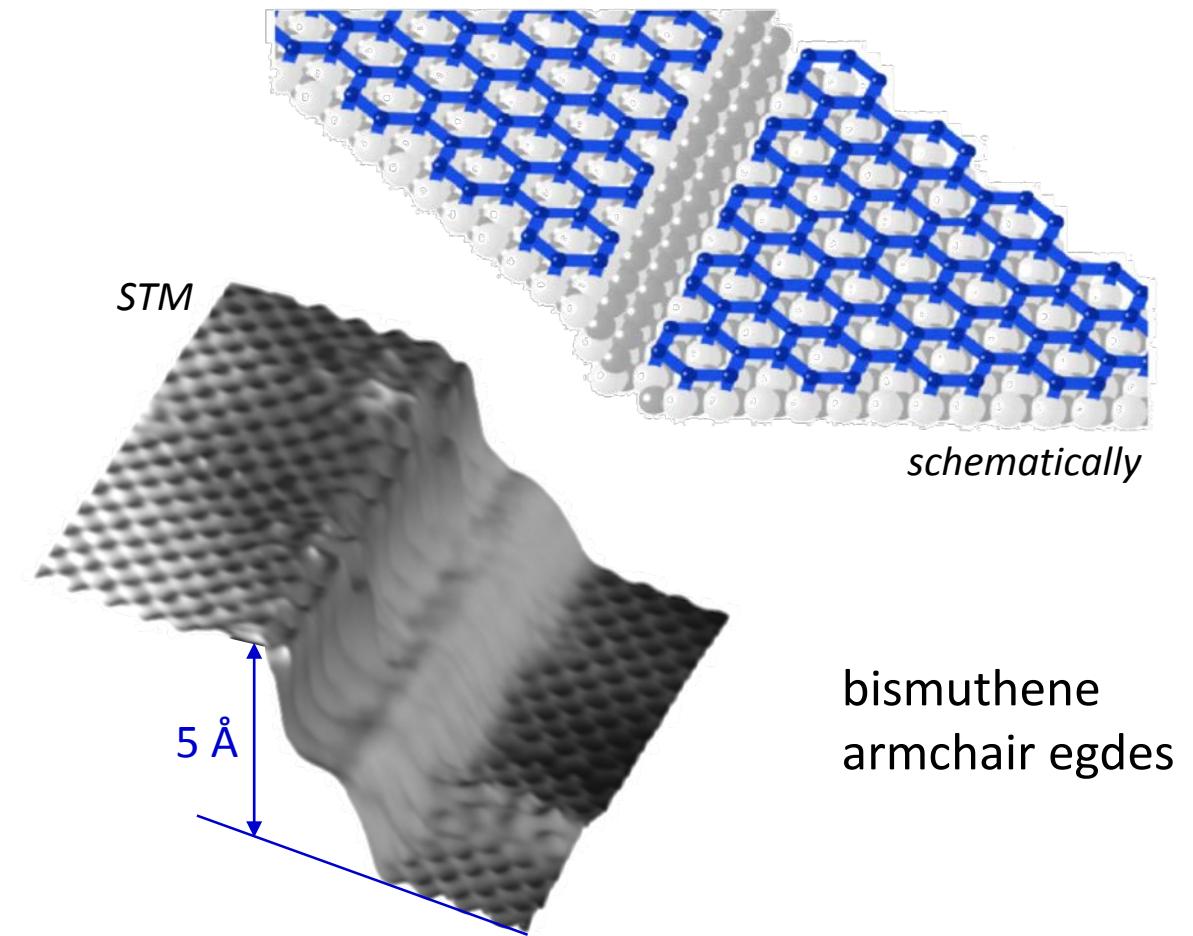
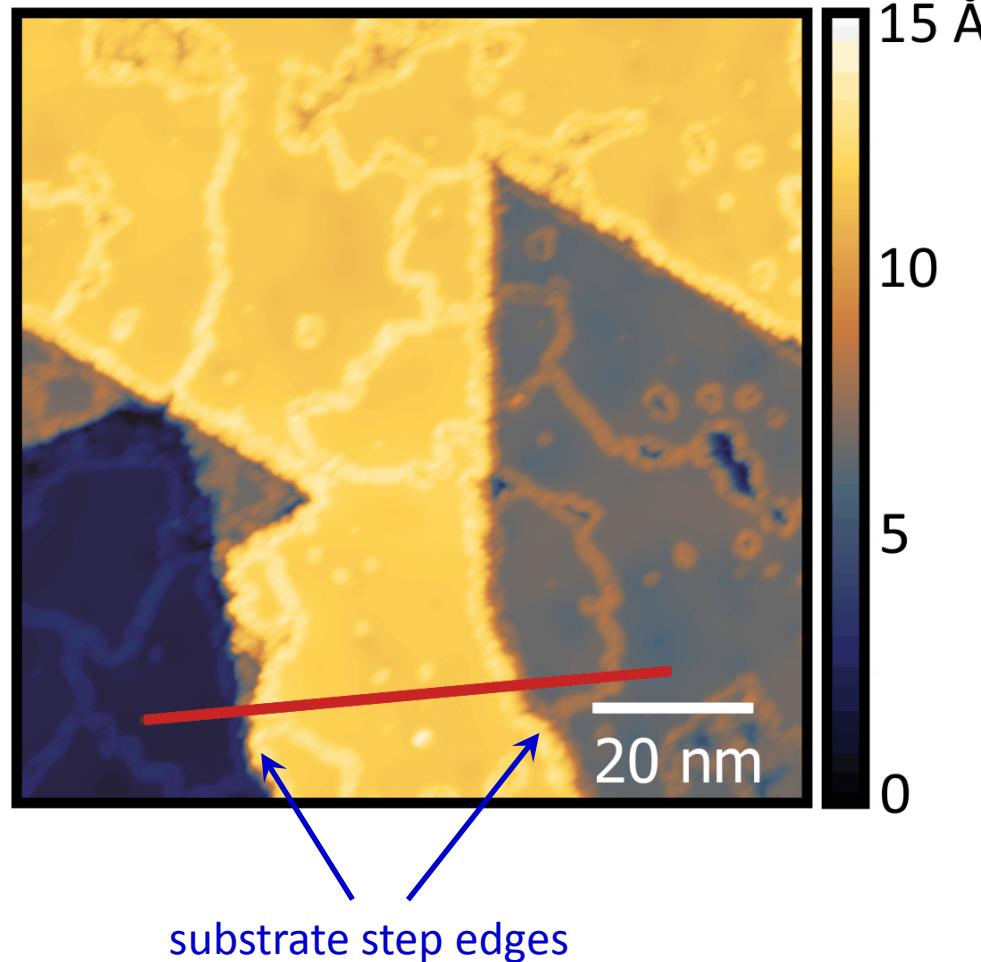
→ **bismuthene/SiC is a QSH system !**

(cf. Hsu et al., NJP 2015)

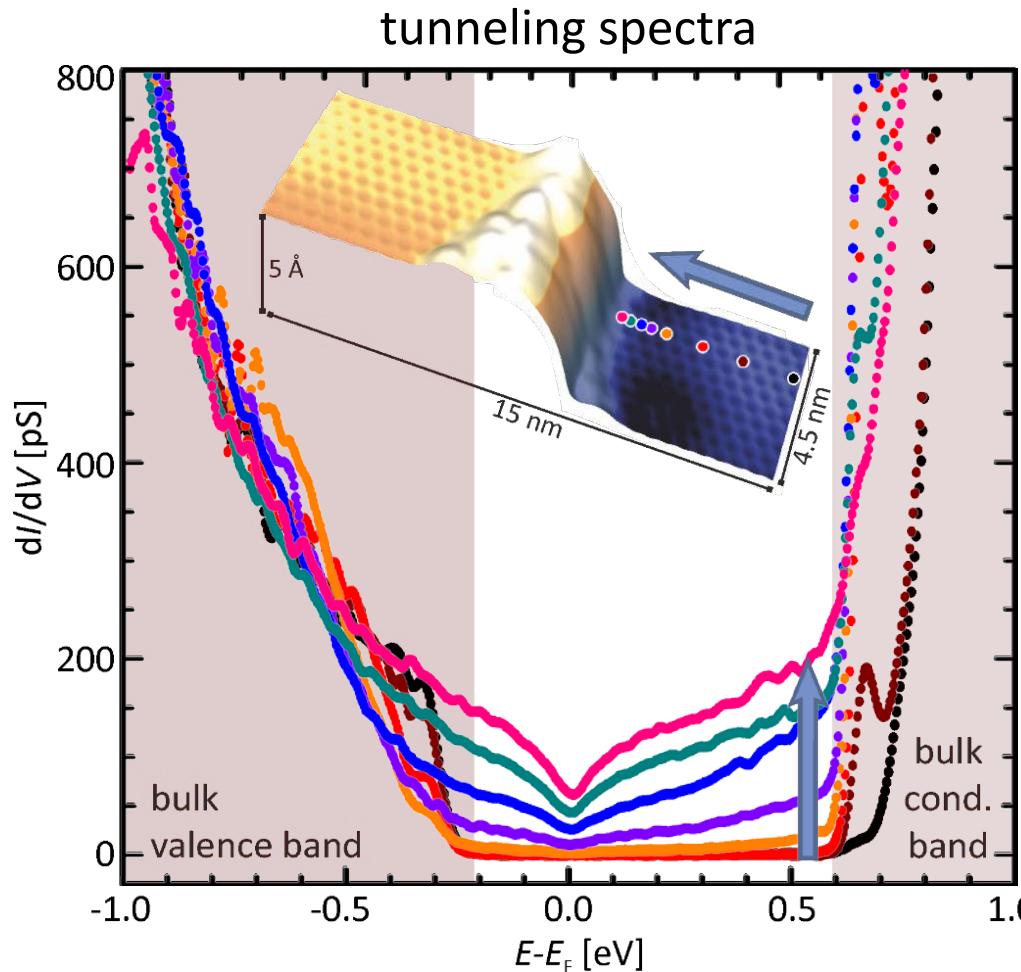




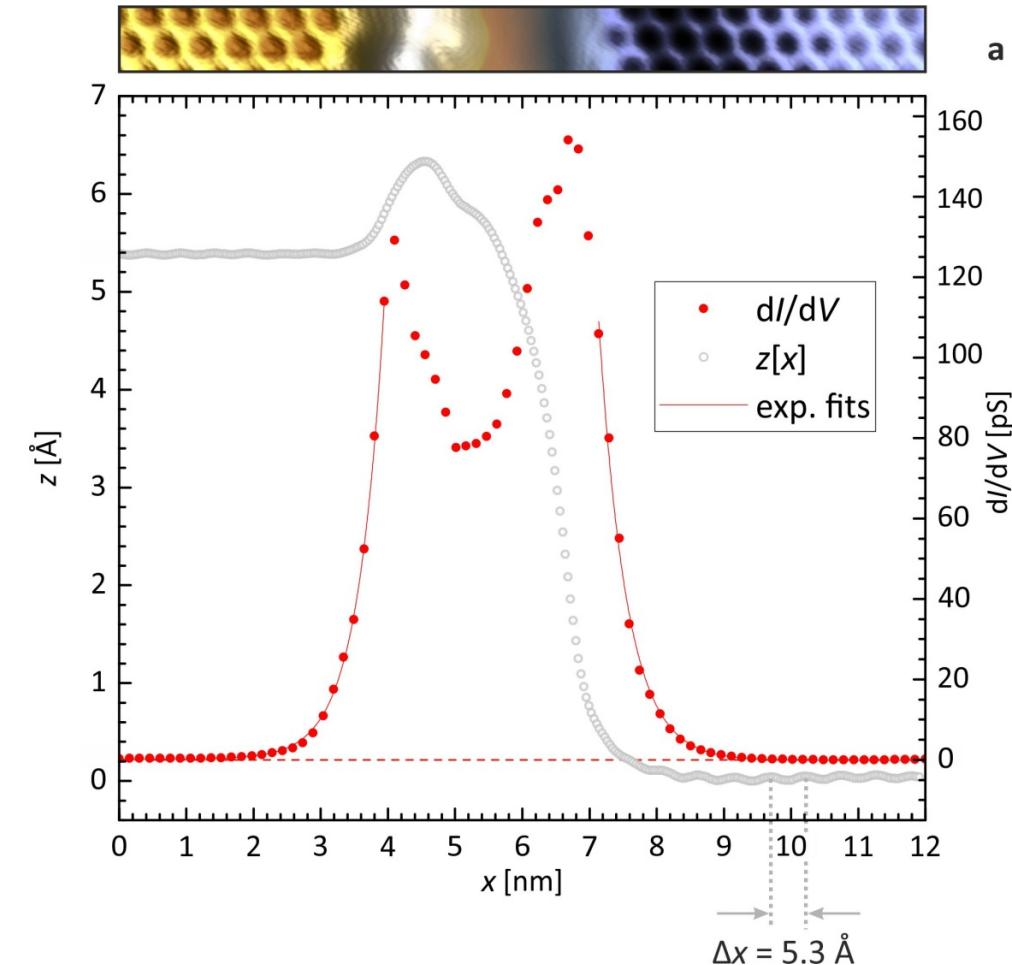
STM topography



Bismuthene/SiC(0001): edge state spectroscopy



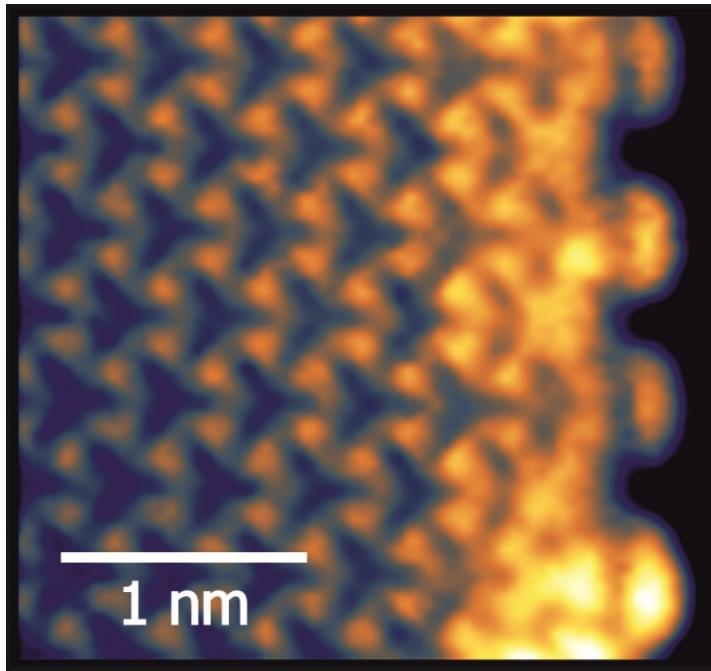
- metallic edge states in bulk gap



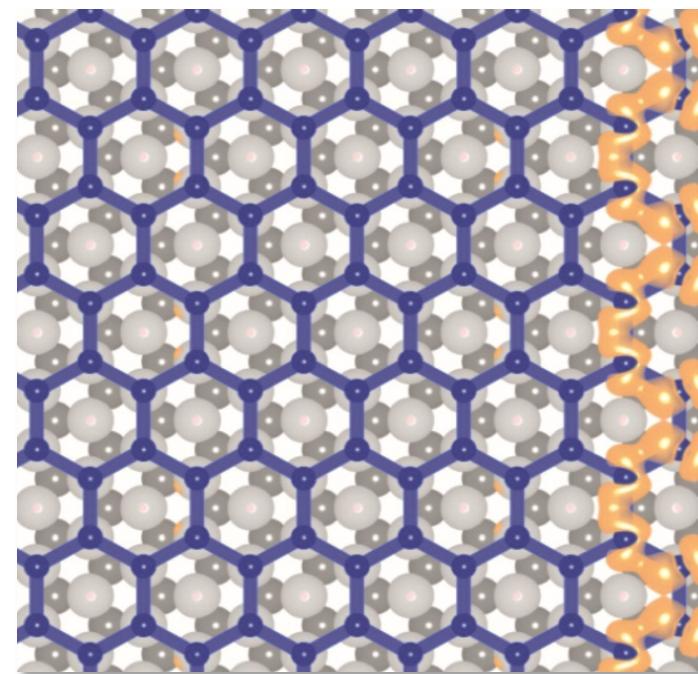
- exponential decay into bulk quantitatively consistent with theory

charge density @ arm chair edge

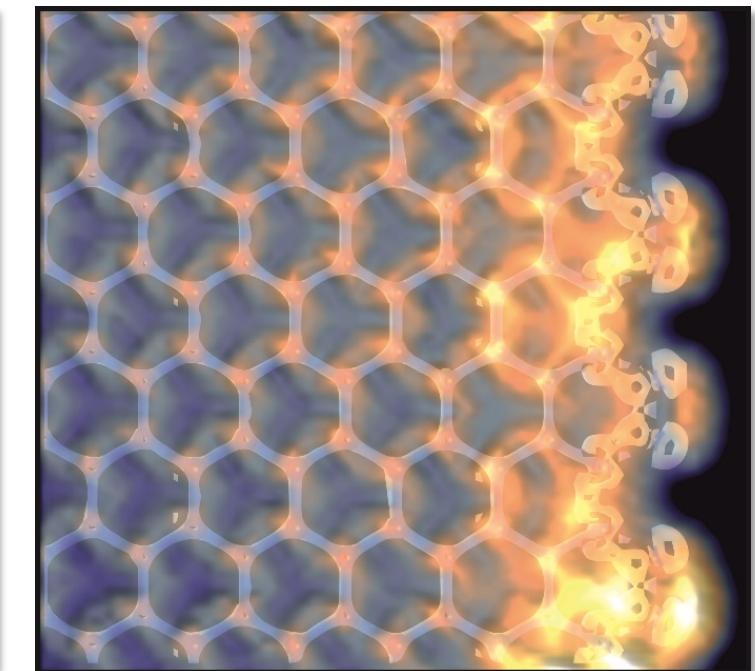
STM



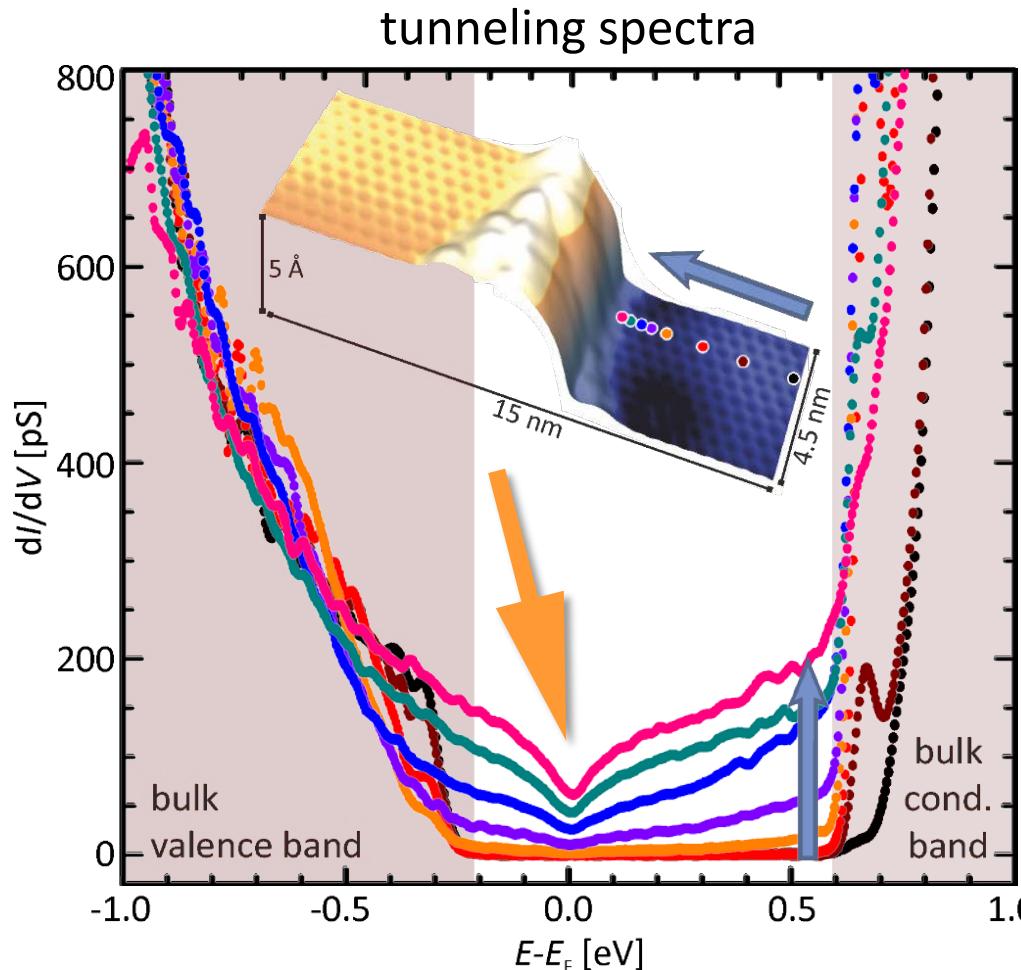
nano-ribbon calculation



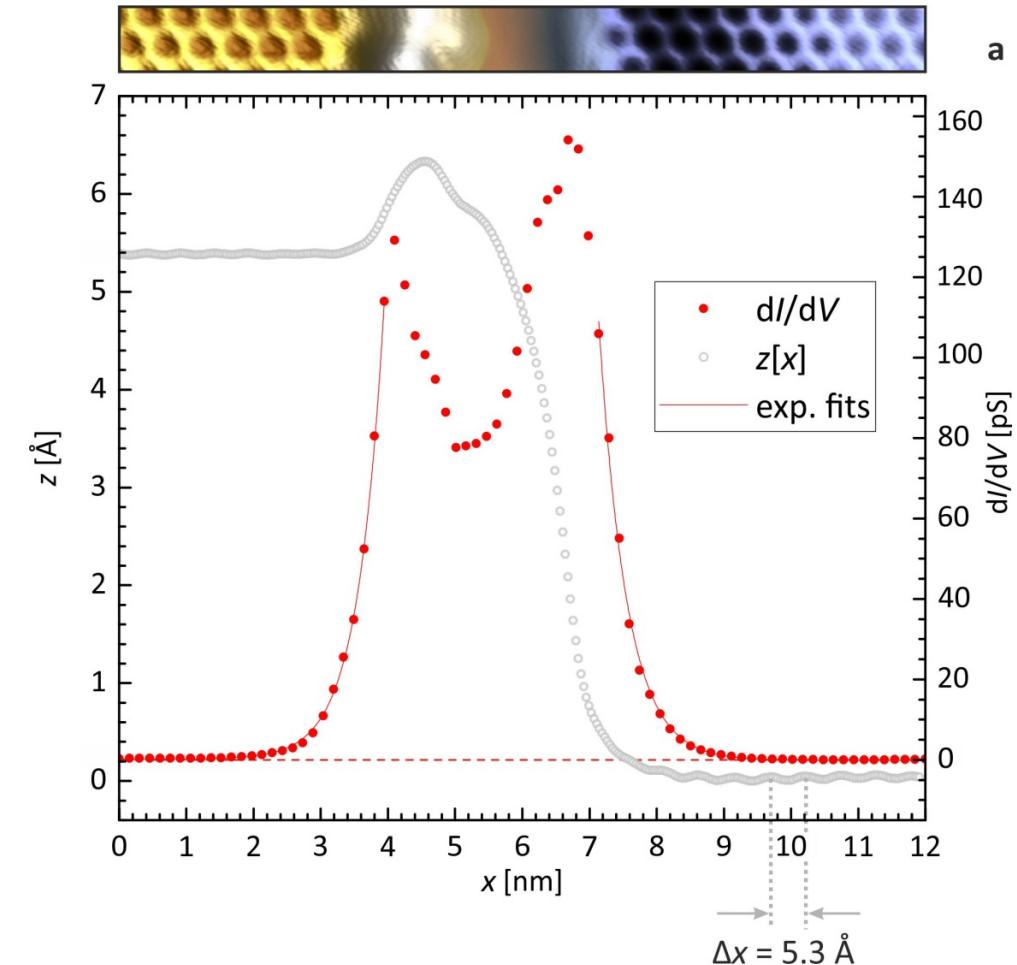
overlay theory/STM



Bismuthene/SiC(0001): edge state spectroscopy



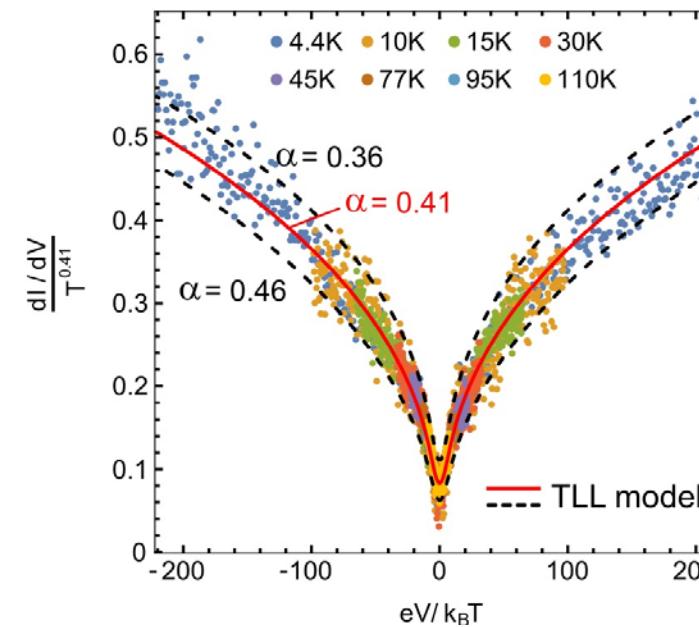
- metallic edge states in bulk gap
- zero bias anomaly (ZBA)



- exponential decay into bulk quantitatively consistent with theory

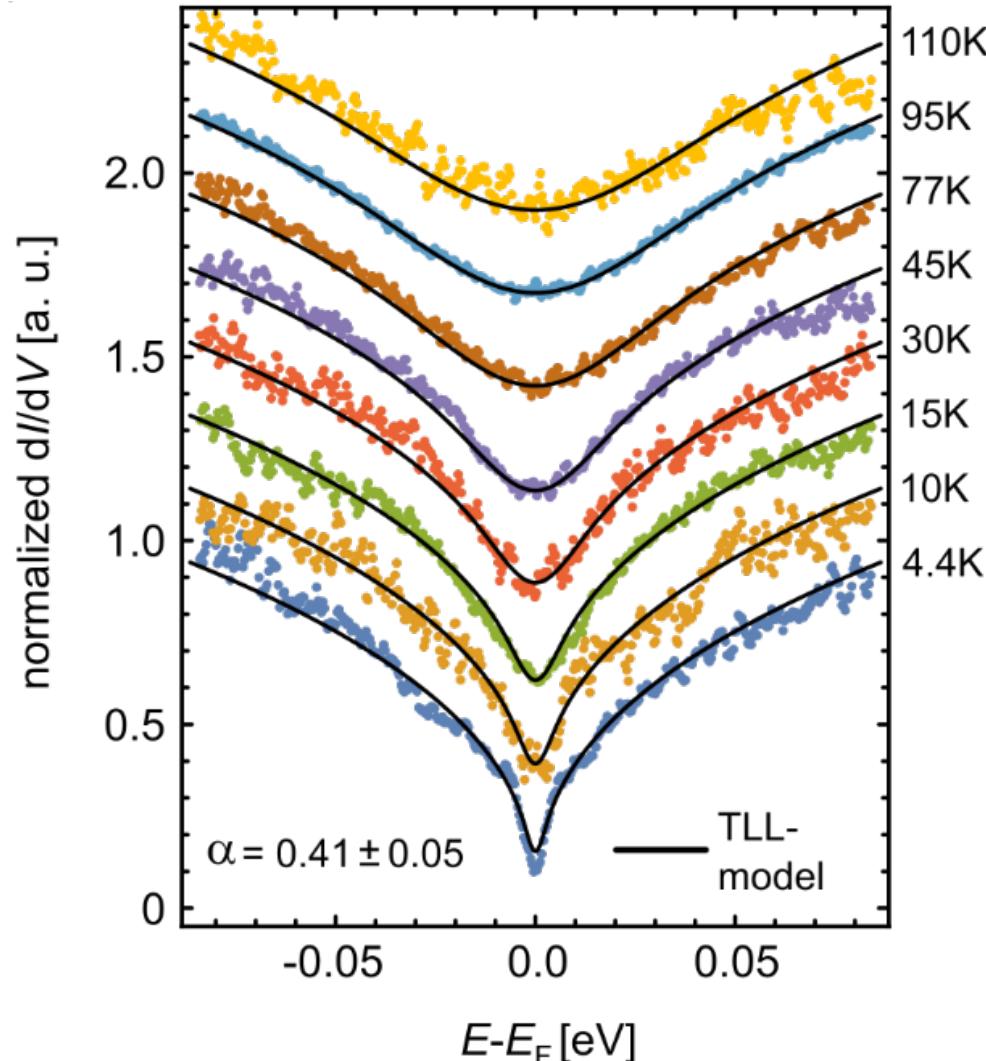


Bismuthene/SiC(0001): edge states as helical Tomonaga-Luttinger liquid

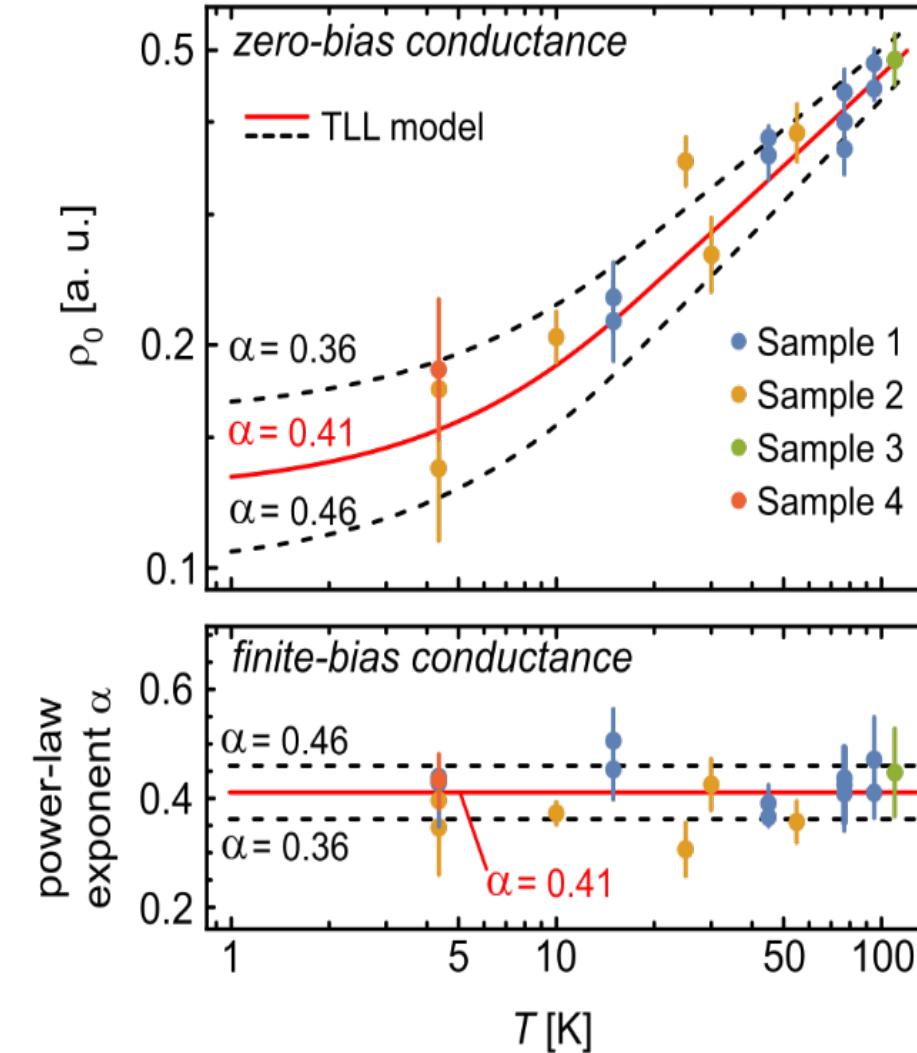
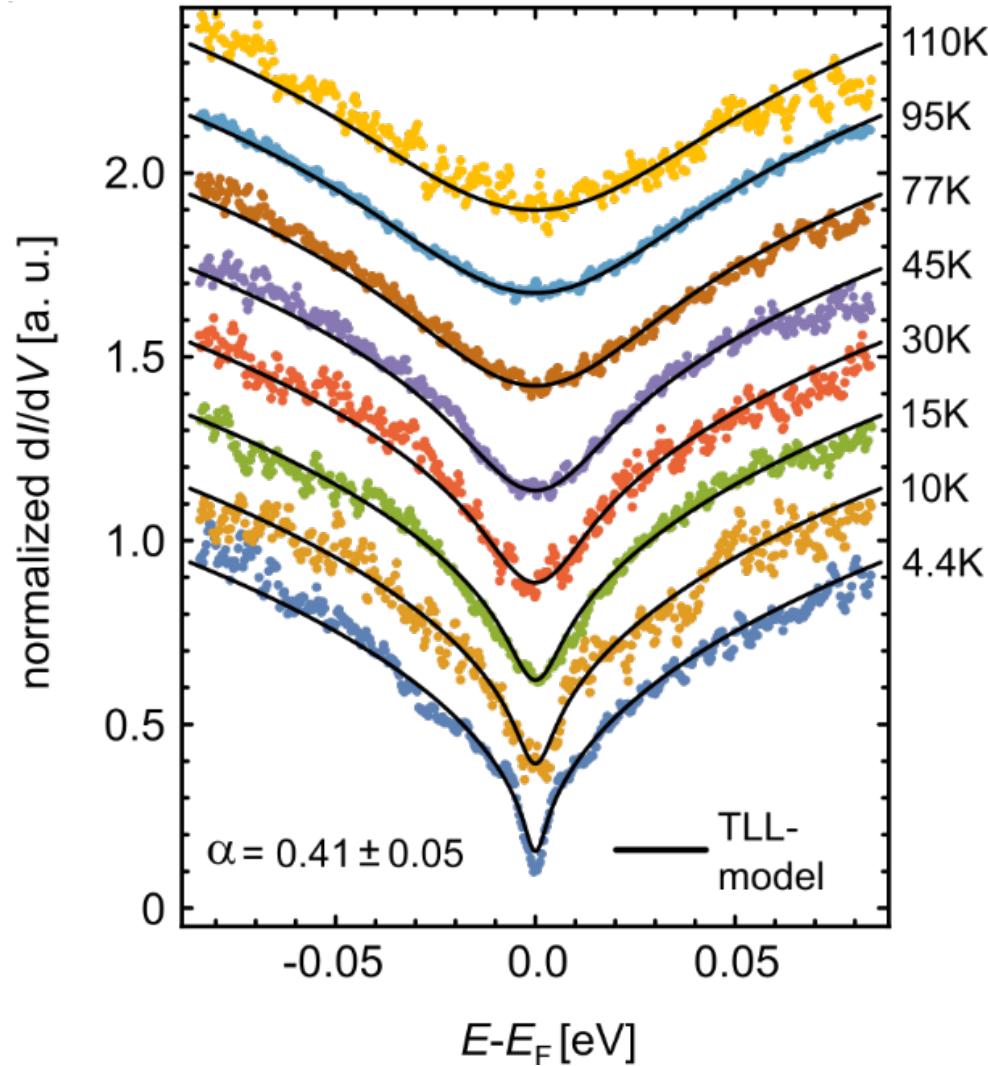


Bismuthene/SiC(0001): zero bias anomaly

ZBA: bias and T dependence



Bismuthene/SiC(0001): zero bias anomaly

ZBA: bias and T dependence

→ power laws

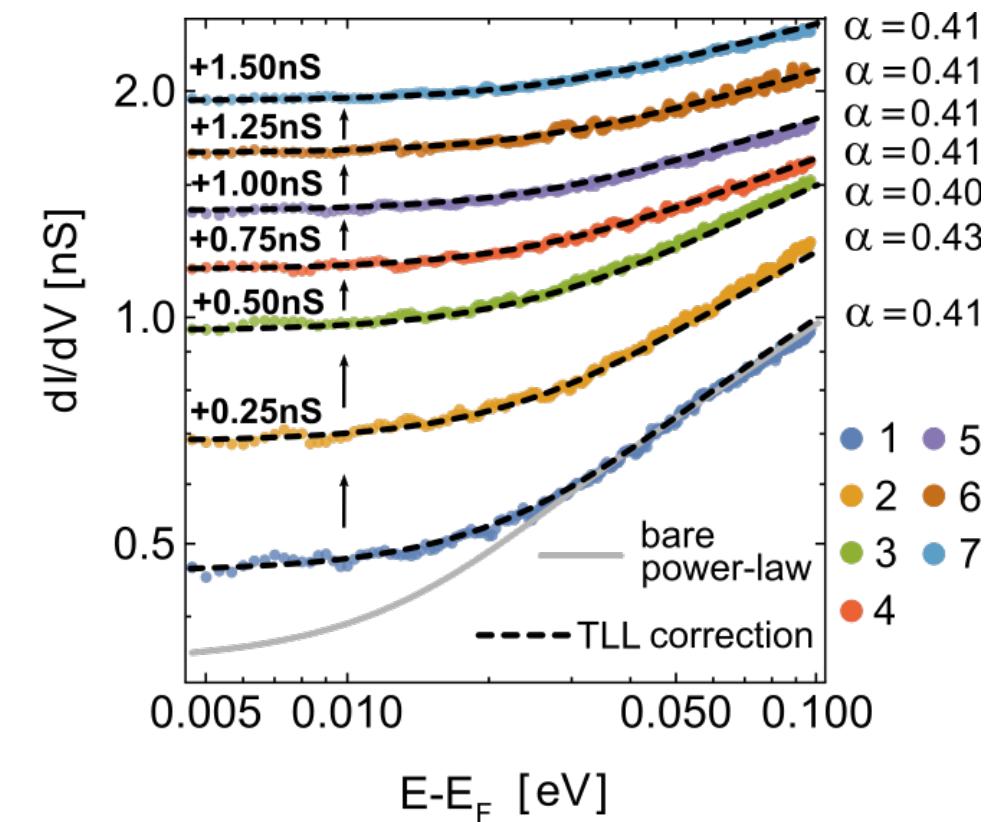
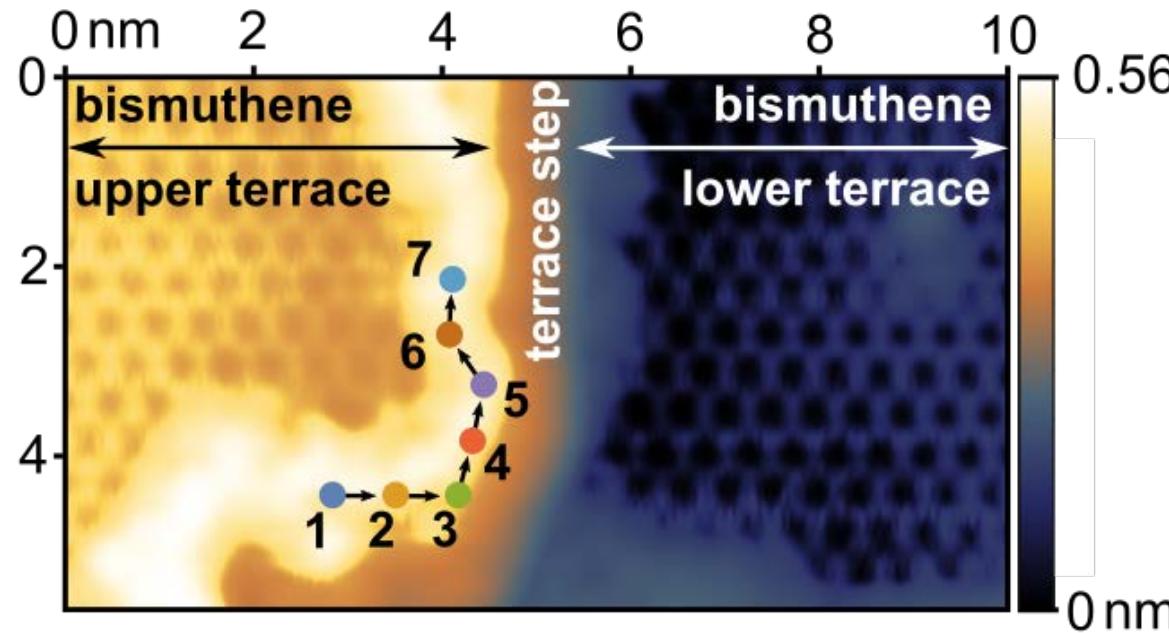
$\propto T^\alpha$

$\propto V^\alpha$

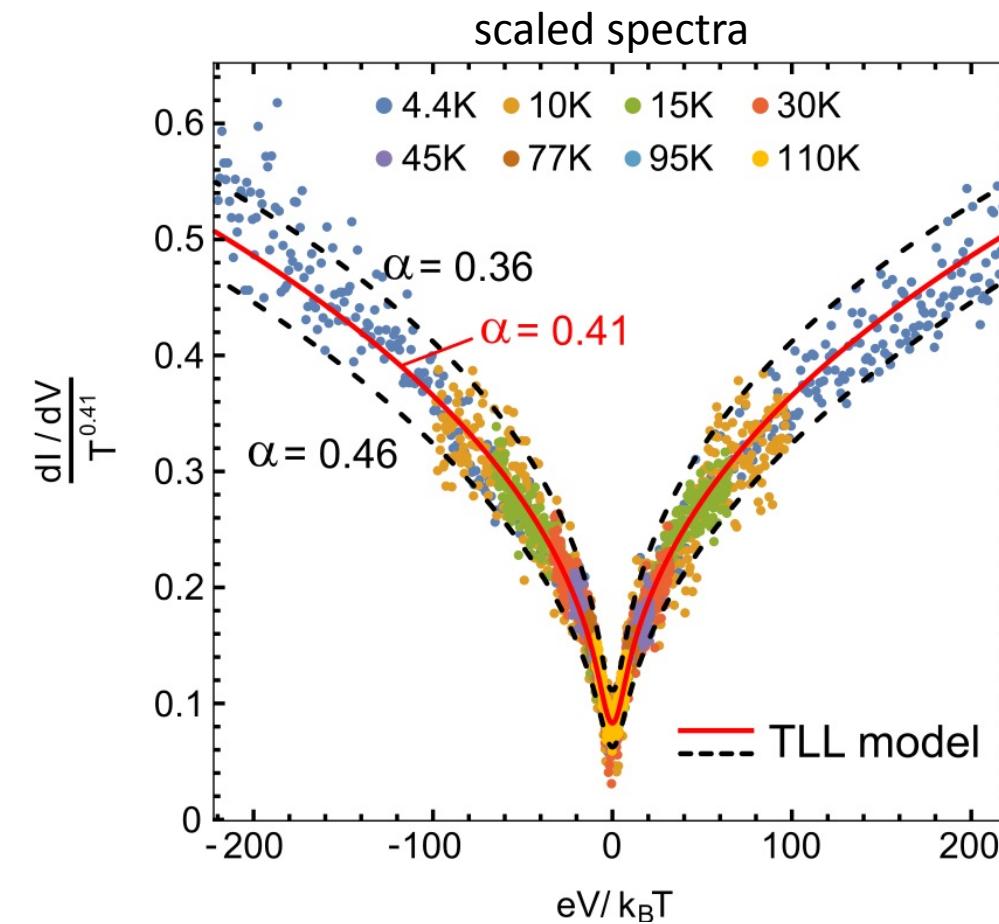
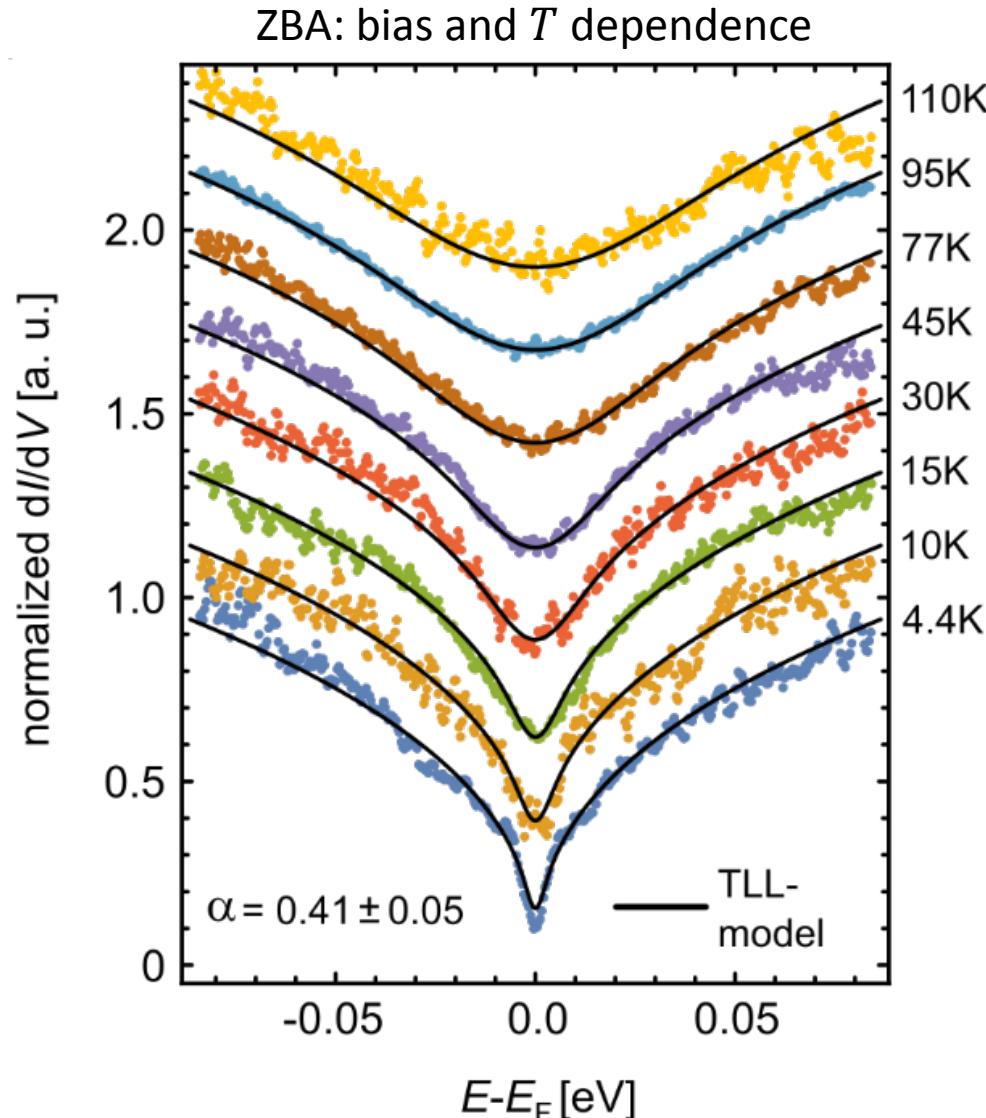
same exponent
 $\alpha = 0.41$

Bismuthene/SiC(0001): zero bias anomaly

persistent power-law dependence of ZBA even in "kinky" edge



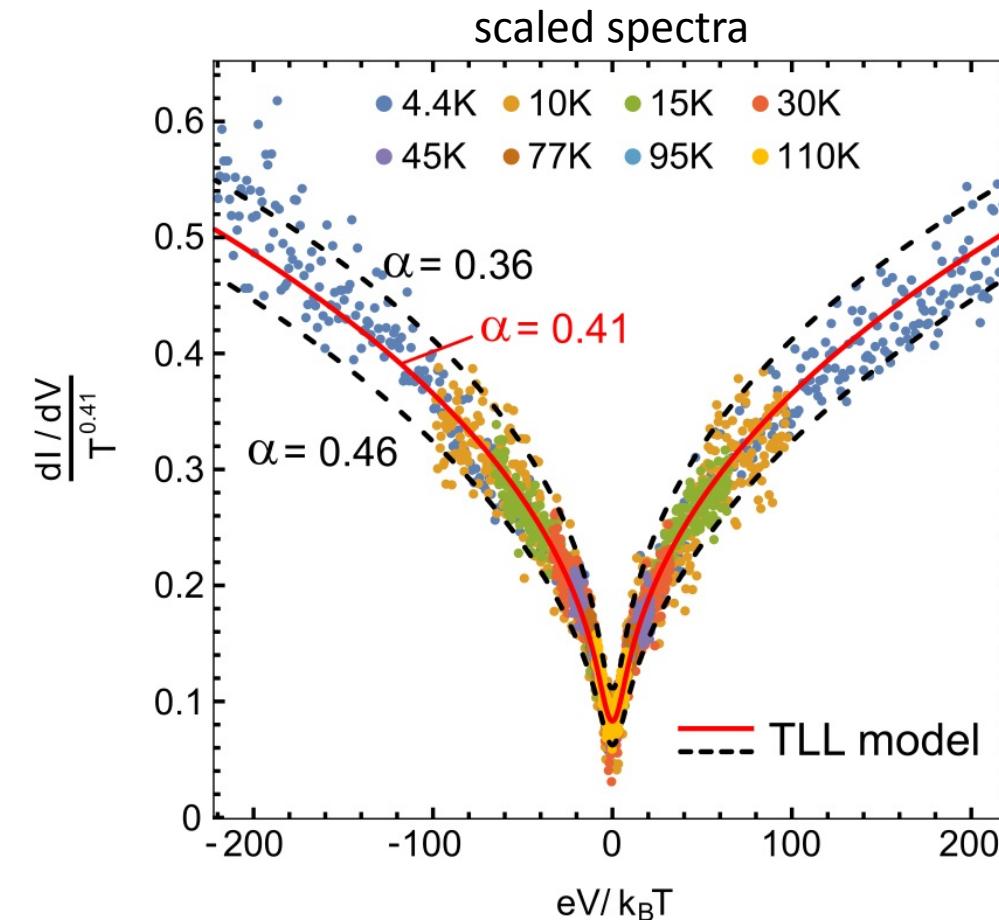
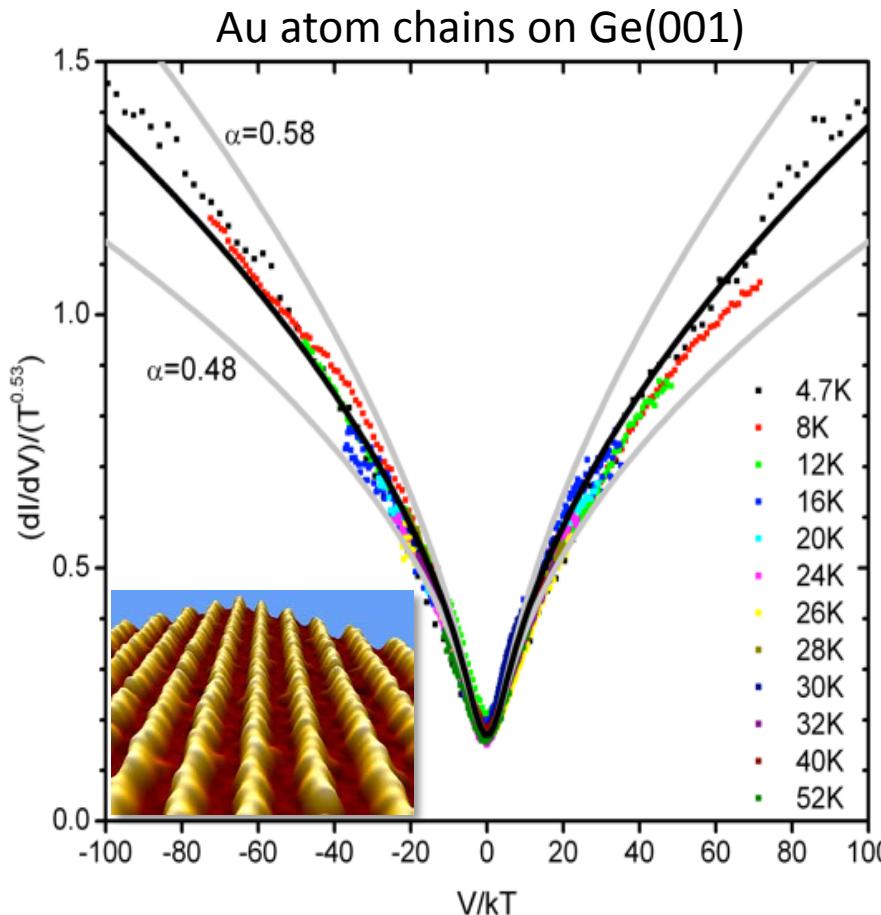
Bismuthene/SiC(0001): zero bias anomaly



same power law in V and T : **universal scaling**

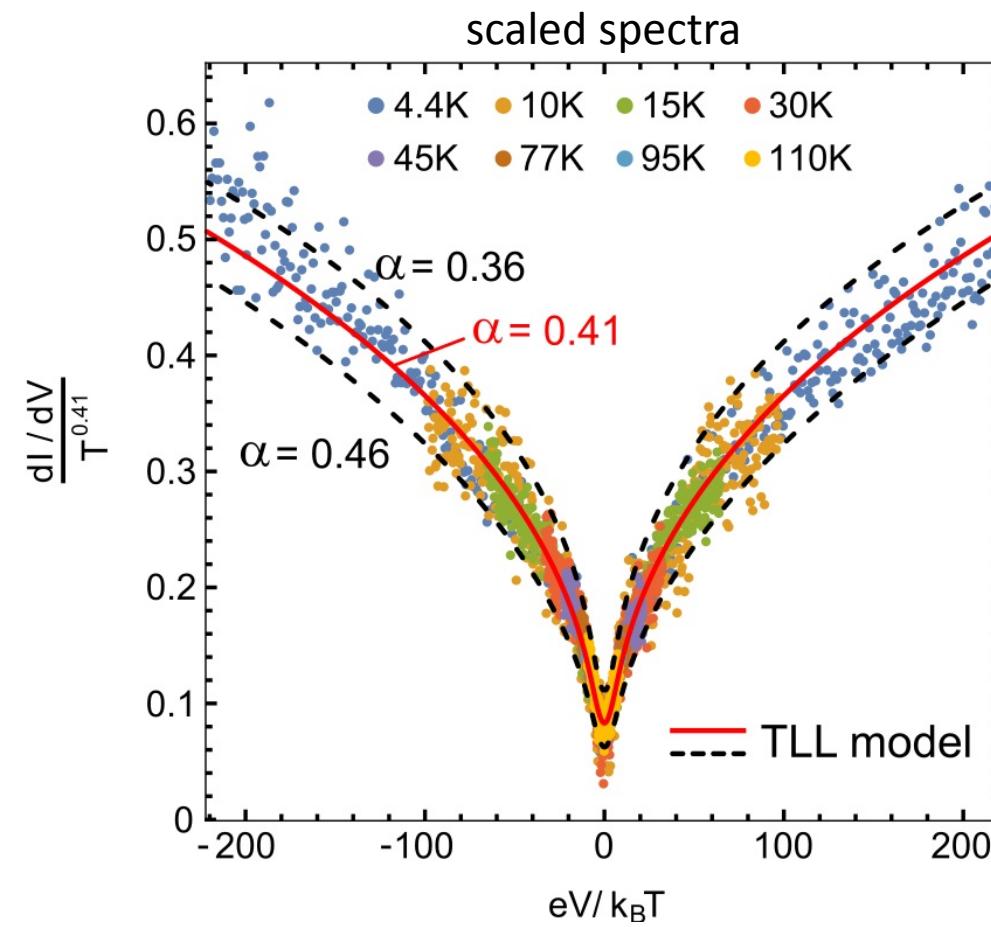
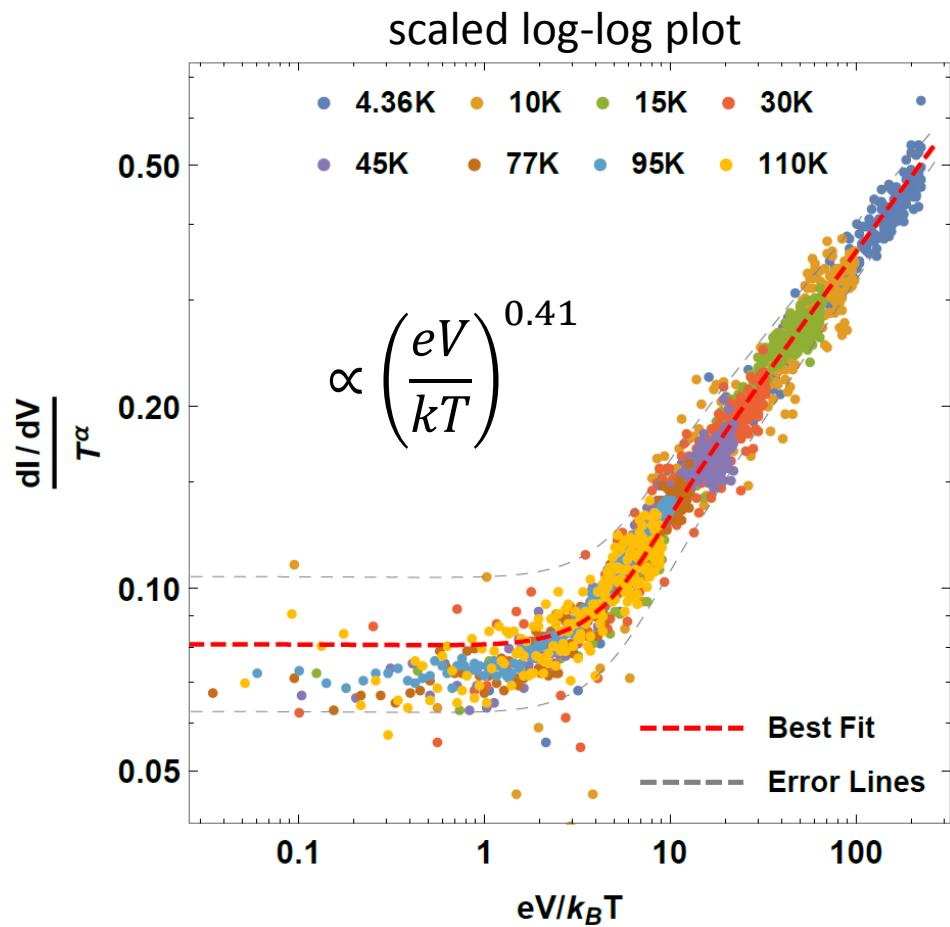
→ el. correlations in 1D: Tomonaga-Luttinger liquid

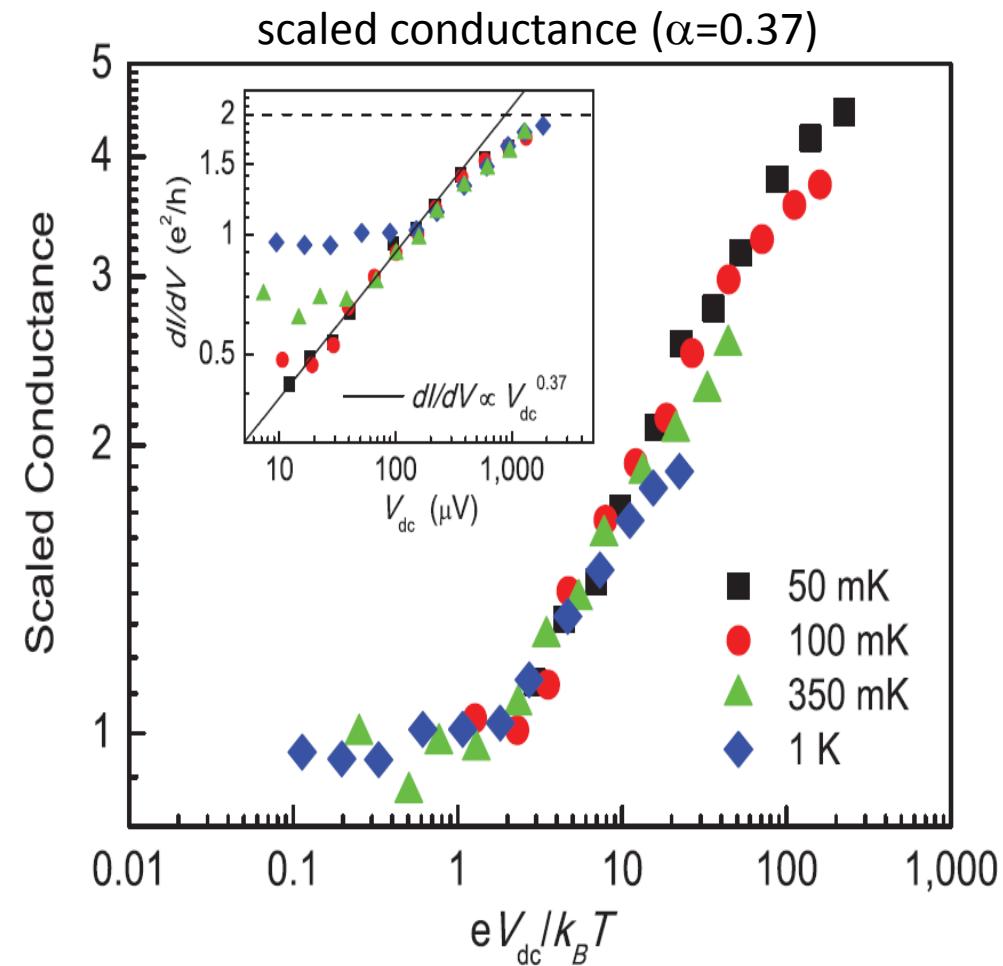
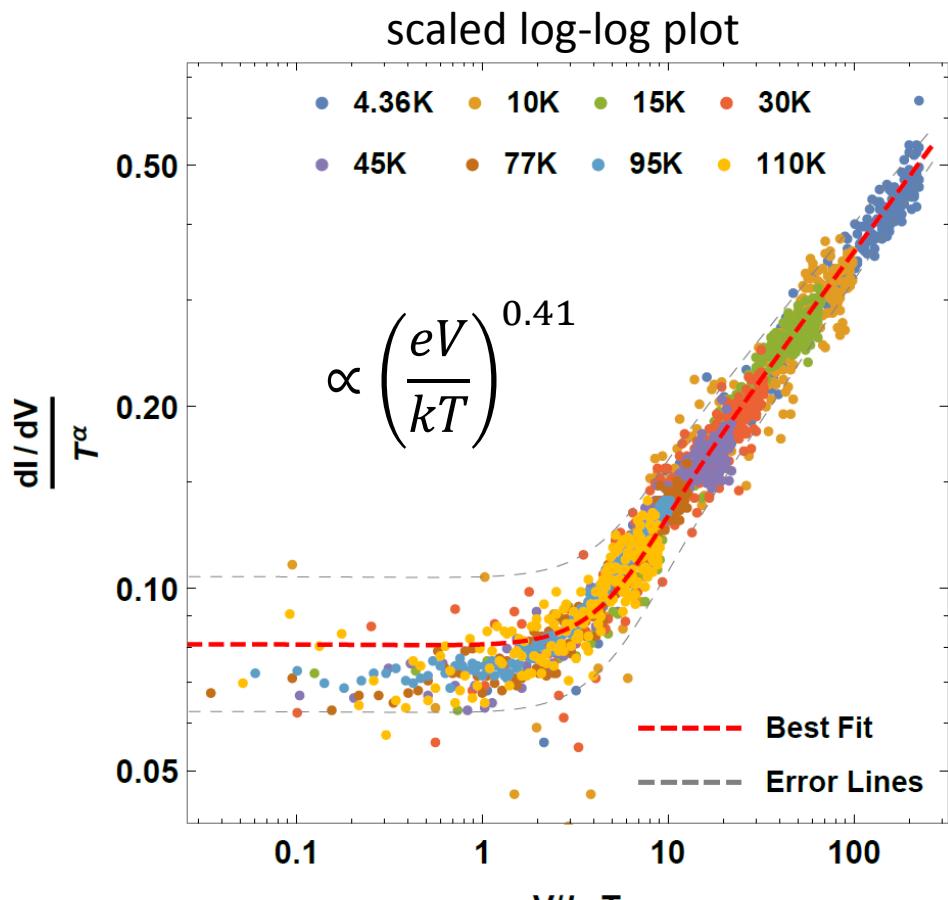
Bismuthene/SiC(0001): 1D edge states as helical TLL



$$\frac{dI}{dV} \propto T^\alpha \cosh\left(\frac{eV}{kT}\right) \left| \Gamma\left(\frac{1+\alpha}{2} + i\frac{eV}{2\pi kT}\right) \right|^2 \otimes \frac{df}{dE}$$

→ el. correlations in 1D: Tomonaga-Luttinger liquid



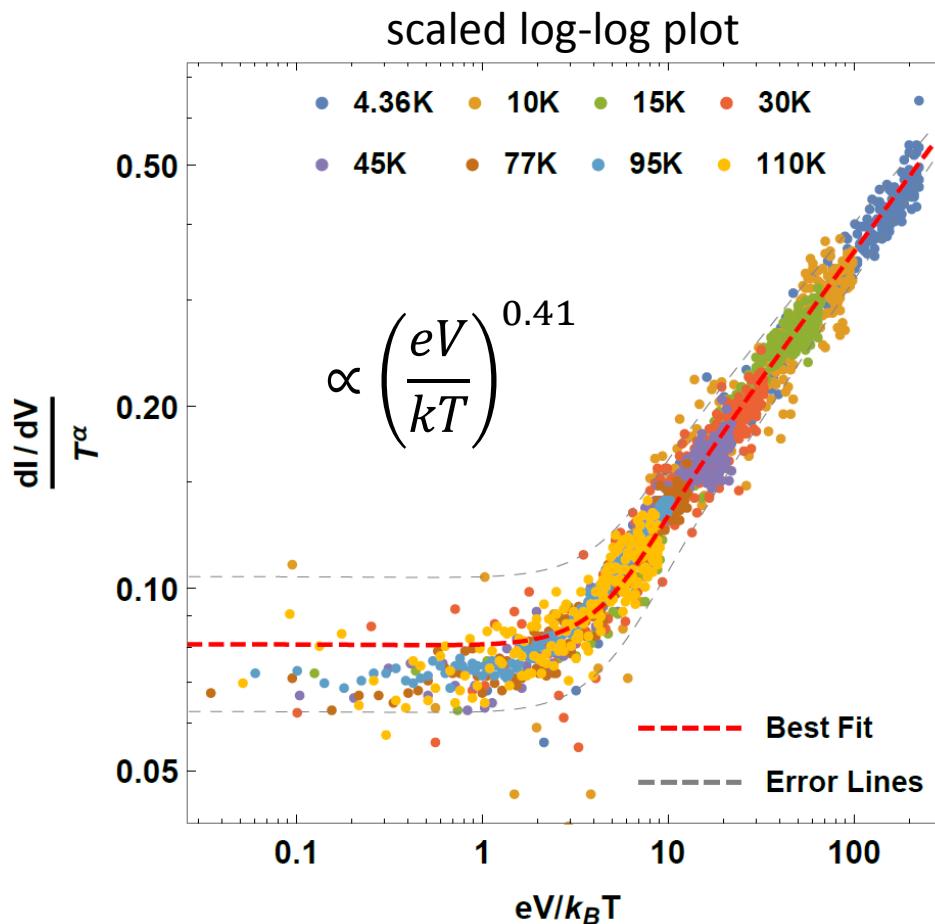


TLL scaling behavior in
Bi^{ene}/SiC

for comparison:
InAs/GaSb QWs

Li et al., PRL 115, 136804 (2015)

Bismuthene/SiC(0001): Tomonaga-Luttinger parameter



TLL scaling behavior in
Bi^{ene}/SiC

$$\text{TLL parameter } K = \sqrt{\frac{1+\tilde{g}_4-\tilde{g}_2}{1+\tilde{g}_4+\tilde{g}_2}} \text{ with } \tilde{g}_i = \frac{g_i}{2\pi\hbar v_F}$$

$$= 0.42 \pm 0.05$$

$$\text{from } \alpha = \frac{1}{2} \left(K + \frac{1}{K} - 2 \right)$$

↑
for helical TLL, see
Braunecker et al., PRB (2012)

	K	source
non-interacting limit	1	
HgTe/CdTe QW	~0.5...0.9	theor. estimate Teo & Kane (PRB 2009)
bismuthene/SiC	0.42	tunneling DOS <i>this work</i>
InAs/GaSb QW	0.21 vs. 0.8	exp. conductance Li et al. (PRL 2015) vs. Väryrynen et al. (PRB 2016)

α -Sn: gap too small for practicable purposes, but nice (Hg-free!) 3D model TI for spectroscopic studies

stanene: potentially interesting QSH insulator, but still needs to be synthesized

bismuthene/SiC:

- **novel QSH paradigm**, extend to other group V elements
- **verify helical nature** of edge states:
 - in \vec{B} -field (characteristic gap openings: *arXiv:1803.02648*)
 - magnetic impurities (local breaking of TRS)
 - ultimately:
QSHE, conductance quantization ($\frac{2e^2}{h}$) at room temperature !

role of 1D correlations?

