

Congratulations to KITS

Interaction Driven Topological Phases in Two-Dimensional Systems

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Research is supported by NSF (magnetic systems)
and DOE (topological matter)

Research collaborations

Shoushu Gong (FS), Wei Zhu (LANL)
Tiansheng Zeng (CSUN), Wenjun Hu (RU)

Y. C. He and Yan Chen, Balents on CSL;
Liang Fu on Spontaneous QHE

Duncan Haldane, F. Becca
Kun Yang, O. Starykh

Outline

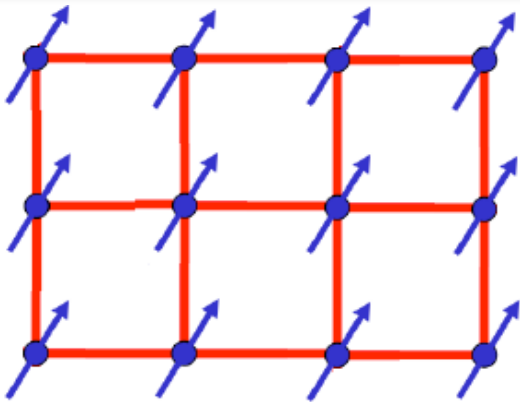
- I. **Introduction:** Quantum spin liquid (SL) and topological order;
Time reversal symmetry (TRS) spontaneously breaking Chiral spin liquid (CSL) proposed 30 years ago ---no known examples until our recent work

- II. **Establish robust CSL in Kagome SU2 model (J1-J2-J3) small perturbation J2-J3, ~ 0.07)**
comparing to other spin models (triangular etc)
VMC simulations, indication of Dirac to CSL transition

- III. **Frustration in Electron Systems:
Spontaneous Quantum Hall Effect**

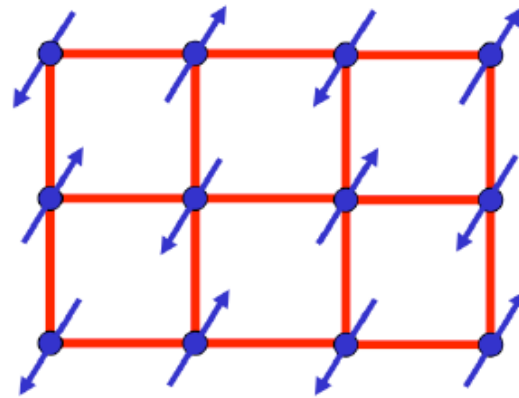
Magnetic systems can develop orders by breaking symmetry

$J < 0$

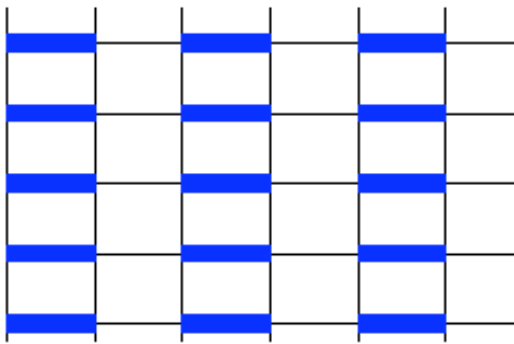


Ferromagnetic order

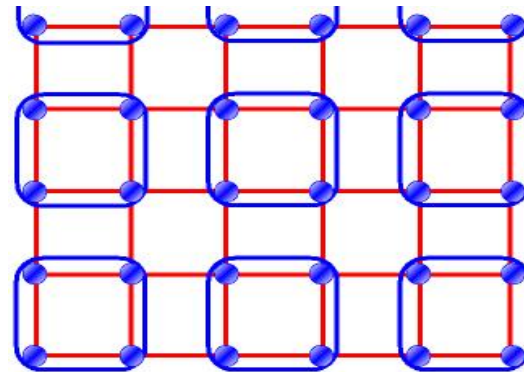
$J > 0$



Antiferromagnetic order with NN J_1

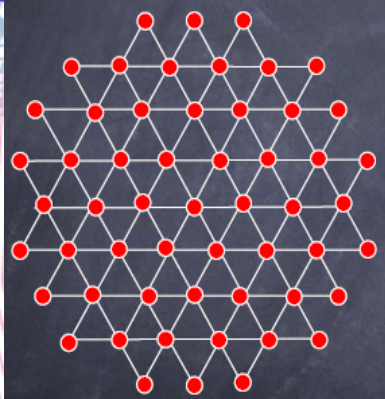


Valence Bond Solid (VBS) order

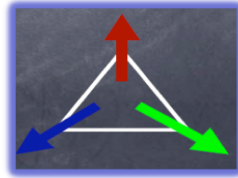


Plaquette VBS (possibly identified in J_1 - J_2 square model)

Spin disordered phase in Heisenberg models with increasing frustration



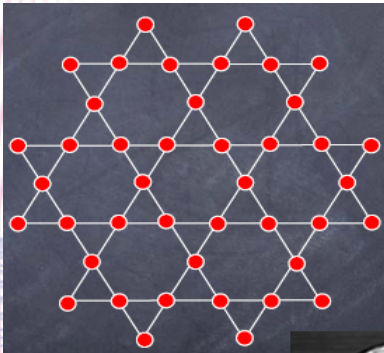
Triangular lattice



Three-sublattice AF order

Resonating valence bond (P. W. Anderson),
RVB is a quantum spin
liquid (SL) 1973, 1989

NN Kagome lattice:
A primary candidate
for Z₂ SL



Spin Liquid

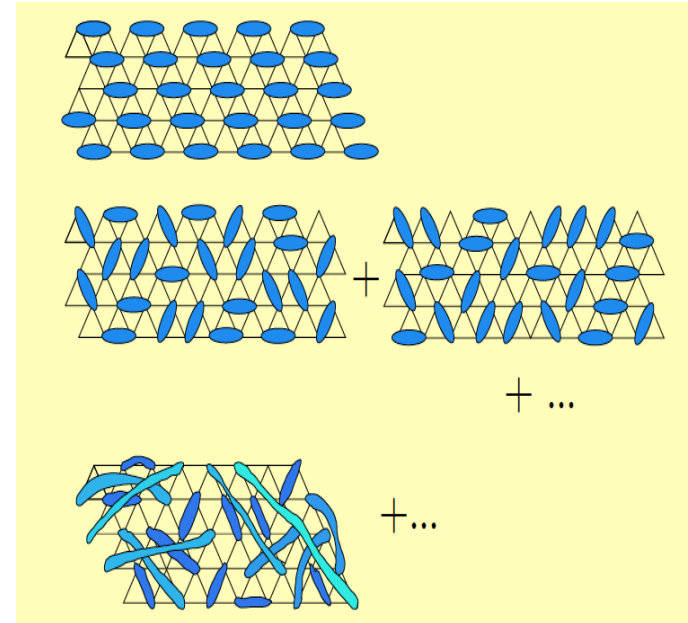
H.C. Jiang, Z.Y. Weng, DNS
(2008), not converged DMRG
for wider systems

Yan, Huse, White,
Science (2010)---new
milestone: A possible
gapped spin liquid

Jiang, Wang,
Balents (2012)
Depenbrock et al.
(2012).

Spin Liquid (SL) State

A new state of matter with no symmetry broken, with topological order /entanglement, and fractionalized excitations X. G. Wen (1990, 1991)



Z₂ SL in contrived theoretical models

Wen (1990,1991), Kivelson, Rokhsar, Sethna (1987), Senthil, Fisher (2000), Balents, Fisher, Girvin (2002) Moessner and Sondhi, (2001), Senthil, Motrunich (2002) Senthil, Vishwanath, Balents, Sachdev, Fisher (2004) PA Lee (2008), Balents (2010)

Gapped SL State

Gaps to all spin excitations
exponential decay correlations
J₁-J₂-J₃ (Ising) kagome, Balents et al and J₁ kagome model

Time reversal symmetry (TRS) broken chiral Spin liquid (CSL): a bosonic FQHE state

Kalmeyer and Laughlin 1987

Wen, Wilczek, Zee 1989

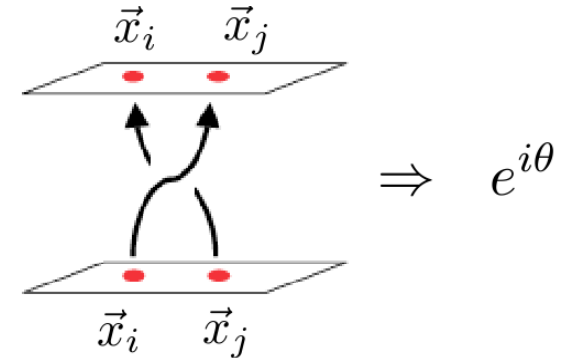
anyon quasiparticles obey
fractional statistics

(bosonic)FQHE without B field

Haldane and Arovas 1995

Yang, Warman and Girvin 1993.

Yao and Kivelson: a contrived
CSL state (for Kitaev model)
2007



Induce CSL using
(S1xS2)*S3 interaction

Schroeter et al 2007, Thomale et al.
2009, Nielsen et al. 2012, Bauer et al.
2013-2014 (Mott materials)

Approx. Methods find CSL
Hermele et al (2009), Messio et al.
For kagome J1-J2-J3 (2012) Cuboc

Experimental Candidates

Kagome magnets

Bilayer magnets

Triangular organic compounds

Rare-earth spin-orbit coupling
inorganic compounds

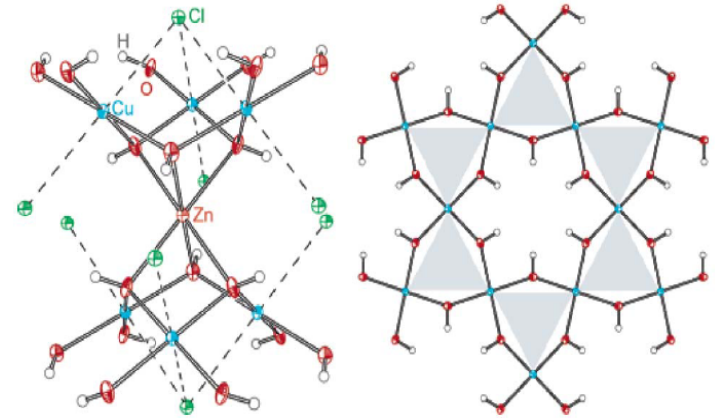
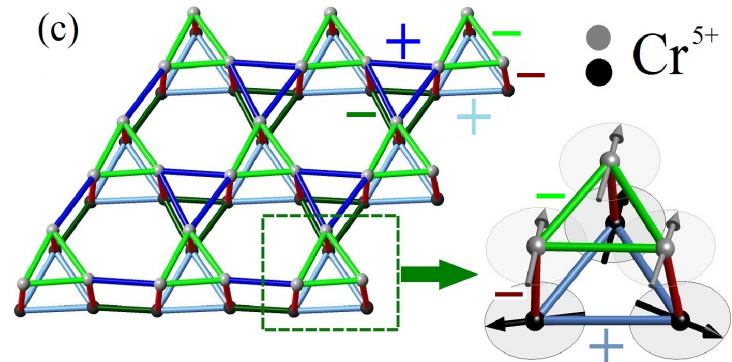
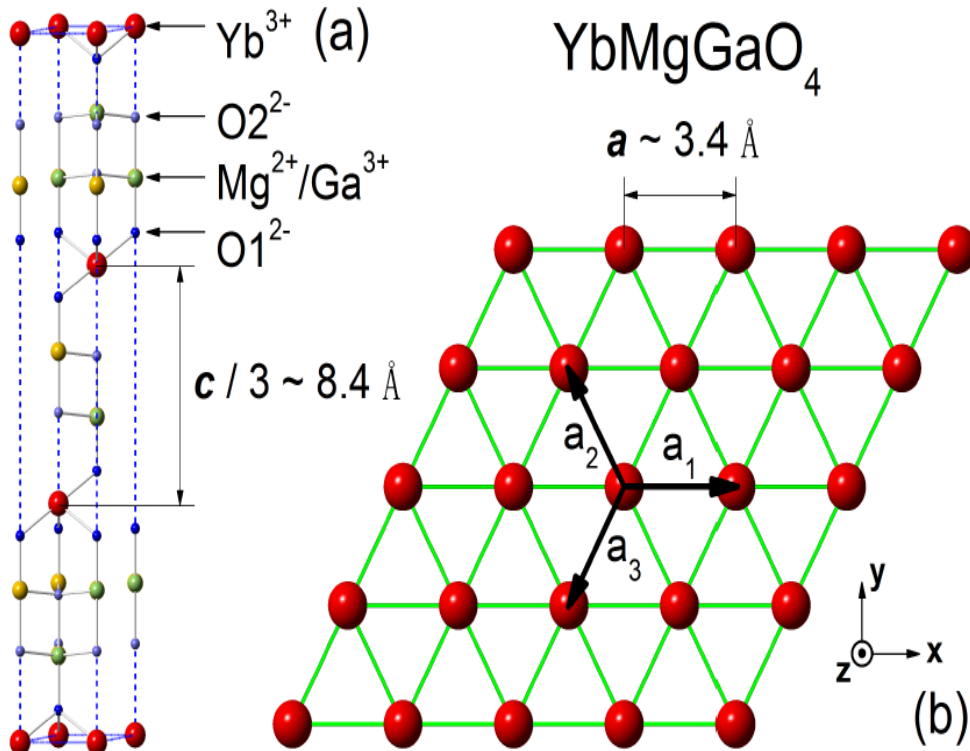


Figure 1. Crystal structure of Zn-paratacamite (1), $\text{Zn}_{0.33}\text{Cu}_{3.67}(\text{OH})_6\text{Cl}_2$.

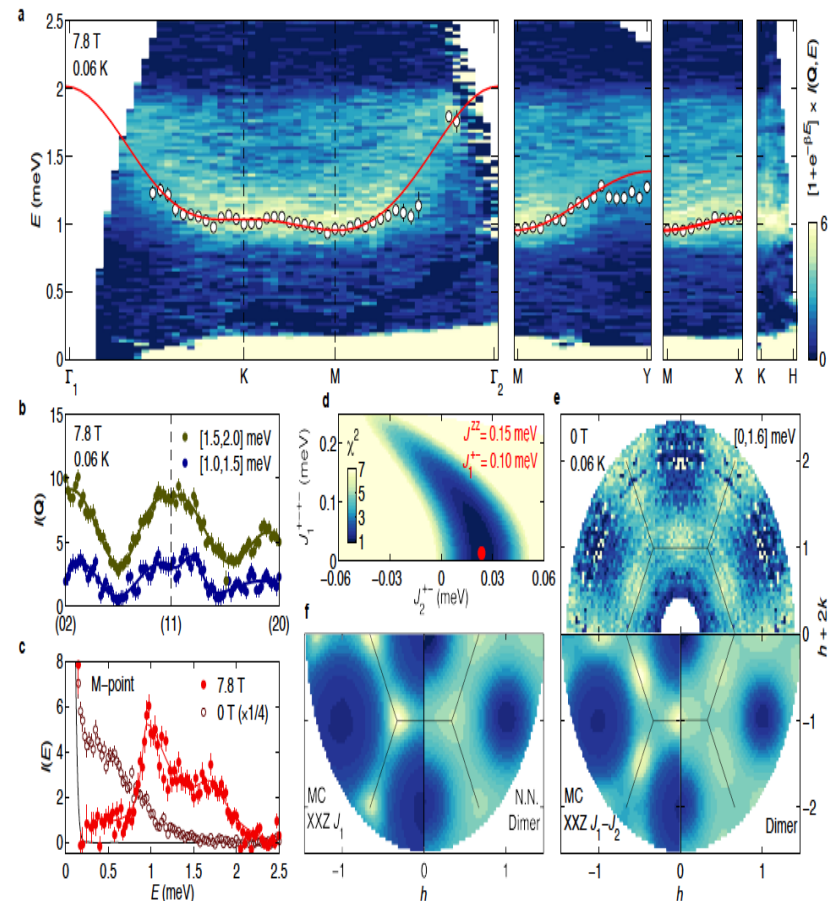


Rare-earth spin-orbit coupling triangular inorganic compounds

Y. Li et al 2015, PRL



Paddison et al 2016



J1-J2 Heisenberg models with spin nonconserved terms (S+S+)

Numerical Studies

What phases are realized in realistic simple spin models with strong frustration ?

--- density matrix renormalization group (DMRG) numerical studies

DMRG convergence is controlled by kept states (m)
 $m=6000\text{--}30000$ (U1 equivalent) states
---utilizing **SU2 symmetry** (Gong et al.)

S. R. White 1992

1D DMRG algorithm



Mapping a 2D lattice onto a 1D

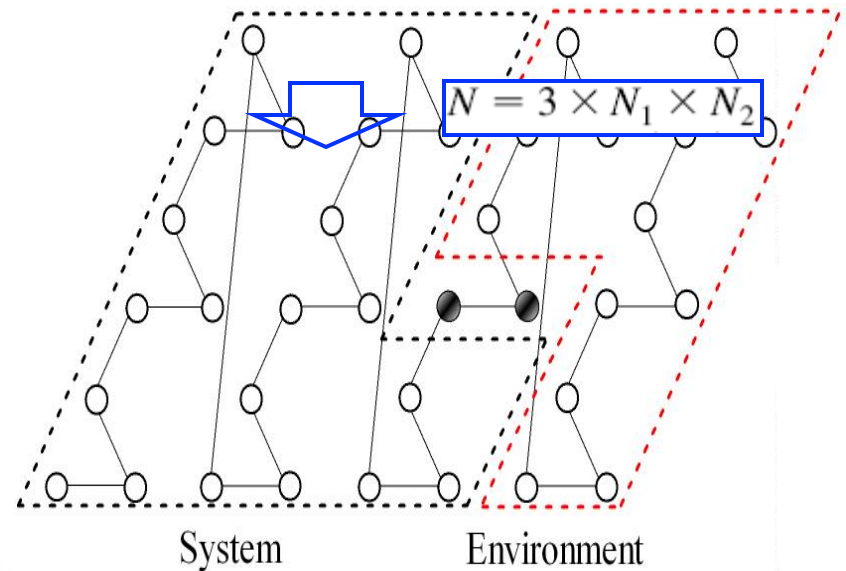
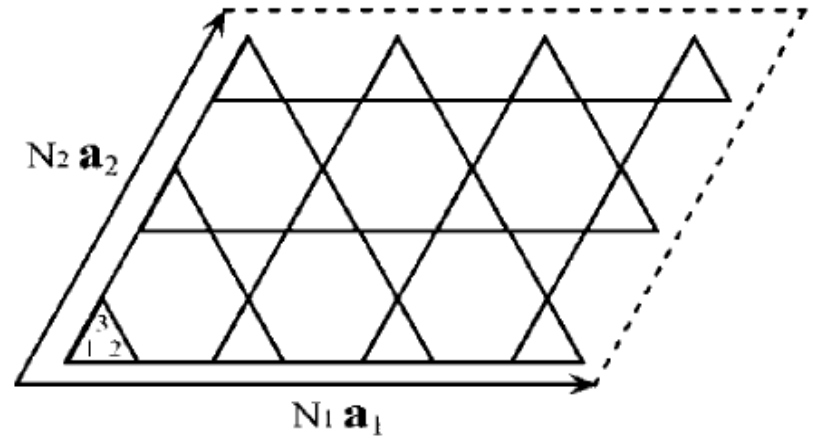
Jiang, ZY Weng, DNS(2008)

DMRG with Su2 invariant

Depenbrock et al. 2012

Shoushu Gong 2012—2014

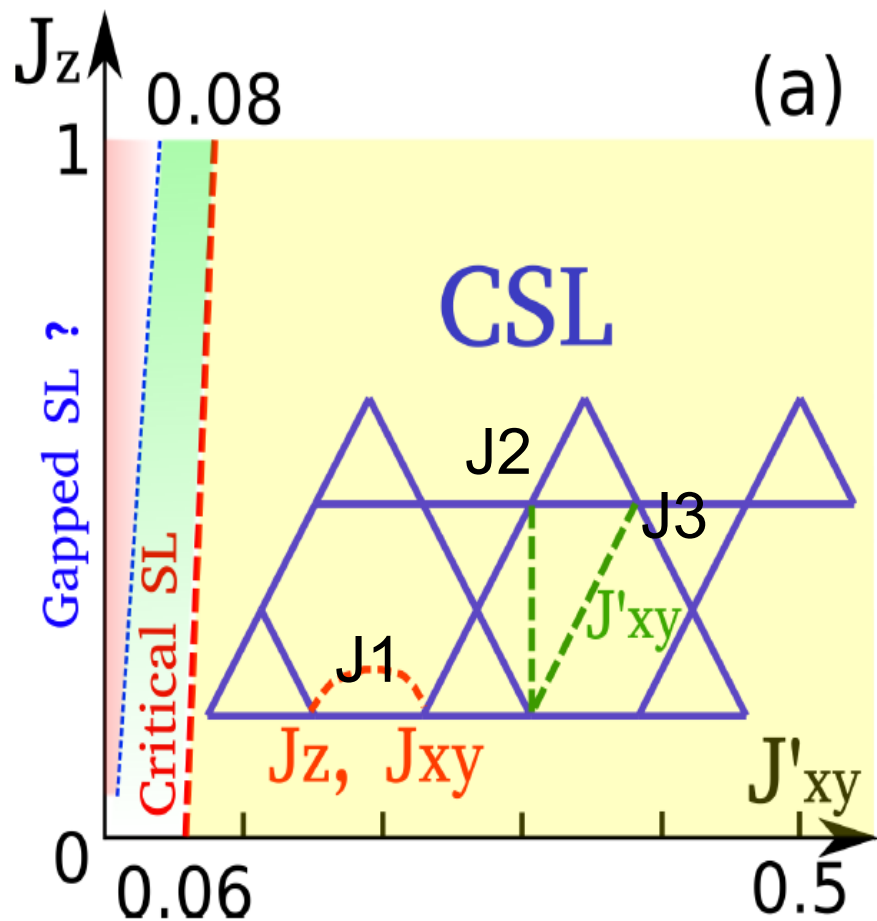
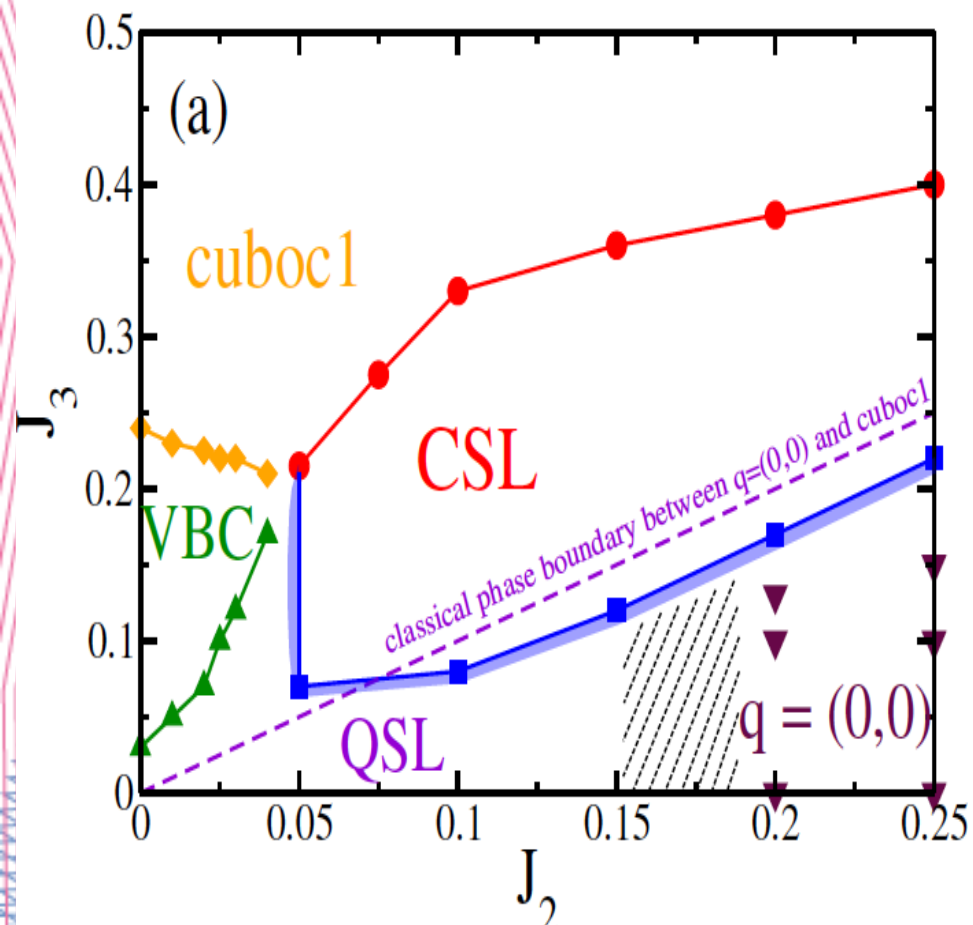
~30,000 states



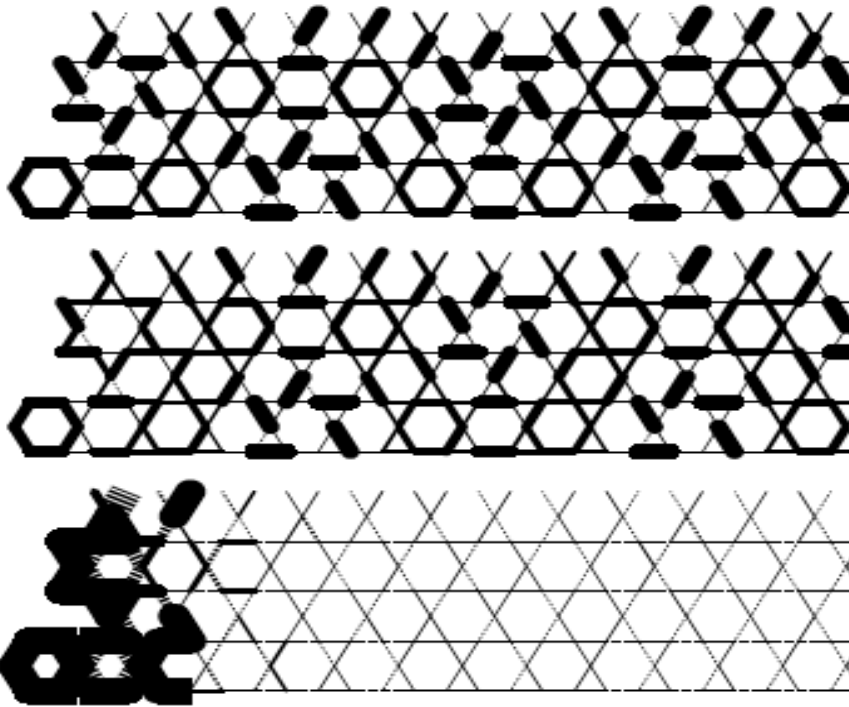
Discovering of Chiral Spin Liquid (CSL) in kagome spin $\frac{1}{2}$ models

J1-J2-J3
Heisenberg
model

$$H = J_1 \sum_{\langle i,j \rangle} S_i \cdot S_j + J_2 \sum_{\langle\langle i,j \rangle\rangle} S_i \cdot S_j + J_3 \sum_{\langle\langle\langle i,j \rangle\rangle\rangle} S_i \cdot S_j$$

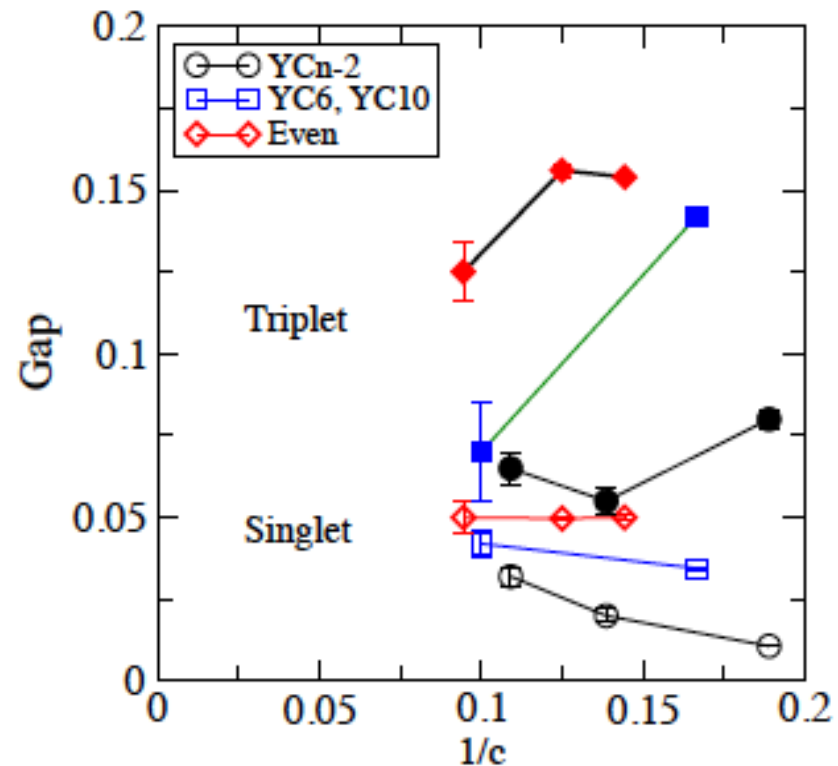


Main evidence for possible SL for J1 kagome spin1/2 (DMRG, Yan et al)

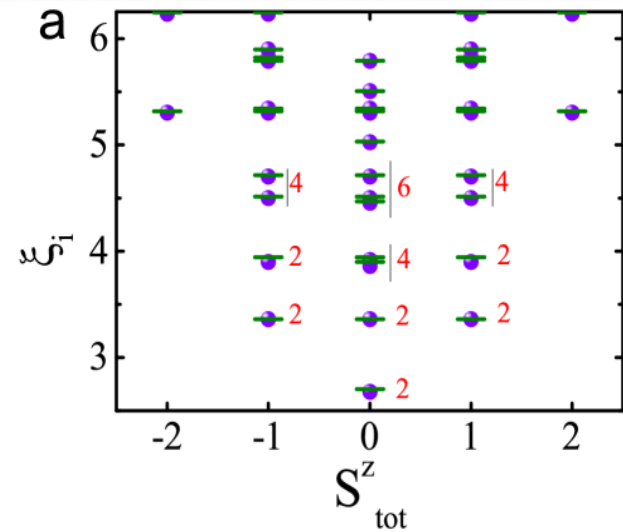
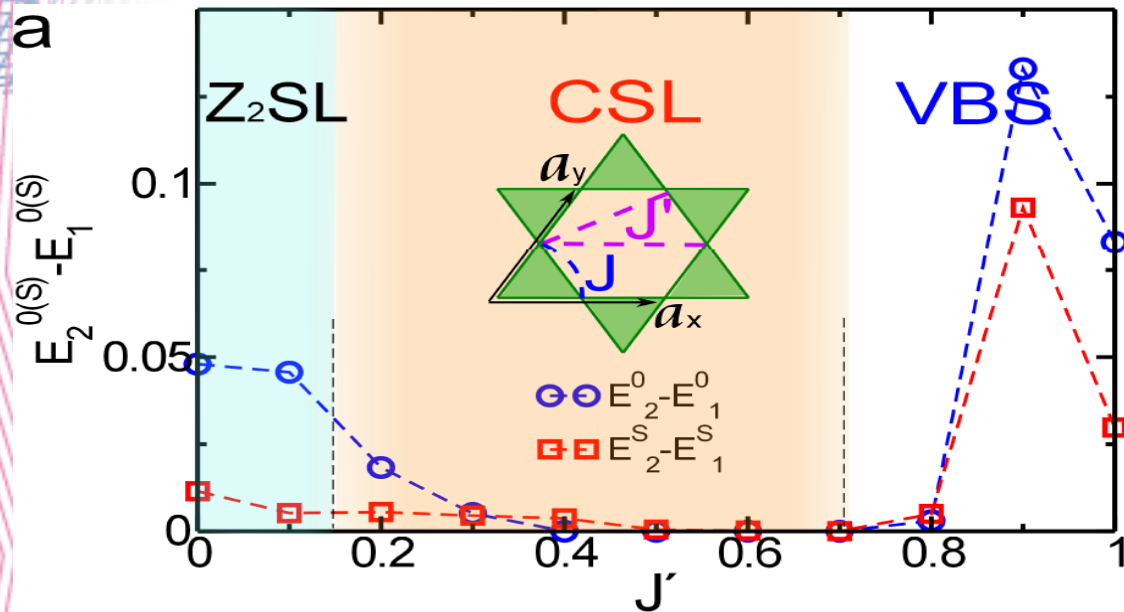


Yan, Huse, White, 2010 found two topological sectors in DMRG; Challenge to find four sectors to fully establish topological nature of the phase

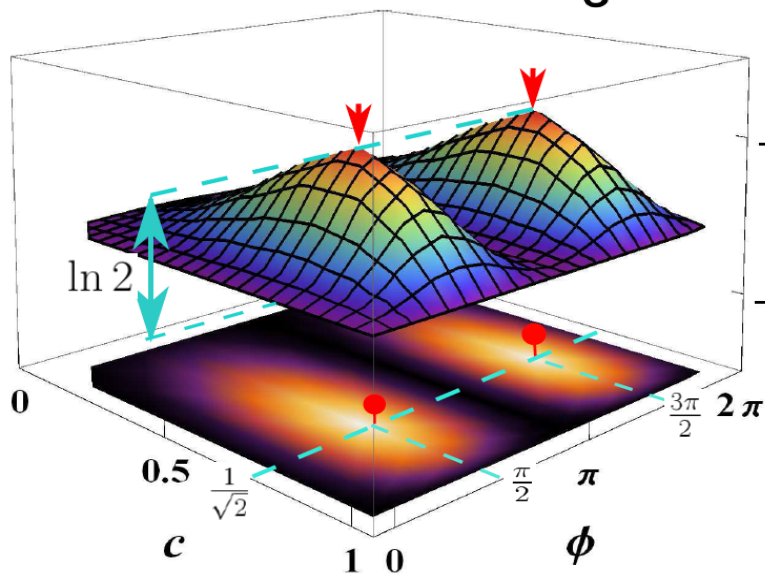
Evolving to uniform SL from VBS on kagome



Four fold ground state degeneracy and entanglement spectra

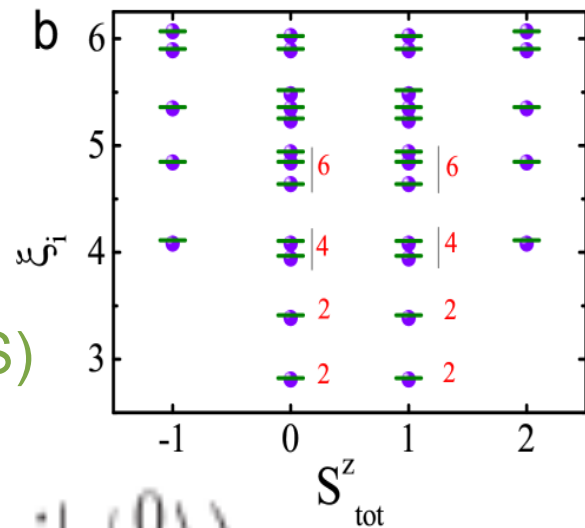


c

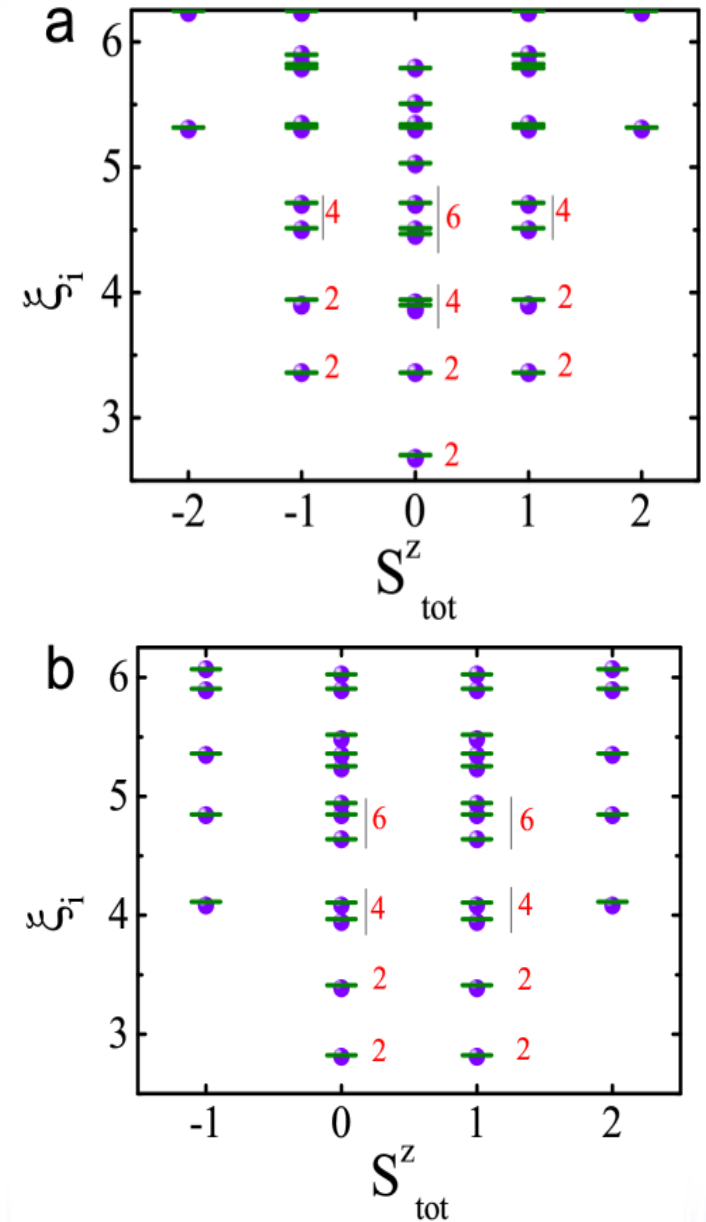
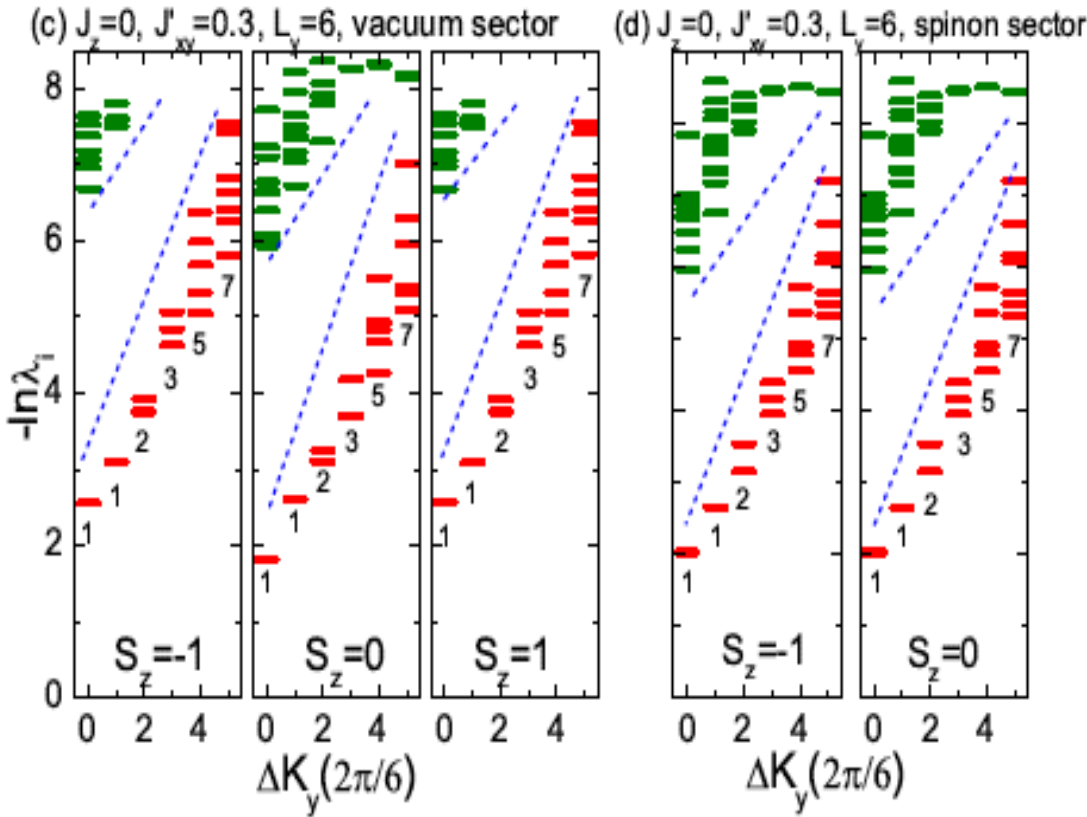


Minimum Entangled State (MES)

$$\frac{1}{\sqrt{2}}(|\psi_1^0\rangle \pm i|\psi_2^0\rangle).$$

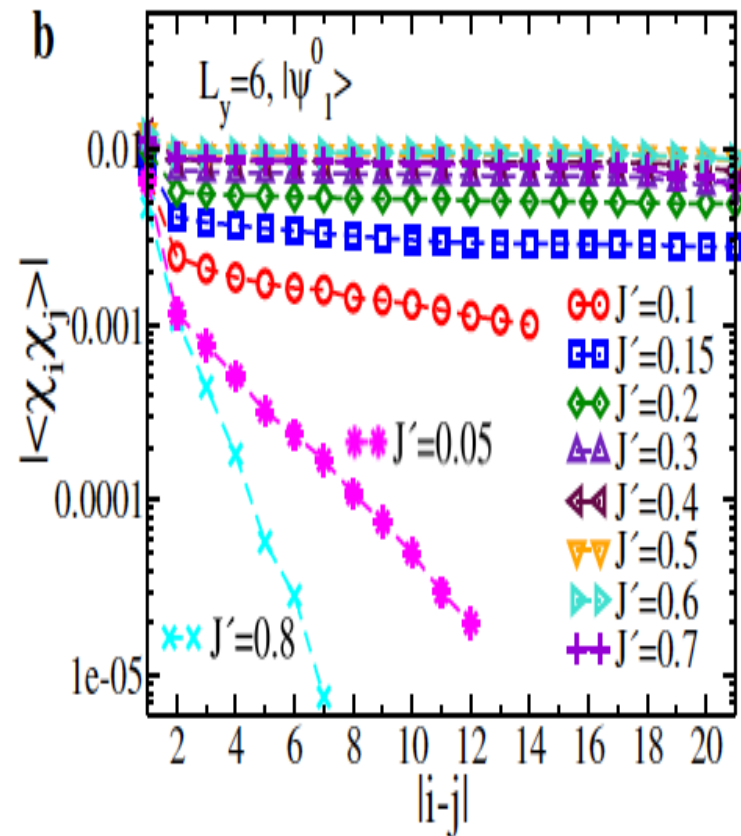
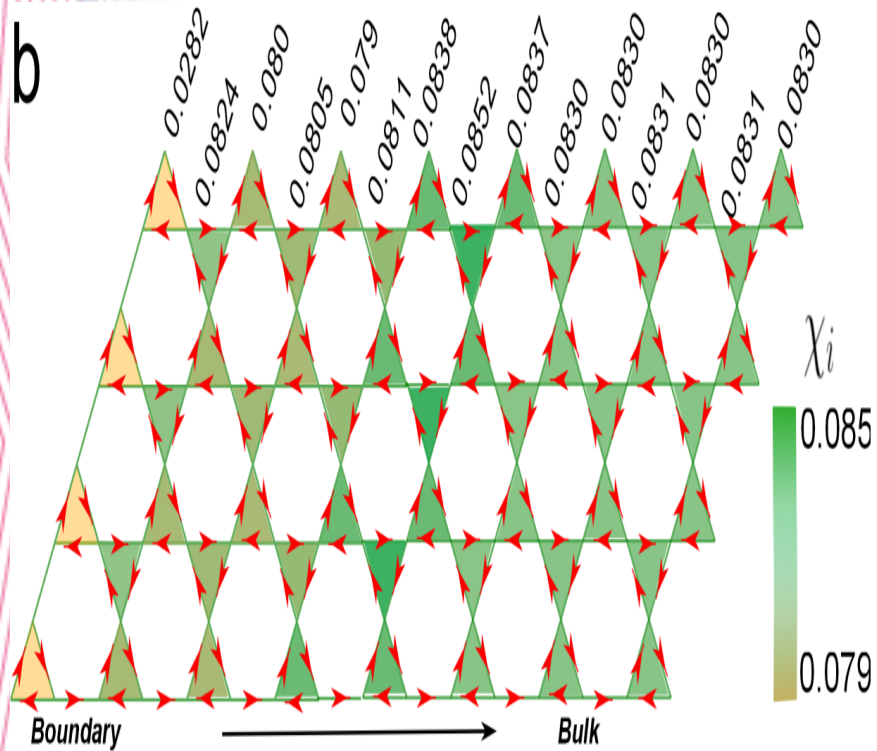


Entanglement spectra of CSL: emerging Laughlin $\nu=1/2$ FQHE

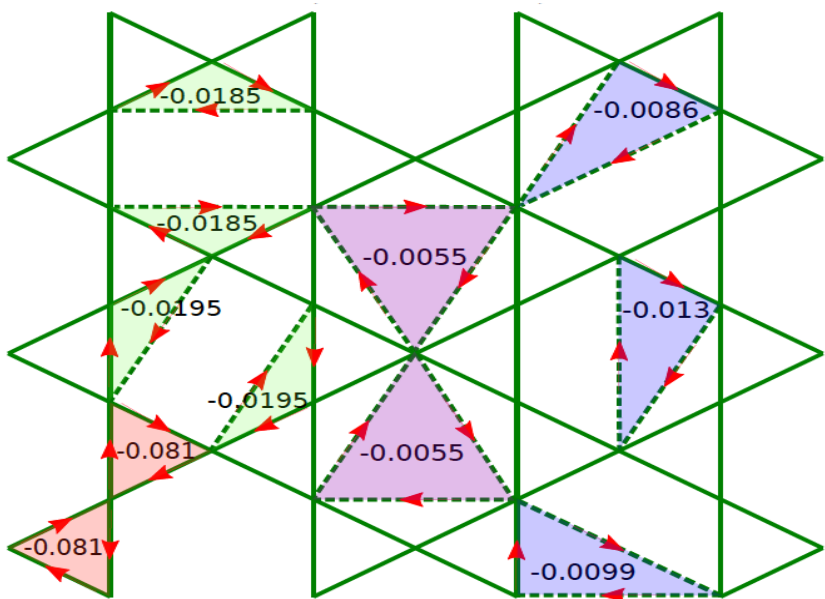
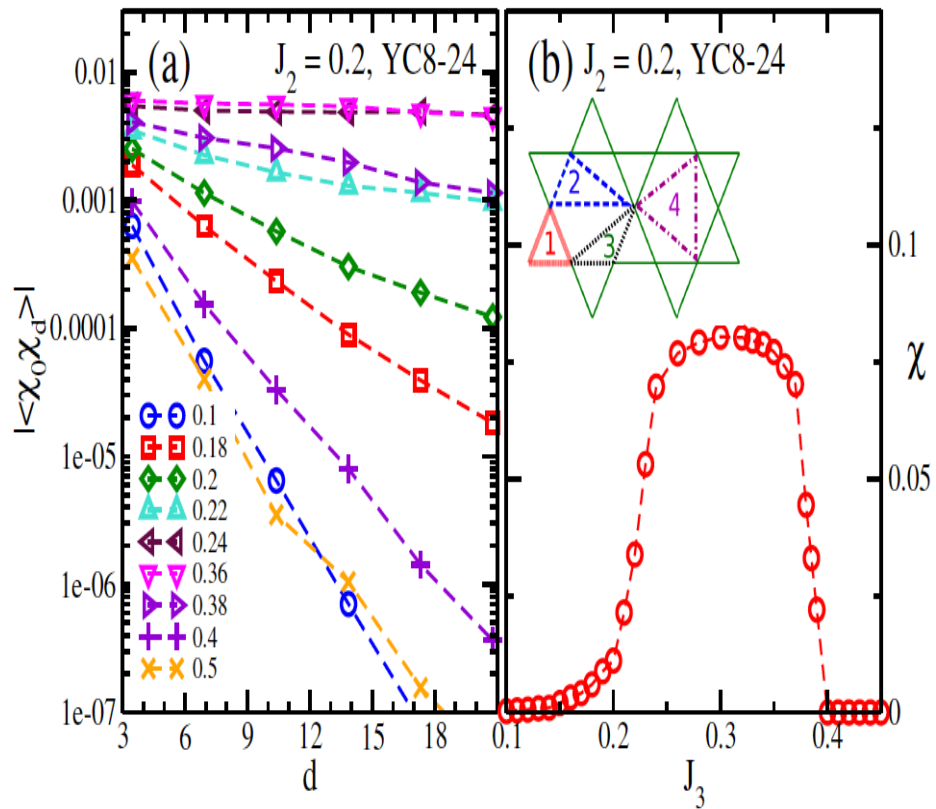
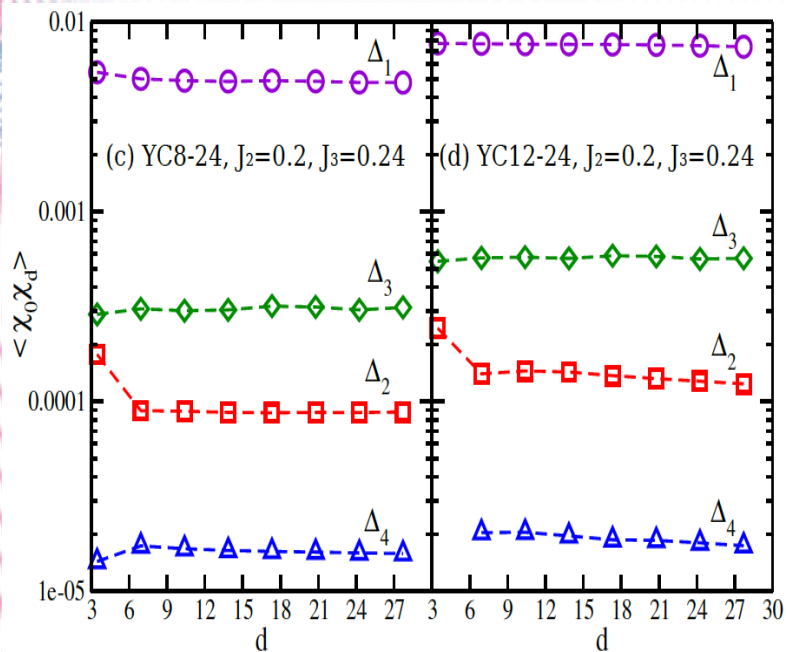


Inserting flux turns the ground state from vacuum to spinon sector

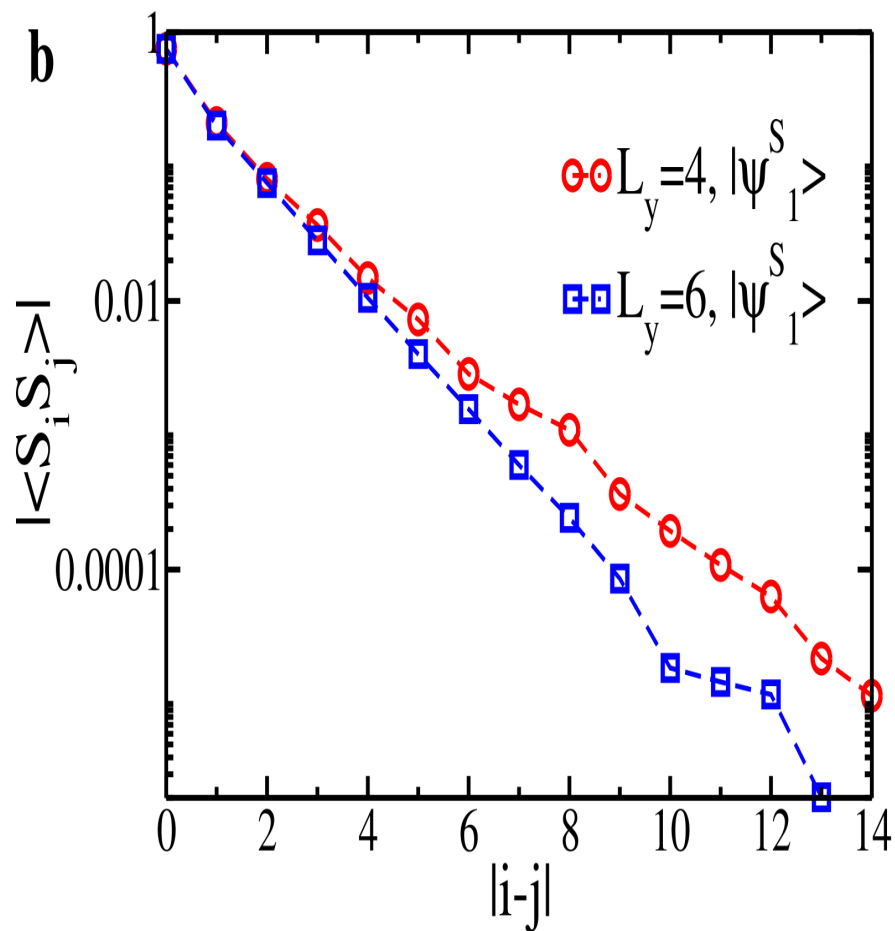
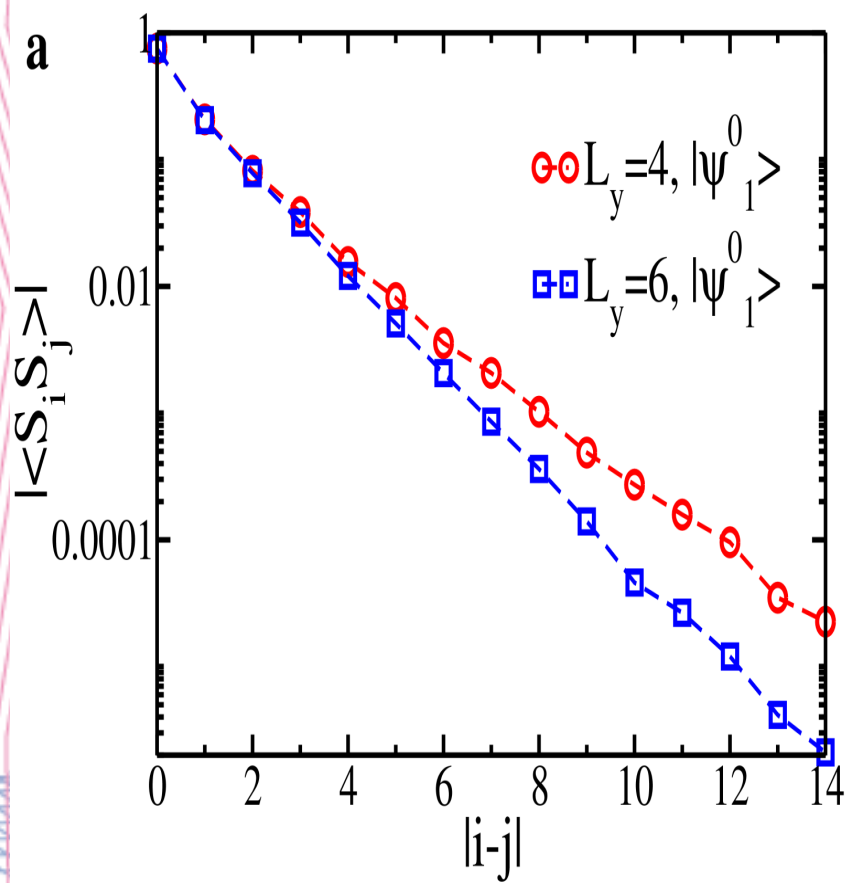
Chiral SL with nonzero spin chirality $\chi_i = (\mathbf{S}_1 \times \mathbf{S}_2) \cdot \mathbf{S}_3$



Chiral SL with nonzero spin chirality $\chi_i = (S_1 \times S_2) \cdot S_3$



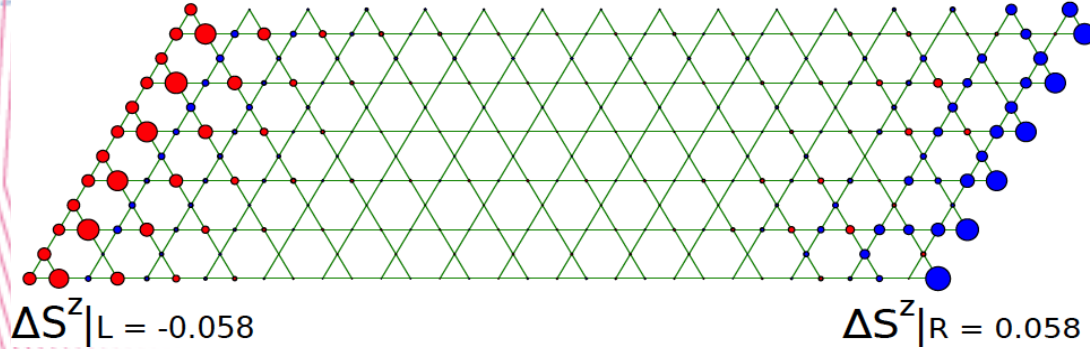
Exponential decay of the spin correlations— $J'=0.5$



Inserting flux adiabatically, spin pump, and $C=1/2$ quantization

Signature of $\nu=1/2$ Laughlin state and CSL

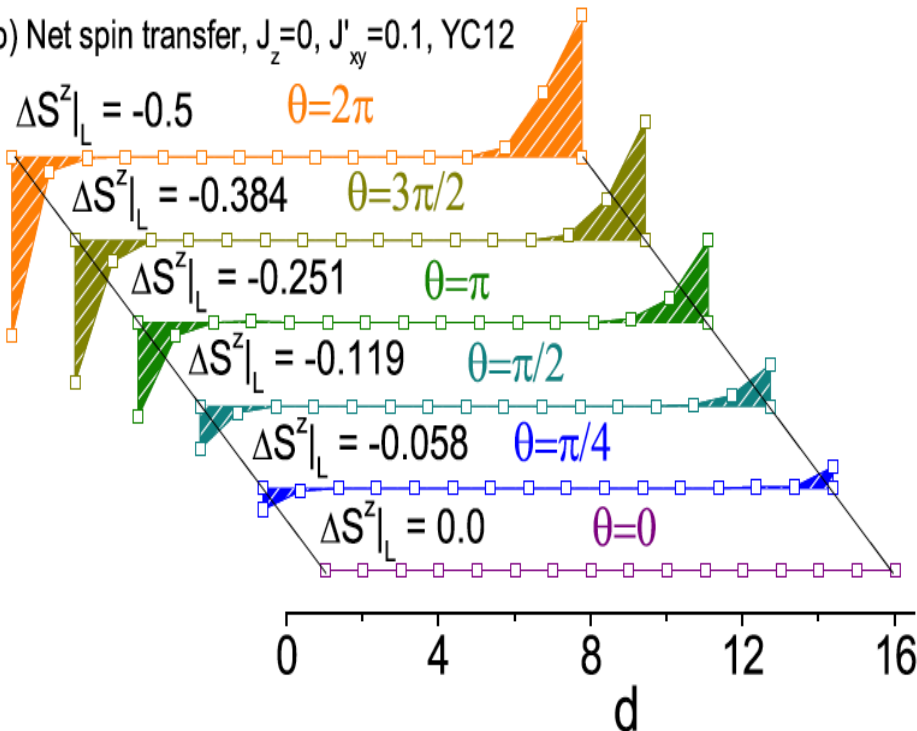
(a) $J_z=0, J'_{xy}=0.1, \text{YC12}, \theta = \pi/4$



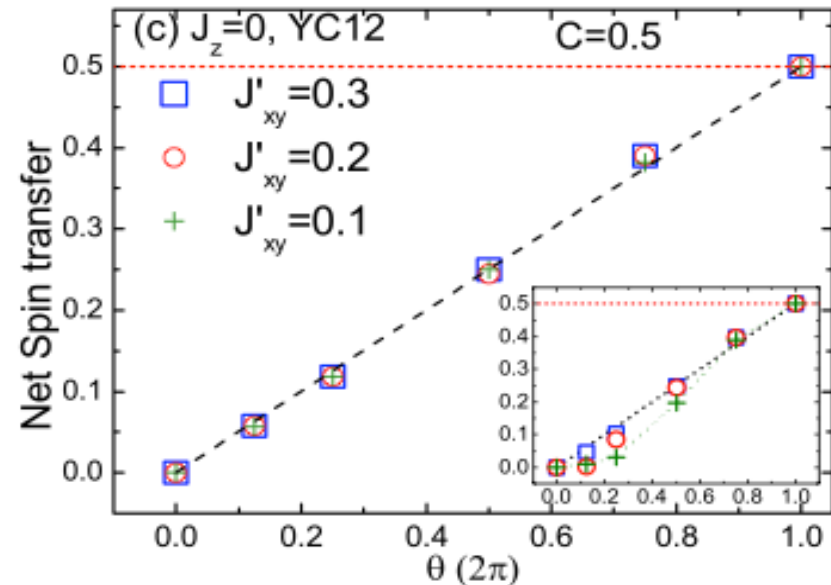
Flux=twist boundary
Phase $S^+_i S^-_j \exp(i\theta)$

Haldane and Arovas 1995

(b) Net spin transfer, $J_z=0, J'_{xy}=0.1, \text{YC12}$



$L_y=12$



Entanglement spectrum and modular matrix: matching to 1/2 FQHE (model independent)

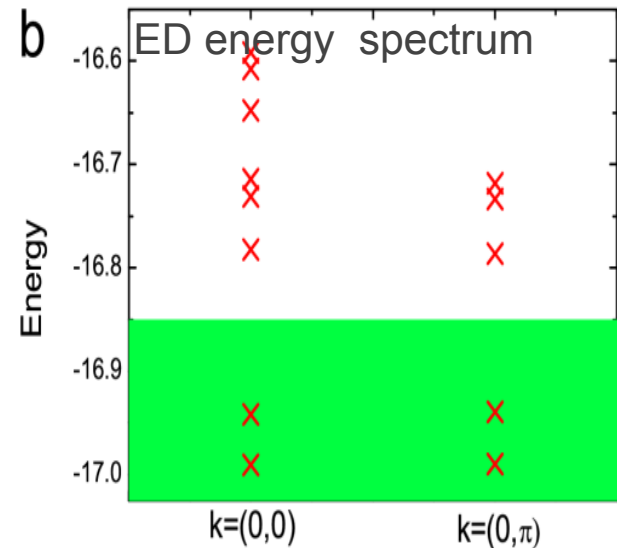
DMRG “torus” modular matrix

$$\mathcal{S} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} + \frac{10^{-2}}{\sqrt{2}} \begin{pmatrix} -0.42 & -2.2 \\ -1.26 & 0.76 - 0.15i \end{pmatrix}$$

and

$$U = e^{-i(2\pi/24)} \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \times \begin{pmatrix} e^{0.011i} & 0 \\ 0 & e^{-0.006i} \end{pmatrix}$$

Semion with
topological spin $h=1/4$



Wen 1990; Li, Haldane 2008

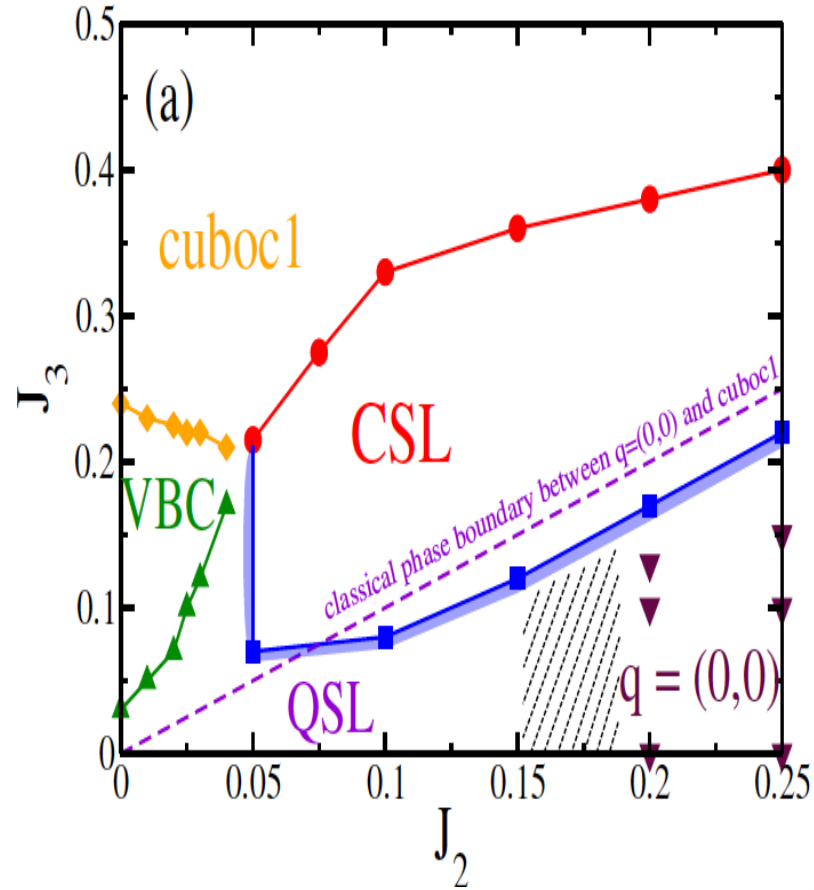
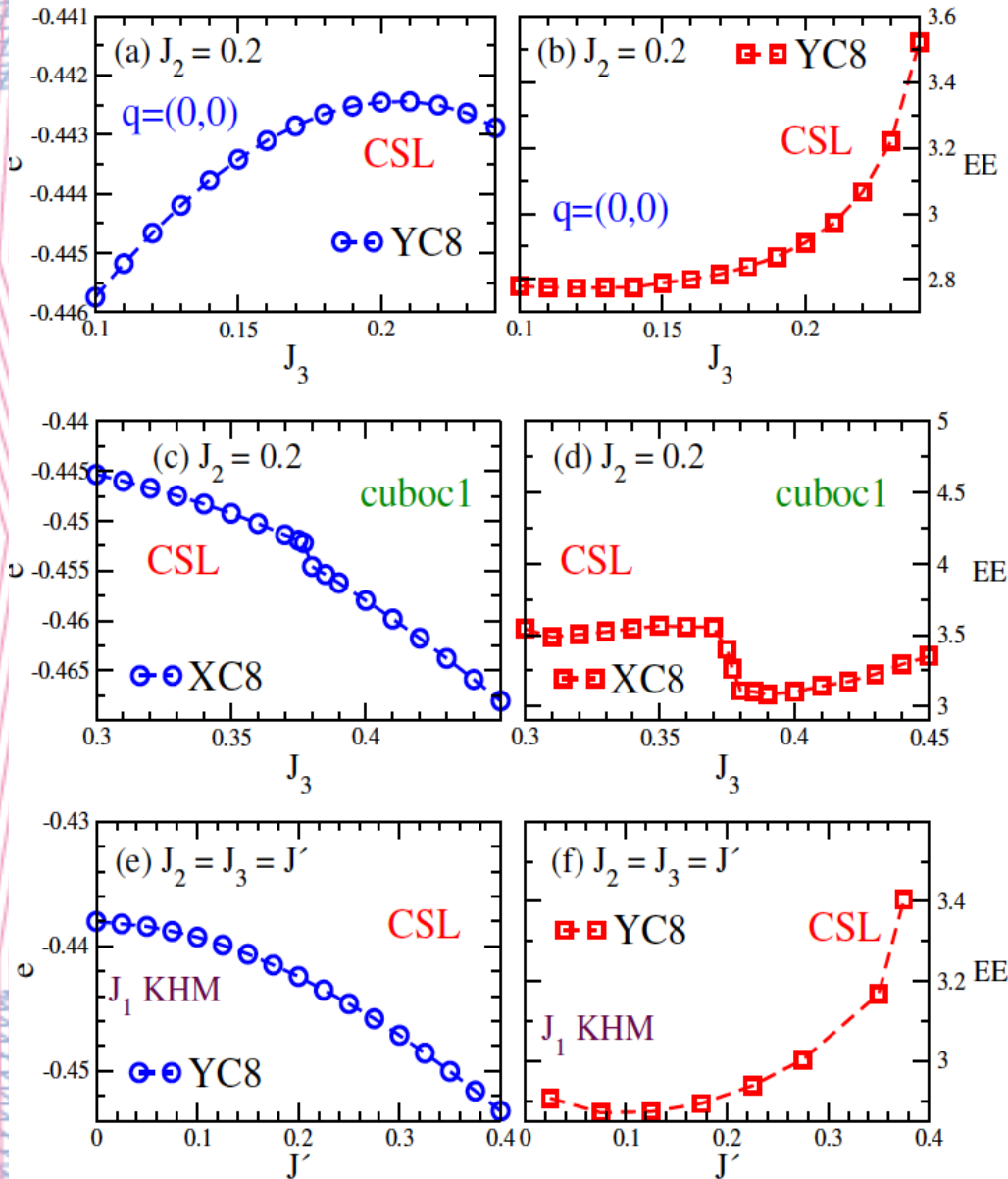
Y. Zhang et al 2012

Cincio & Vidal 2013

Zhu, Gong, Haldane, DNS 2013-2014

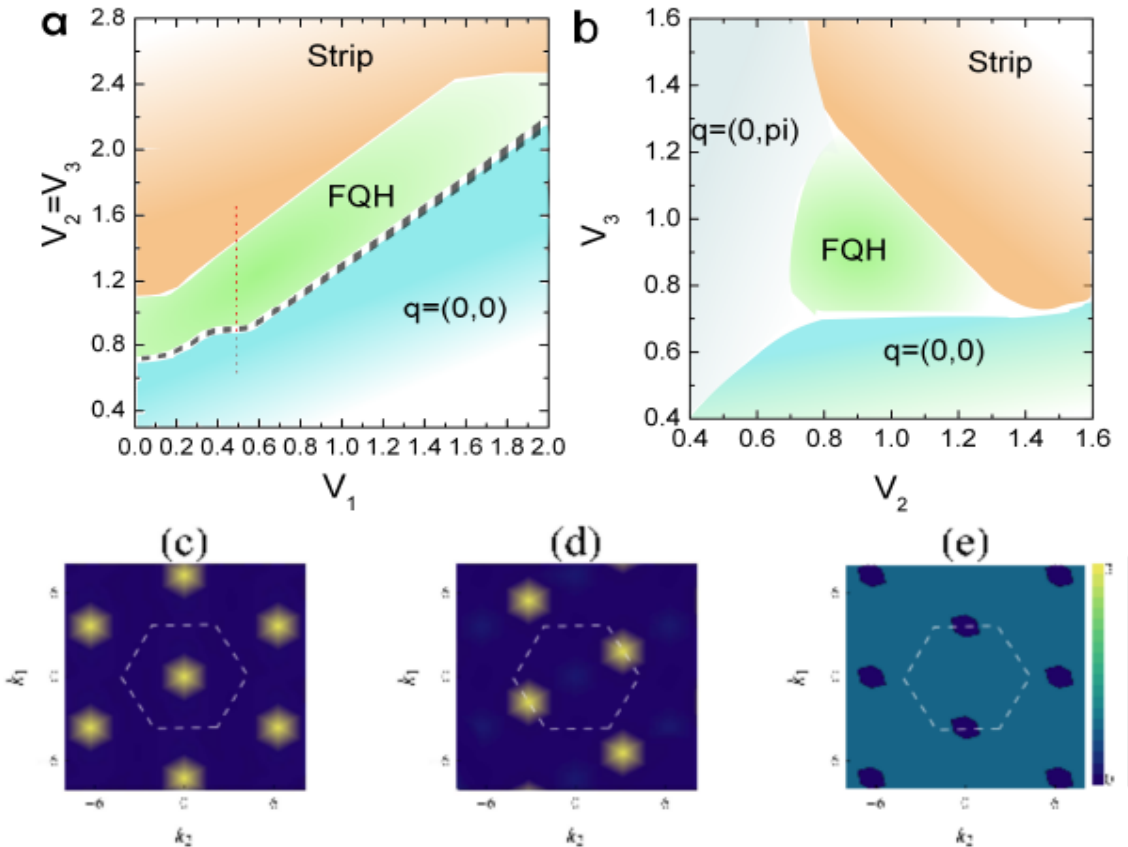
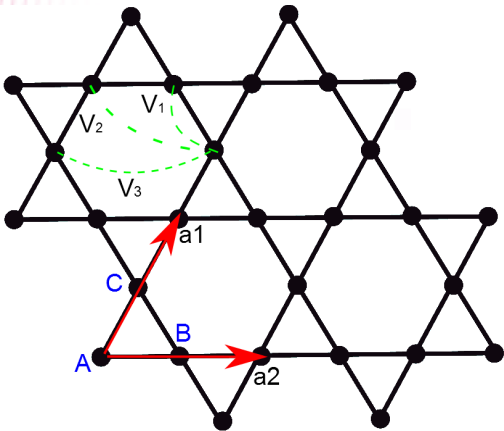


Quantum phase transitions



Boson at filling 1/3 ---- CSL (FQH)

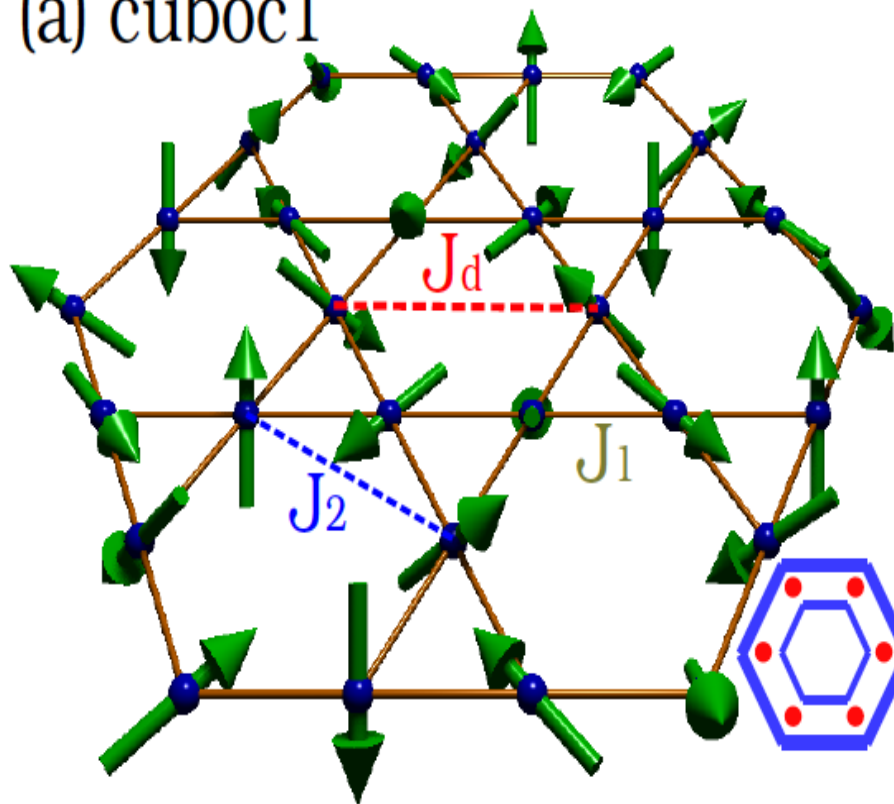
$$H = t \sum_{\langle \mathbf{r}\mathbf{r}' \rangle} [b_{\mathbf{r}'}^\dagger b_{\mathbf{r}} + \text{H.c.}] + V_1 \sum_{\langle \mathbf{r}\mathbf{r}' \rangle} n_{\mathbf{r}} n_{\mathbf{r}'} + V_2 \sum_{\langle\langle \mathbf{r}\mathbf{r}' \rangle\rangle} n_{\mathbf{r}} n_{\mathbf{r}'} + V_3 \sum_{\langle\langle\langle \mathbf{r}\mathbf{r}' \rangle\rangle\rangle} n_{\mathbf{r}} n_{\mathbf{r}'}$$



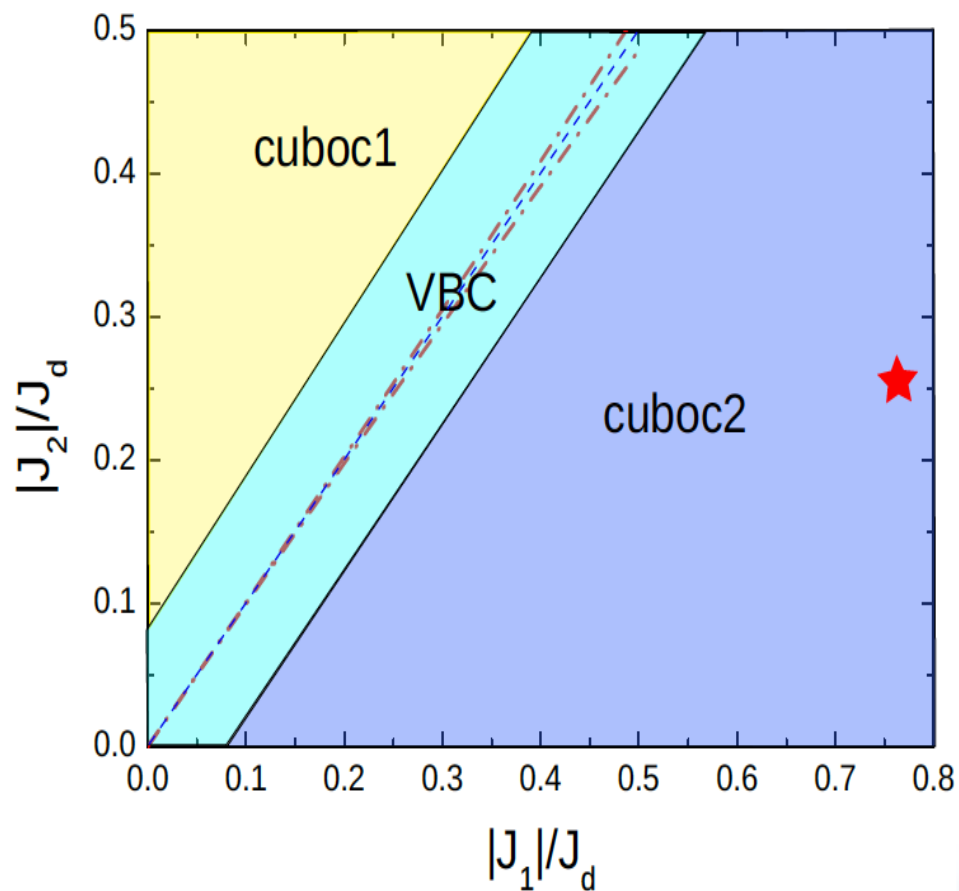
Boson at filling 1/3: W Zhu, et. al, PRL 2016

No SL nearby two cuboc phases (related to Kapellasite), disagree with slave-particle

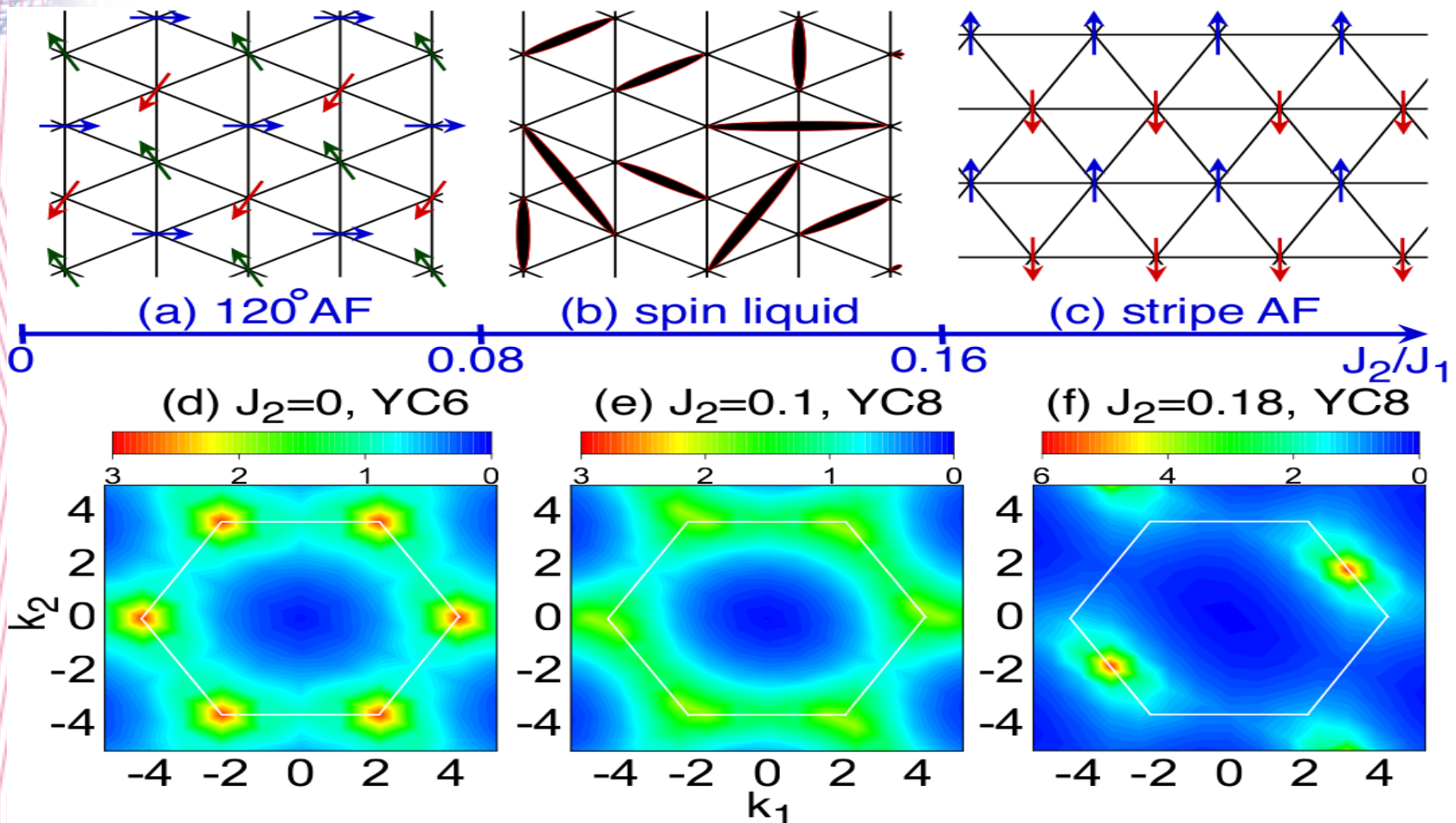
(a) cuboc1



Gong et al 2016



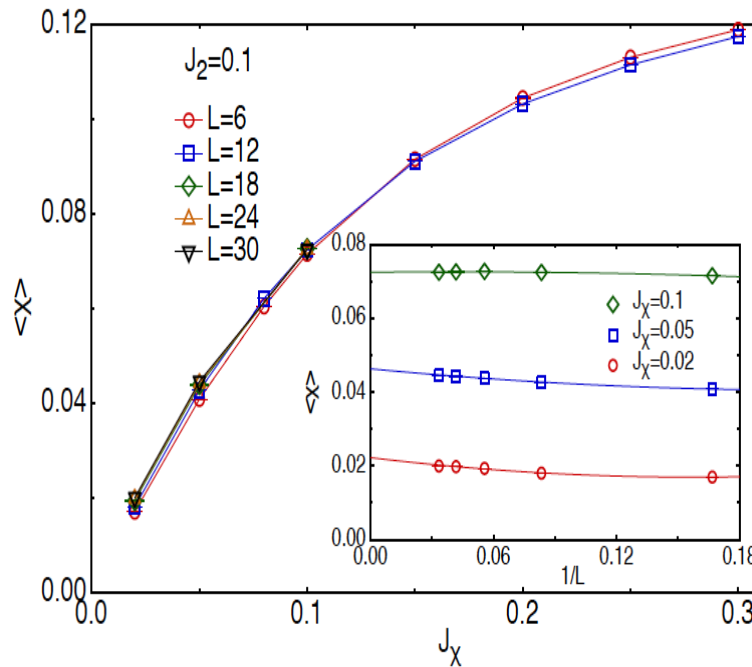
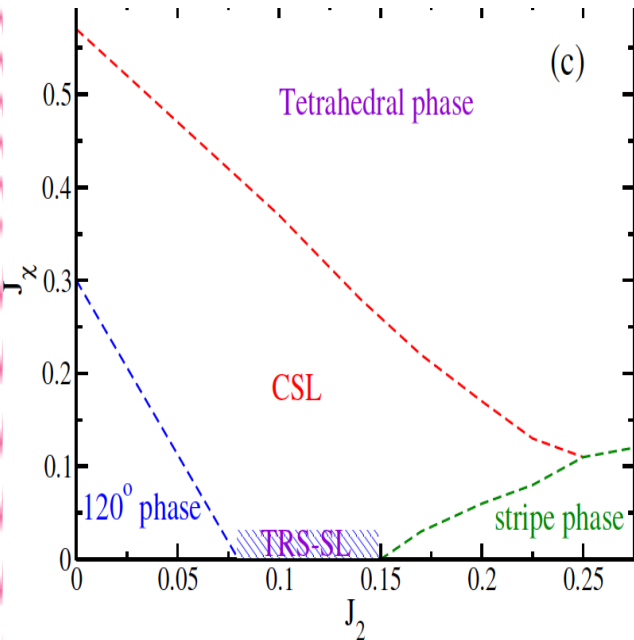
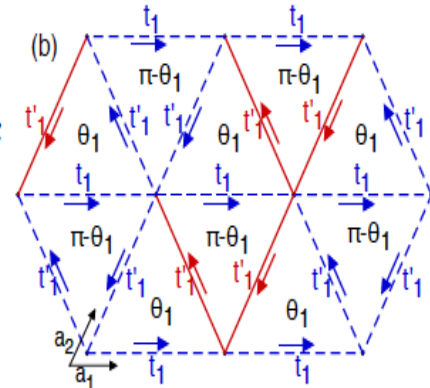
Frustrated triangular lattice spin $\frac{1}{2}$ J1-J2 models



Z. Zhu and S. R. White, PRB; W. J. Hu et al., PRB

DMRG and VMC pictures: Dirac to CSL

$$H = J_1 \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j + J_2 \sum_{\langle\langle i,j \rangle\rangle} \vec{S}_i \cdot \vec{S}_j + J_\chi \sum_{\Delta/\nabla} (\vec{S}_i \times \vec{S}_j) \cdot \vec{S}_k$$



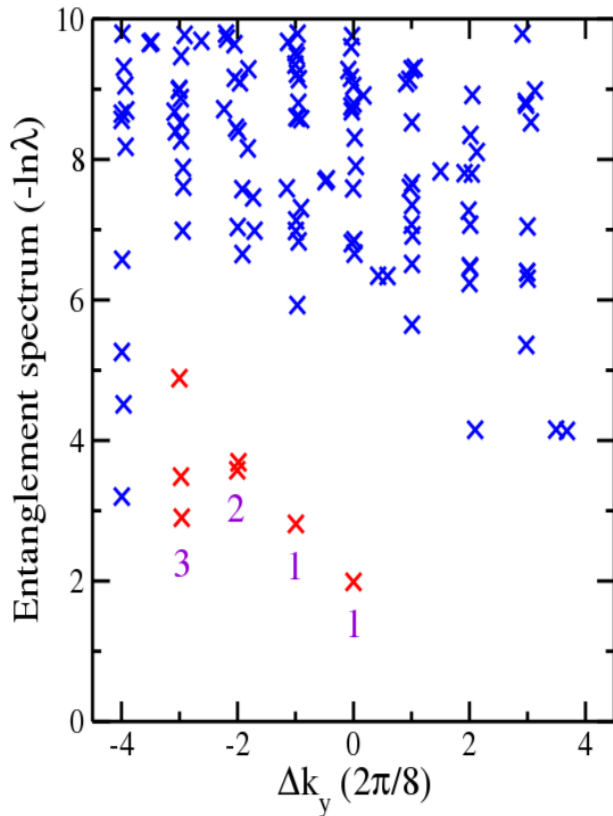
Agree with ED phase diagram, Lauchli et al.

Bauer et al. (2014)
 He et al. (2016)
 Zhu et al. 2015
 Hu et al. 2015
 Similar critical phase
 For kagome

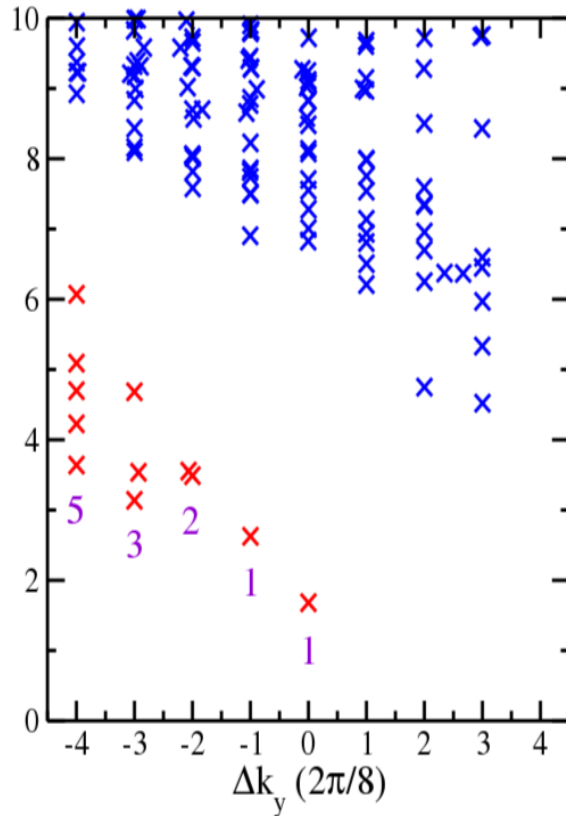
Entanglement spectrum of the vacuum sector and additional J_x $(S1 \times S2) \times S3$

$J_1=1.0, J_2=0.1, N=8 \times 20$, vacuum sector

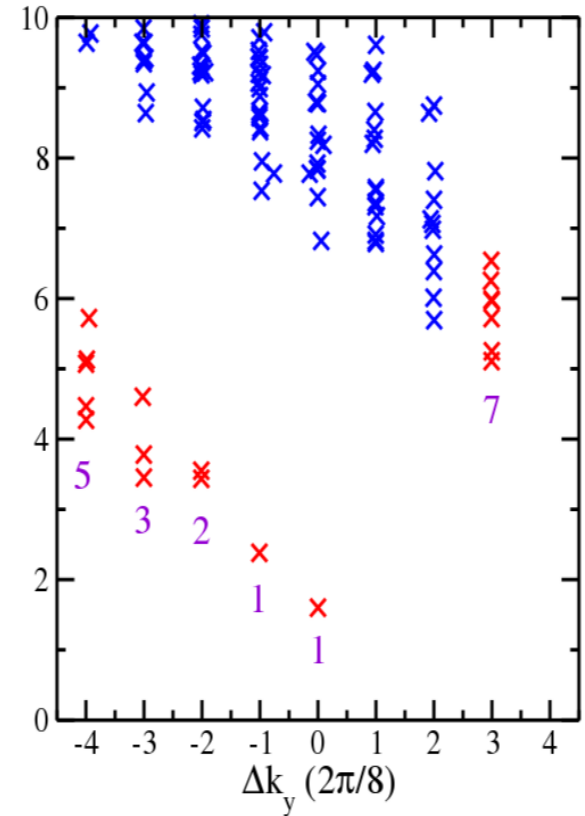
(a) $J_\chi = 0.0$



(b) $J_\chi = 0.05$

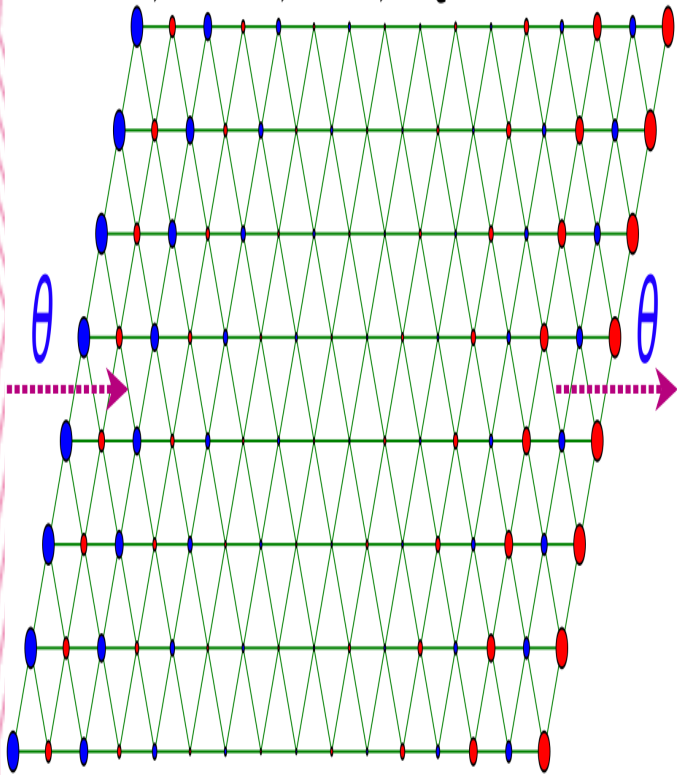


(c) $J_\chi = 0.2$

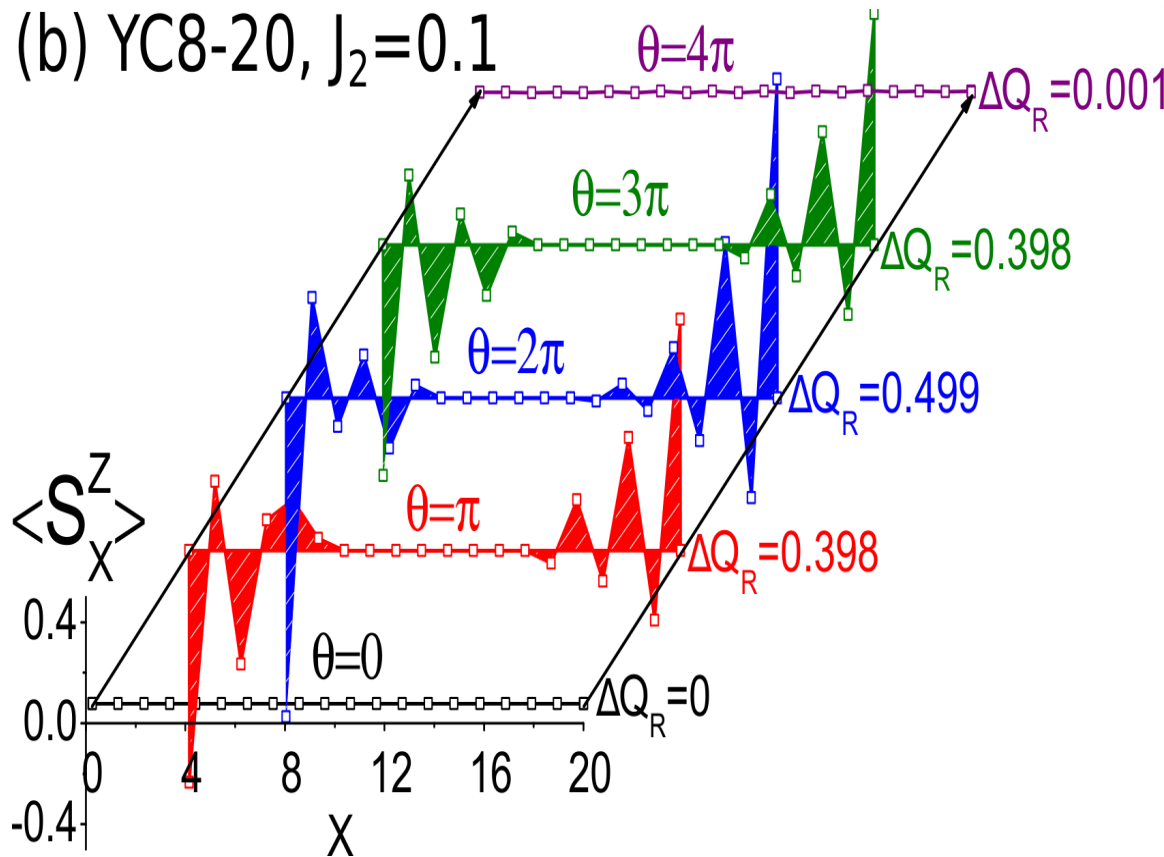


SL phase around $J_2=0.1$, inserting flux 2π gives fermionic spinon sector

(a) $J_2=0.1$, YC8-16, $\theta=2\pi$, $\Delta Q_R=0.48$

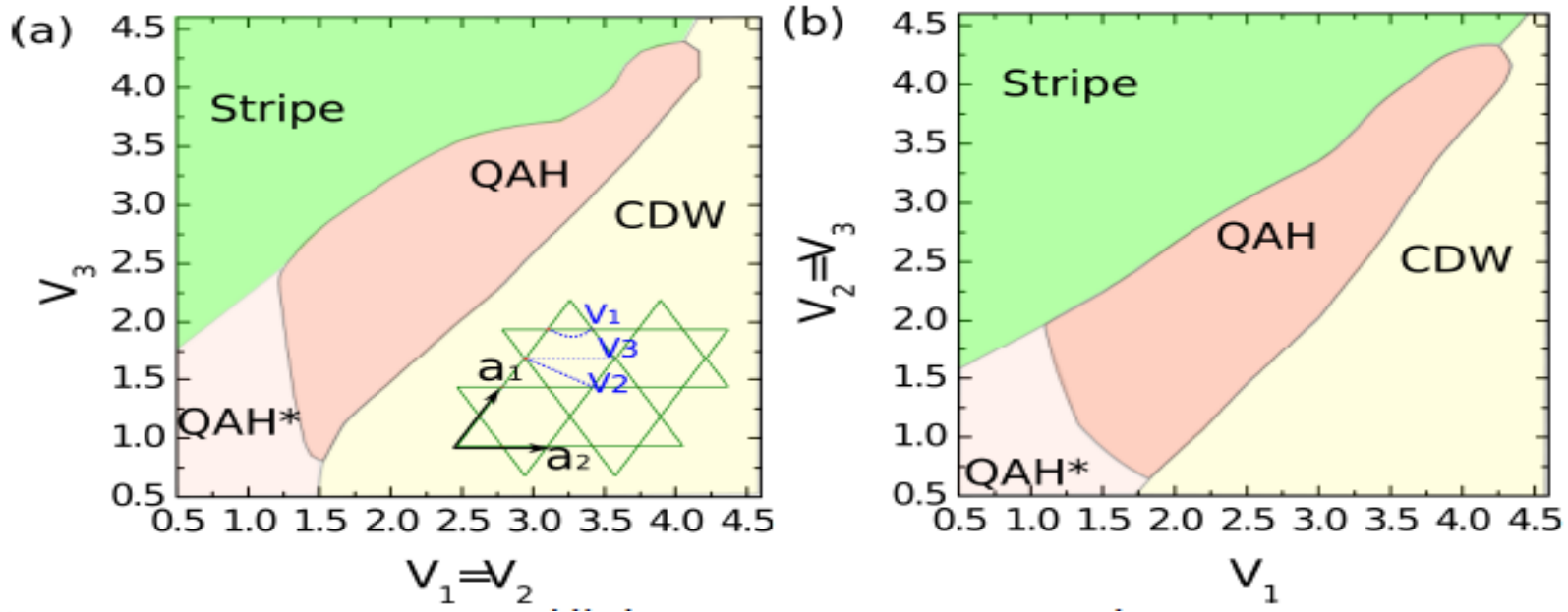


(b) YC8-20, $J_2=0.1$



Interaction-driven QAH ---- Fermion-Hubbard on kagome at 1/3

$$H = t \sum_{\langle rr' \rangle} [c_r^\dagger c_{r'} + \text{H.c.}] + V_1 \sum_{\langle rr' \rangle} n_r n_{r'} + V_2 \sum_{\langle\langle rr' \rangle\rangle} n_r n_{r'} + V_3 \sum_{\langle\langle\langle rr' \rangle\rangle\rangle} n_r n_{r'}$$

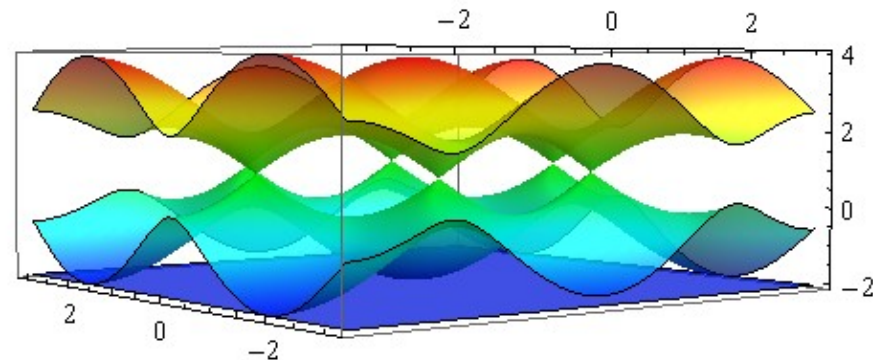
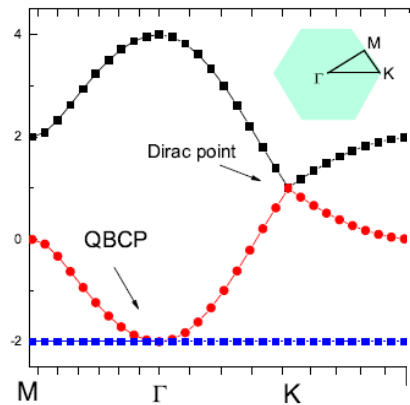
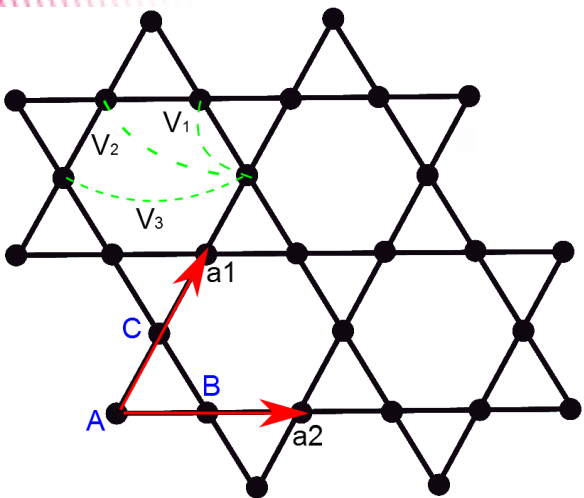


How to identify QAH? Using ED and DMRG

- Time-Reversal Symmetry Spontaneously Breaking
Emergent Staggered Magnetic Flux
- Topological Chern Number (Hall Conductance)

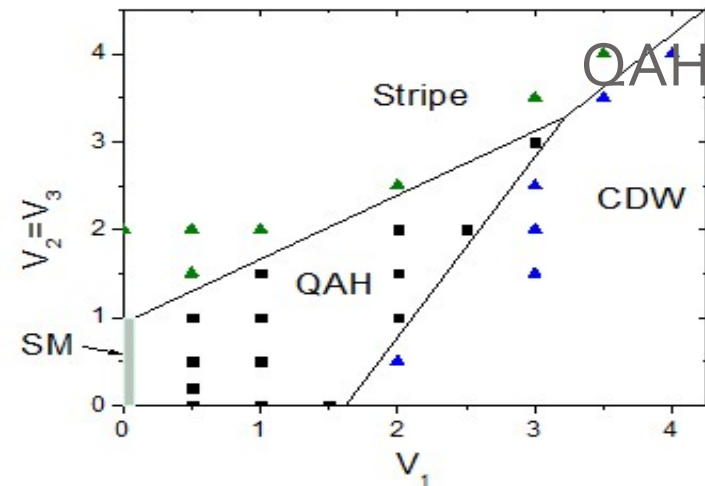
Interaction-driven QAH ---- Fermion-Hubbard on kagome at 1/3

$$H = t \sum_{\langle rr' \rangle} [c_{r'}^\dagger c_r + \text{H.c.}] + V_1 \sum_{\langle rr' \rangle} n_r n_{r'} + V_2 \sum_{\langle\langle rr' \rangle\rangle} n_r n_{r'} + V_3 \sum_{\langle\langle\langle rr' \rangle\rangle\rangle} n_r n_{r'}$$



Mean-Field
Phase Diagram

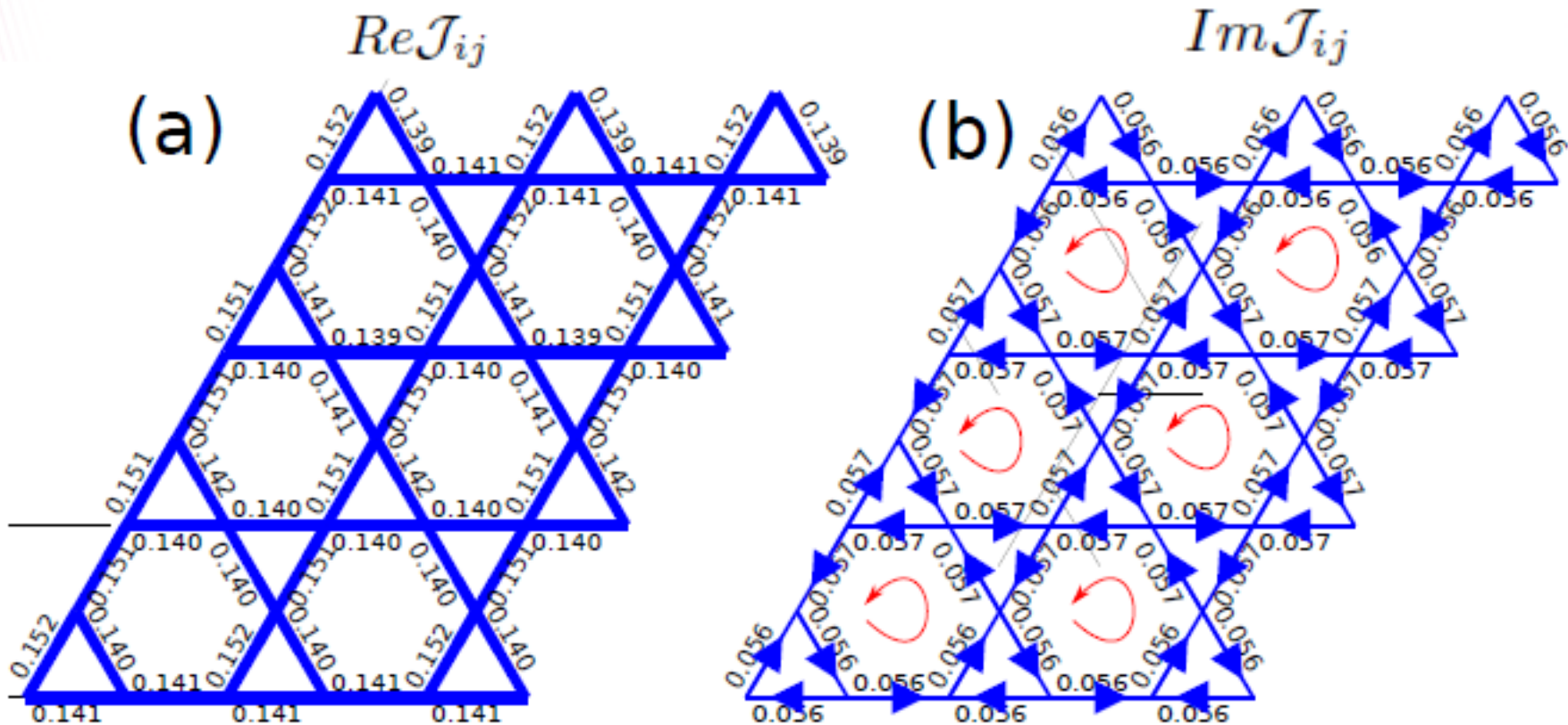
Wu et al. (2007), Bermann et al (08)
Sun et al (09)
Zhu, Gong, Zeng, Fu, Sheng (16)



Interaction-driven QAH ---- Emergent Magnetic Flux

$$H = t \sum_{\langle rr' \rangle} [c_r^\dagger c_{r'} + \text{H.c.}] + V_1 \sum_{\langle rr' \rangle} n_r n_{r'} + V_2 \sum_{\langle\langle rr' \rangle\rangle} n_r n_{r'} + V_3 \sum_{\langle\langle\langle rr' \rangle\rangle\rangle} n_r n_{r'}$$

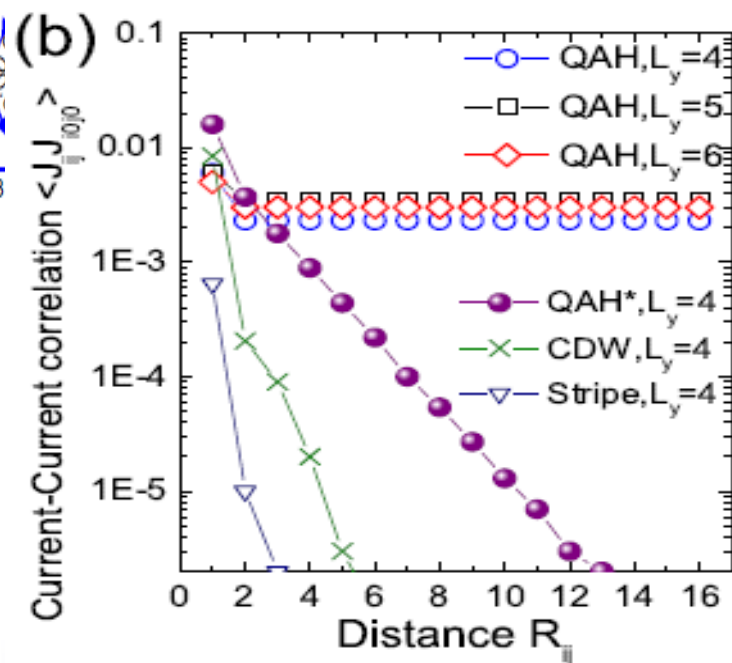
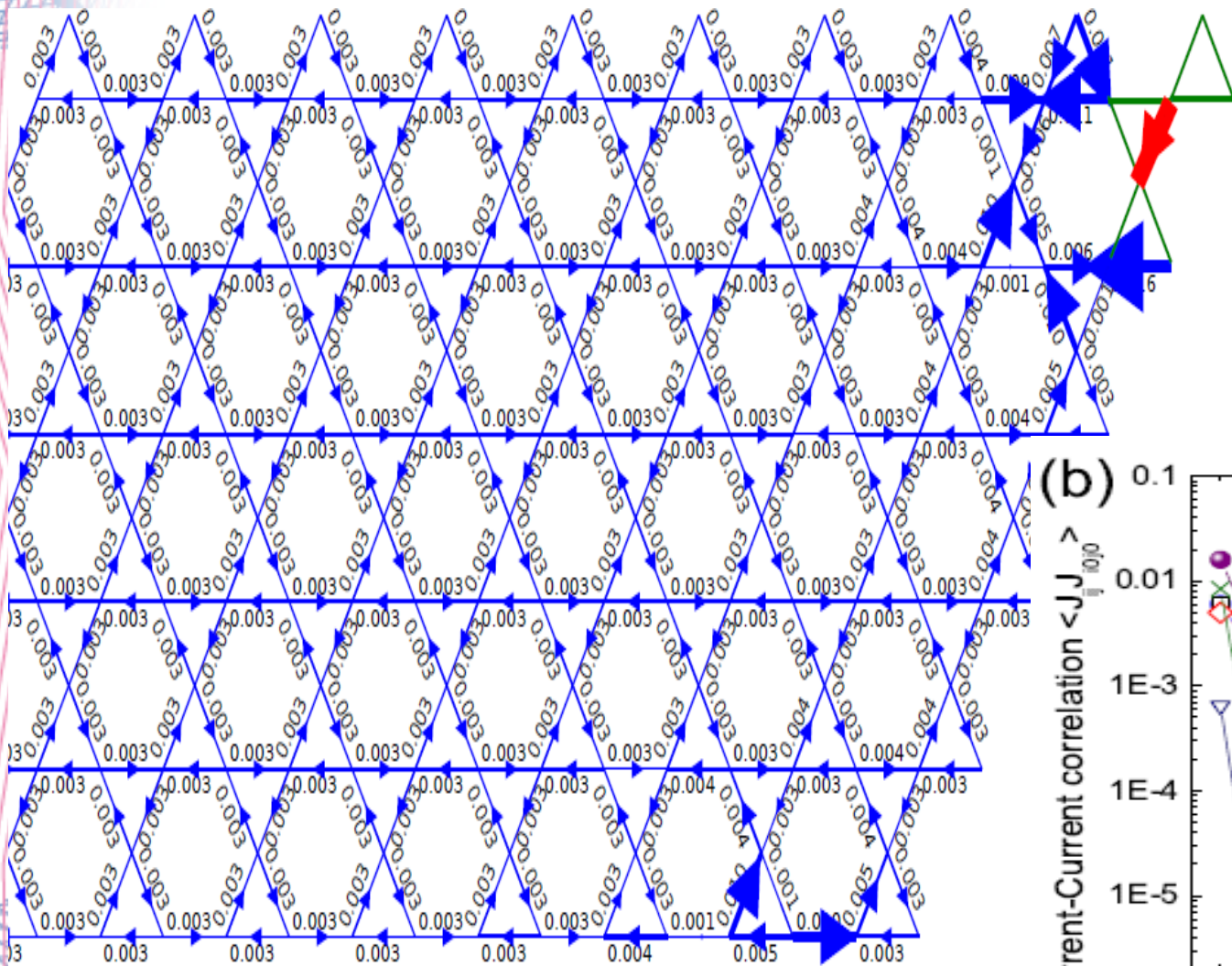
$$\mathcal{J}_{ij} = \langle \Psi^L | c_i^\dagger c_j | \Psi^L \rangle$$



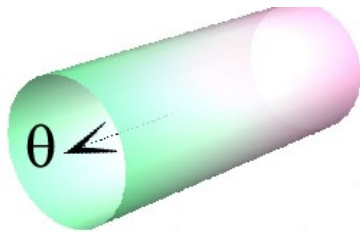
Staggered Magnetic Flux!!

Current-Current correlation

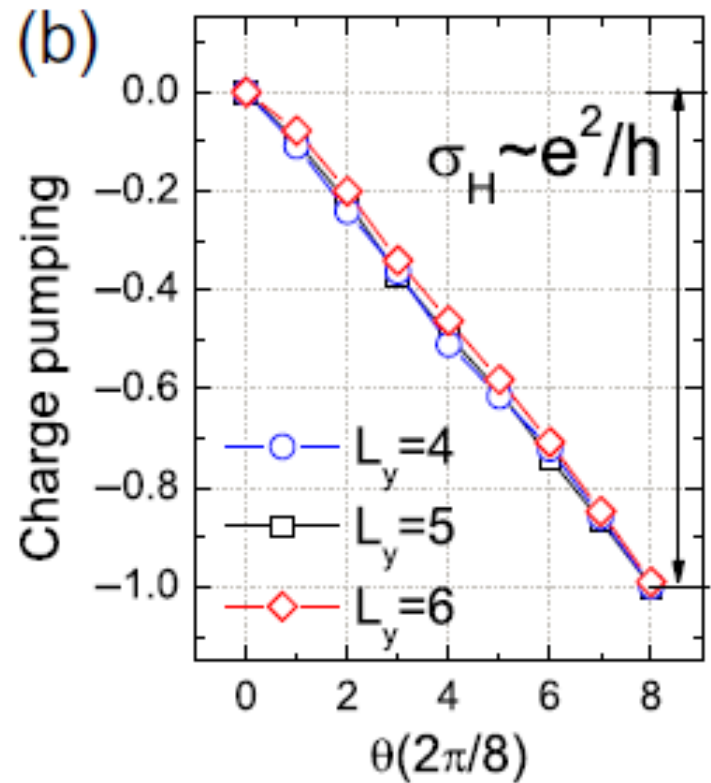
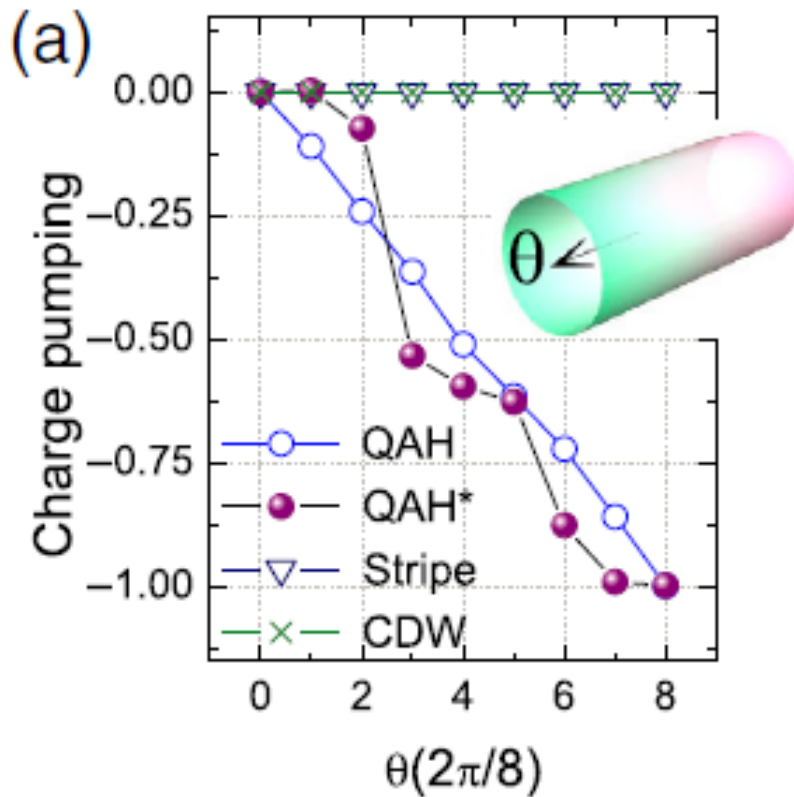
$$J_{ij} = i \langle \Psi^{L(R)} | c_i^\dagger c_j - c_j^\dagger c_i | \Psi^{L(R)} \rangle$$



Interaction-driven QAH ---- Chern Number



Chern Number (Hall conductance):
 $C = 1$ (Left chirality)
 $C = -1$ (Right chirality)



Summary and discussions

Most experimental SL candidates have triangular or kagome lattices.

Numerically we also found frustrated triangles in these lattices play essential role in stabilizing SLs.

We discovered long-sought chiral spin liquid (CSL) in kagome lattice model and establish the instability toward CSL for triangular J1-J2 model.

VMC/DMRG studies indicate a continuous Dirac SL to CSL transition

Fermionic systems with such frustrations can realize Spontaneous Quantum Hall Effect and other exotic phases (ongoing work).

Research collaborations

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Tiansheng Zeng (CSUN), Wenjun Hu (RU)

Y. C. He and Yan Chen, Balents on CSL;
Liang Fu on Spontaneous QHE

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Thank you !