# Gapless Spin Liquid Ground State in the S=1/2 Kagome Antiferromagnet



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# Acknowledgment

#### 1. H. J. Liao, et al, PRL 118, 137202 (2017)

#### 2. H. J. Liao, et al, PRB 93, 075154 (2016)

#### 3. Z. Y. Xie, et al, PRX 4, 011025 (2014).





Bruce Normand

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**Ruizhen Huang** 



## **Questions to Address**



S=1/2 Kagome Heisenberg

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j, \quad J > 0$$

Is the ground state

- 1. gapped or gapless?
- 2. a quantum spin liquid?

Herbertsmithite: ZnCu<sub>3</sub>(OH)<sub>6</sub>Cl<sub>2</sub>

#### ✓ Valence bond crystal

	Singh & Huse, PRB	2008	series expansion
	Evenbly & Vidal, PRL	2010	MERA
	Iqbal, Becca & Poilblanc, PRB	2011	VMC
$\checkmark$	<b>Gapped Z<sub>2</sub> spin liquid (Topological)</b>		
	Sachdev, PRB	1992	Schwinger boson
	Gotze, et al, PRB	2011	coupled cluster
	Jiang, Weng & Sheng, PRL	2008	DMRG
	Yan, Huse & White, Science	2011	DMRG
	Depenbrock, McCulloch & Schollwock, PRL	2012	DMRG
	Jiang, Wang & Balents, Nature Physics	2012	DMRG
	Nishimoto, Shibata, Hotta Nat. Commun.	2013	DMRG
	Gong, Zhu & Sheng, Scientific Reports	2014	DMRG
	Li, arXiv:1601.02165	2016	VMC
	Mei, Chen, He & Wen, arXiv:1606.09639	2016	SU(2)-TNS
✓	Gapless spin liquid (Algebra)		
	Ran, Hermele, Lee, Wen, PRL	2007	VMC
	Iqbal, Becca, Sorella, Poilblanc, PRB	2013	VMC+Lanczos
	Iqbal, Poilblanc, Becca, PRB & 1606.02255	2015	VMC
	Hu, Gong, Becca & Sheng, PRB	2015	VMC
	Jiang, Kim, Han & Ran, arXiv:1610.02024	2016	SU(2)-PEPS
	Liao et al, arXiv:1610.04727, PRL 2017	2016	PESS
	He, Zaletel, Oshikawa, Pollmann, 1611.06238	2016	DMRG

#### Hints from Experiments

# Fractionalized excitations in the spin-liquid state of a kagome-lattice antiferromagnet

Nature 492 (2012) 406

Tian-Heng Han<sup>1</sup>, Joel S. Helton<sup>2</sup>, Shaoyan Chu<sup>3</sup>, Daniel G. Nocera<sup>4</sup>, Jose A. Rodriguez-Rivera<sup>2,5</sup>, Collin Broholm<sup>2,6</sup> & Young S. Lee<sup>1</sup>

#### Gapless spin liquid



Along the (H, H, 0) direction, a broad excitation continuum is observed over the entire range measured

Herbertsmithite  $ZnCu_3(OH)_6Cl_2$ : Neutron scattering

#### Hints from Experiments

#### Evidence for a gapped spin-liquid ground state in a kagome Heisenberg antiferromagnet

Science 360 (2016) 655

Mingxuan Fu,<sup>1</sup> Takashi Imai,<sup>1,2</sup>\* Tian-Heng Han,<sup>3,4</sup> Young S. Lee<sup>5,6</sup>

#### Gapped spin liquid



NMR Knight shift

 $\Delta(0)/J = 0.03$  to 0.07

## Problems in the theoretical studies

✓ Density Matrix Renormalization Group (DMRG):

strong finite size effect

error grows exponentially with the system size



Depenbrock et al, PRL 109, 067201 (2012)

# Problems in the theoretical studies

✓ Density Matrix Renormalization Group (DMRG):

strong finite size effect

error grows exponentially with the system size

✓ Variational Monte Carlo (VMC)

need accurate guess of the wave function

✓ Quantum Monte Carlo

Minus sign problem

## **Tensor-Network States**

- 1. A variational wave function that satisfies the area law of entanglement entropy
- 2. Control parameter: bond dimension **D**

the wave function is exact in the  $D \rightarrow \infty$  limit







Kagome Frustrated lattice

Local tensors defined on the honeycomb lattice Rank-3 tensors, no frustration

# **Comparison between PEPS and PESS**

Projected Entangled Pair State (PEPS)

tensors defined on the original lattice

Rank-5 tensors High cost

Virtual spins at two neighboring sites form a maximally entangled state Projected Entangled Simplex State (PESS)

tensors defined on honeycomb lattice

Rank-3 tensors Low cost

Virtual spins at each simplex form a maximally entangled state

## Advantage in using tensor-network states

- 1. No finite lattice size effect: PESS is defined on an infinite lattice
- 2. Most accurate method for studying large lattice size systems



Stoudenmire and White, Annu. Rev. CMP 3, 111(2012)

# Advantage in using tensor-network states

- 1. No finite lattice size effect: PESS is defined on an infinite lattice
- 2. Most accurate method for studying large lattice size systems
- 3. The ground state energy converges fast with the increase of the bond dimension *D* 
  - Converge exponentially with D if the ground state is gapped
  - Converge algebraically with D if the ground state is gapless
    We use this property to determine whether the ground state is gapped or gapless

# Disadvantage: Cost is very high



	Double-layer	
Computational Cost	$O(D^{12})$	
Memory Cost	$O(D^8)$	
Limit of D	13	

# Reduce the Cost by Dimension Reduction



	Double-layer	Shifted single-layer
Computational Cost	$O(D^{12})$	$O(D^8)$
Memory Cost	$O(D^8)$	$O(D^6)$
Limit of D	13	<b>25</b> (not use symmetry)

# S=1/2 Kagome Heisenberg: Ground State Energy



Is this *D* large enough?

## Kagome Heisenberg: Gapless



Energy converges algebraically with the bond dimension

#### Results obtained on the Husimi lattice provide good references

Make comparison between Kagome and Husimi results



(b) Husimi Lattice

- ✓ Highly frustrated
- $\checkmark$  D is generally less than 20

- ✓ Tree Structure
- ✓ Tensor renormalization is rigorous, *D* can reach 1000

## S=1/2 Husimi Lattice: Gapless Spin Liquid



Both energy and magnetization converge algebraically with D

### S=1 Husimi: Gapped Ground State



Energy converges exponentially with the bond dimension

### Kagome Magnetization: magnetic order free

 $M_{Kagome} < M_{Husimi}$ 



Magnetization: decays algebraically with D

Stability of the gapless spin-liquid state against other interactions



#### Bond dimension dependence of the magnetic order



#### Bond dimension dependence of the magnetic order



#### Bond dimension dependence of the magnetic order



# Kagome $J_1 - J_2$ model: phase diagram of infinite D limit



✓ We have performed a large scale tensor
 renormalization group calculation for the S=1/2
 Kagome Heisenberg model

✓ Our result suggests that the ground state of this system is a gapless quantum spin liquid