



Kavli ITS workshop on “topological matter & quantum computation”

Majorana zero mode inside vortex of topological superconductors

Jinfeng Jia

School of Phys. & Astronomy, Shanghai Jiao
Tong University, Shanghai 200240, China

Email: jfjia@sjtu.edu.cn

Beijing, May 4-6, 2018

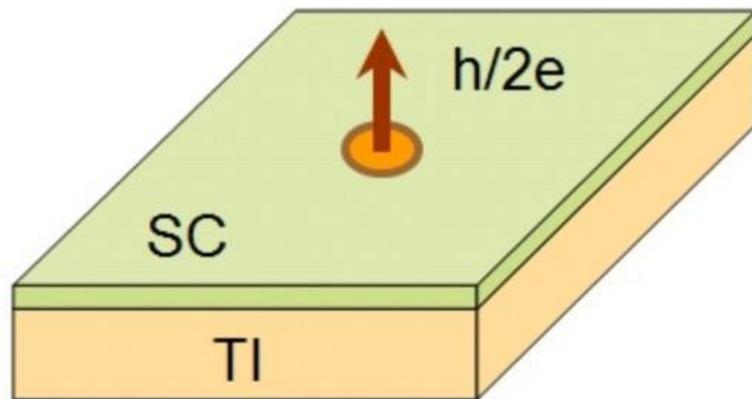
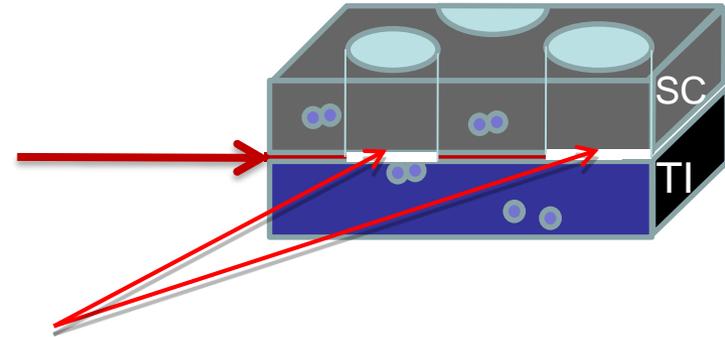


Content

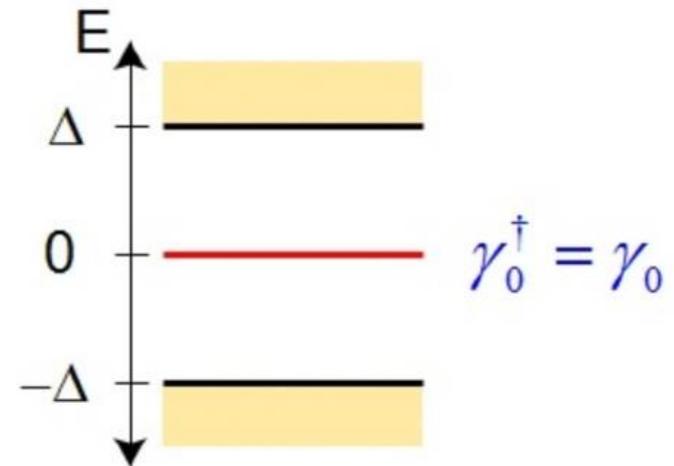
- ① **Introduction**
 - ① **Topological superconductor by proximity effects**
 - ① **Majorana fermion in the vortex**
 - ① **Summary**
-

TSC by Proximity effect

- Proximity effect between SC and TI leads to „ $p_x + ip_y$ “ SC-like-state
- Majorana Bound States (MBS) at magnetic vertices



Quasiparticle Bound state at $E=0$



Majorana Fermion γ_0

MBE growth to obtain sharp interface.

- Pb, Nb, Al,
- NbSe₂, NbS₂.....
-



Sharp interface!

S-SC on TI,
difficult to realize in practice.

Bi-Sb, Bi₂Te₃, Sb₂Te₃, Bi₂Se₃,

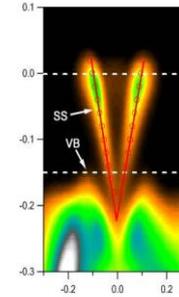
- heavy atoms
- Stable at T < 300 °C

TSS is covered by
superconductor



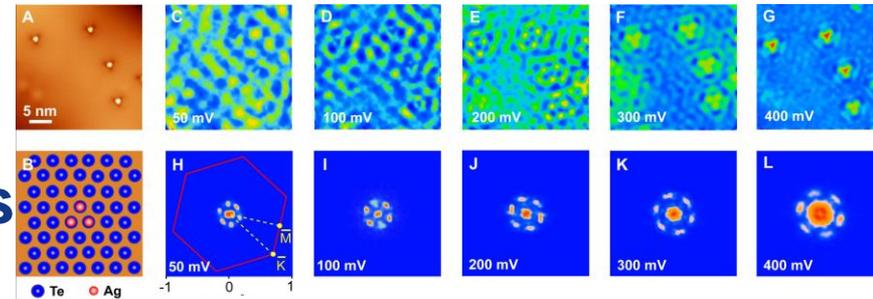
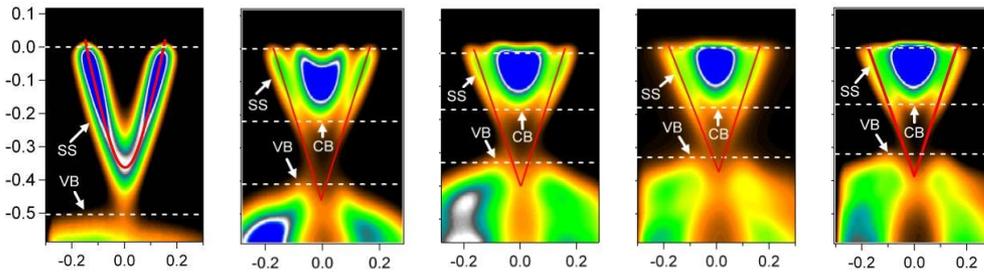
Our previous work on TI films

High quality Bi_2Te_3 , Bi_2Se_3 , Sb_2Te_3 thin films by MBE



Adv. Mater. 22, 4002-4007 (2010)
Adv. Mater. 23, 1162-1165 (2011)

Thinnest limit

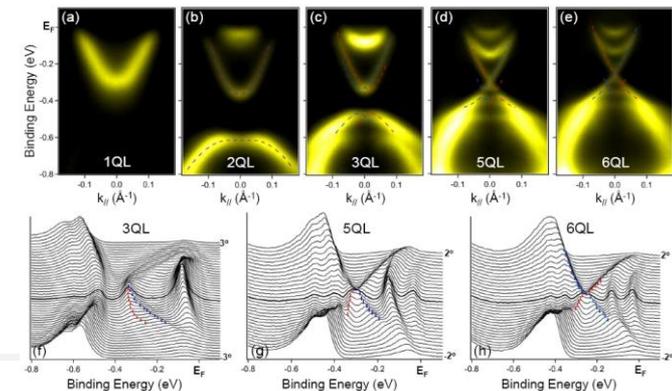


Control the type of Bi_2Te_3 films

Studies by LTSTM & ARPES

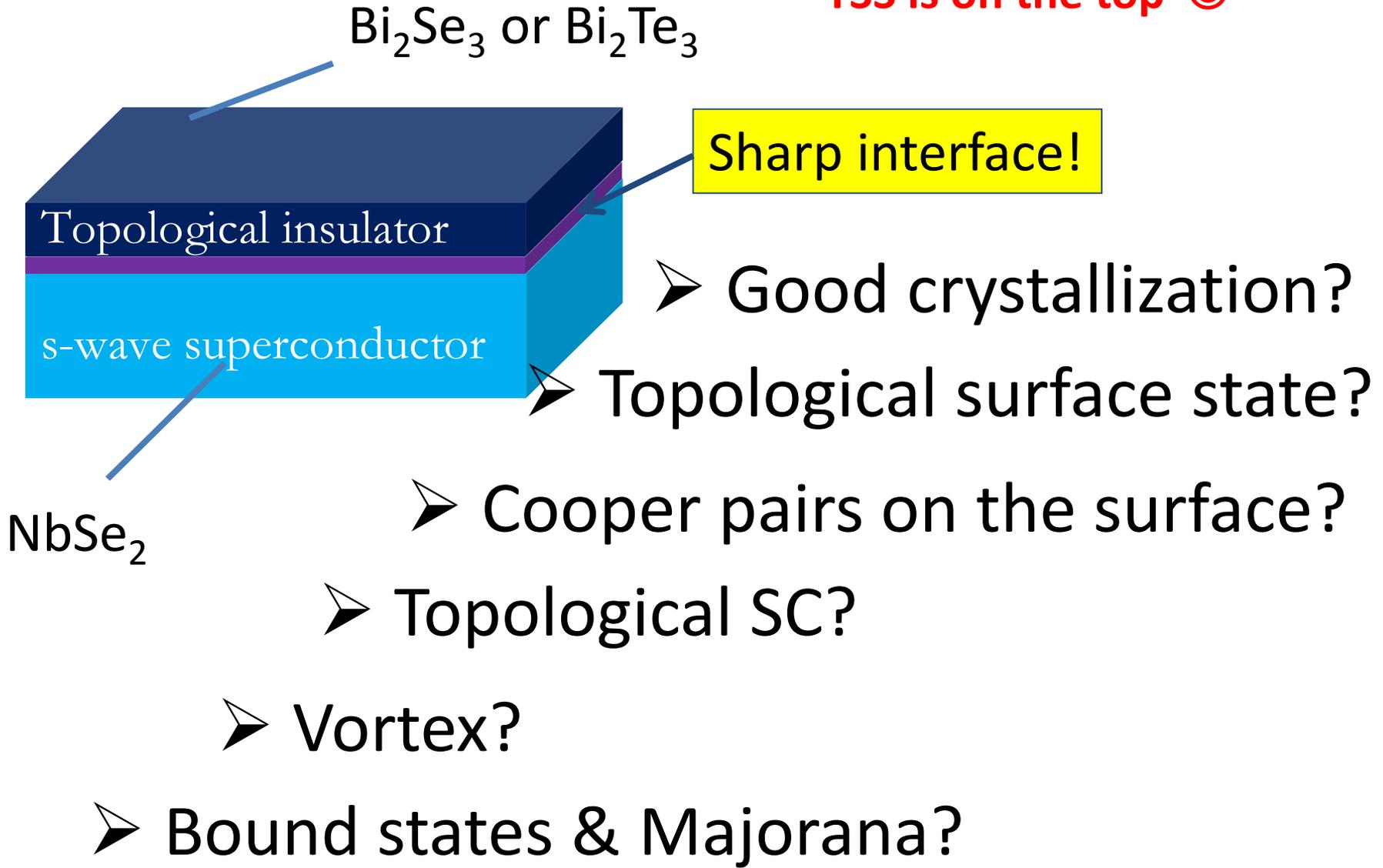
- Standing wave PRL 103, 266803 (2009)
- Landau levels PRL 105, 076801 (2010)
- Formation of DC

– Nature Phys. 6, 584 (2010)

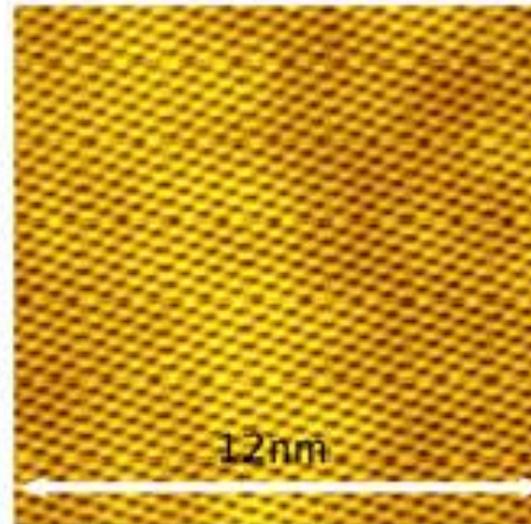
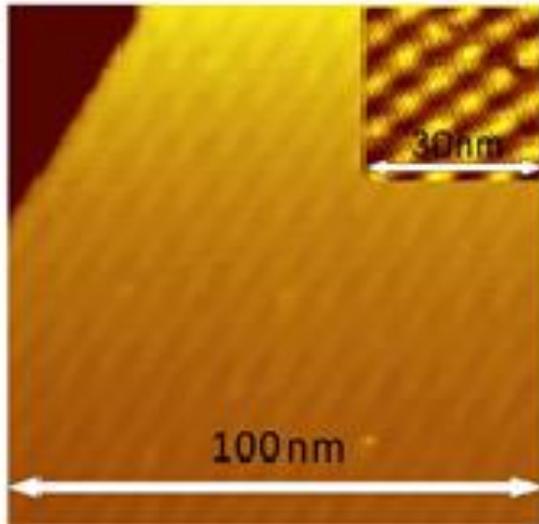
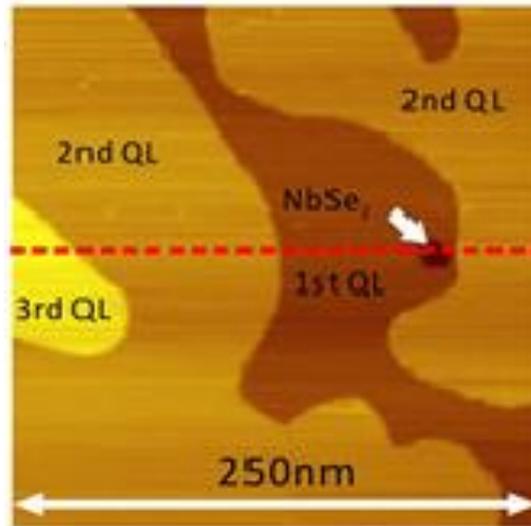
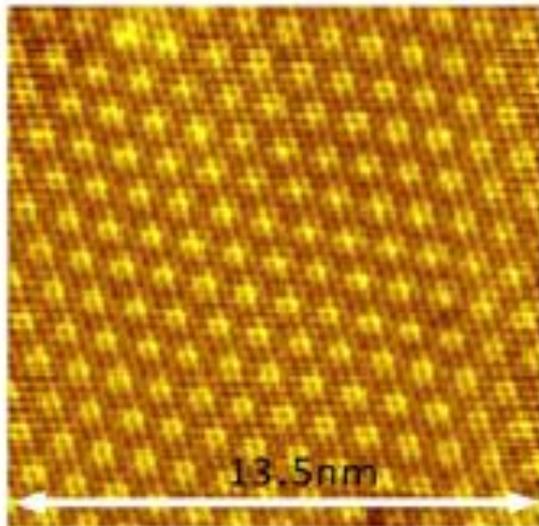


TI on SC, much easier to achieve

TSS is on the top 😊



Bi₂Se₃ films grown on NbSe₂



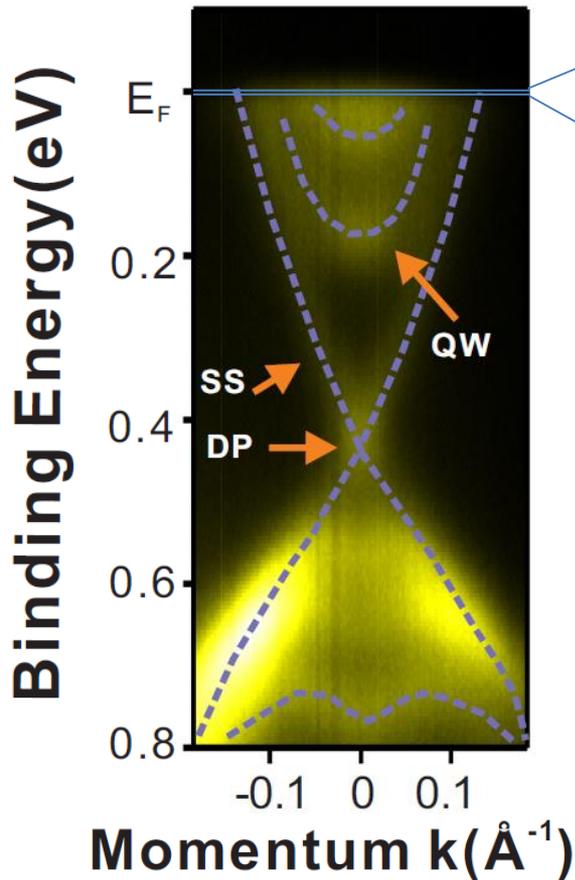
SC

**Sharp
interface**

**Less
defects**

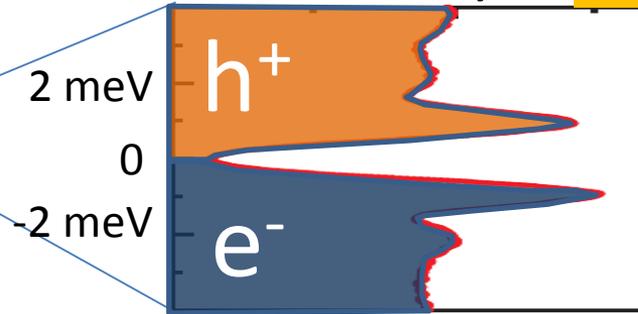
A platform for searching Majorana Fermions

Topological surface states



Intensity

superconductivity



REPORTS

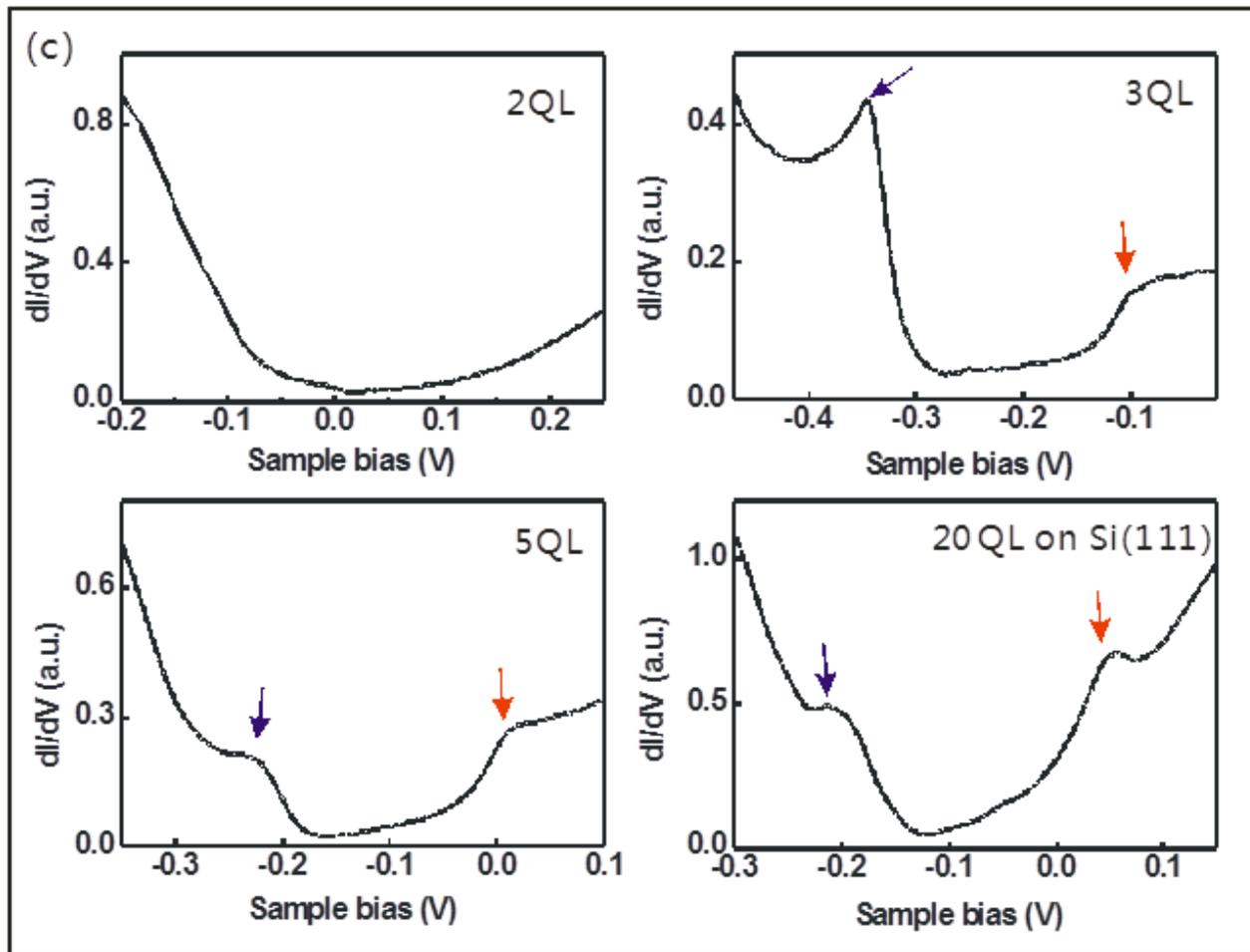
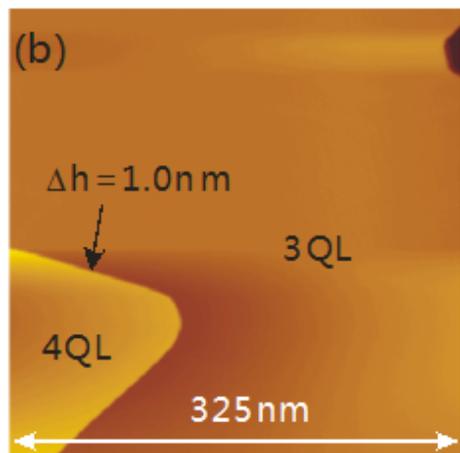
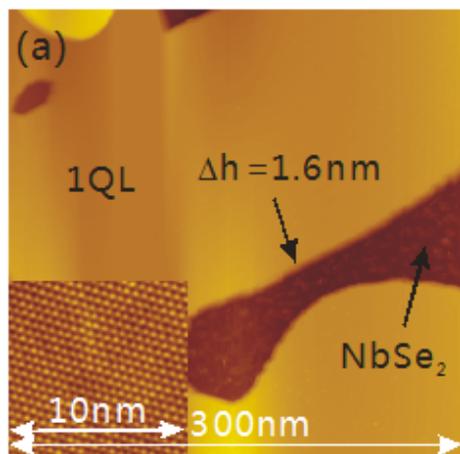
The Coexistence of Superconductivity and Topological Order in the Bi₂Se₃ Thin Films

Mei-Xiao Wang,^{1*} Canhua Liu,^{1*} Jin-Peng Xu,¹ Fang Yang,¹ Lin Miao,¹ Meng-Yu Yao,¹ C. L. Gao,¹ Chenyi Shen,² Xucun Ma,³ X. Chen,⁴ Zhu-An Xu,² Ying Liu,⁵ Shou-Cheng Zhang,^{6,7} Dong Qian,^{1†} Jin-Feng Jia,^{1†} Qi-Kun Xue⁴

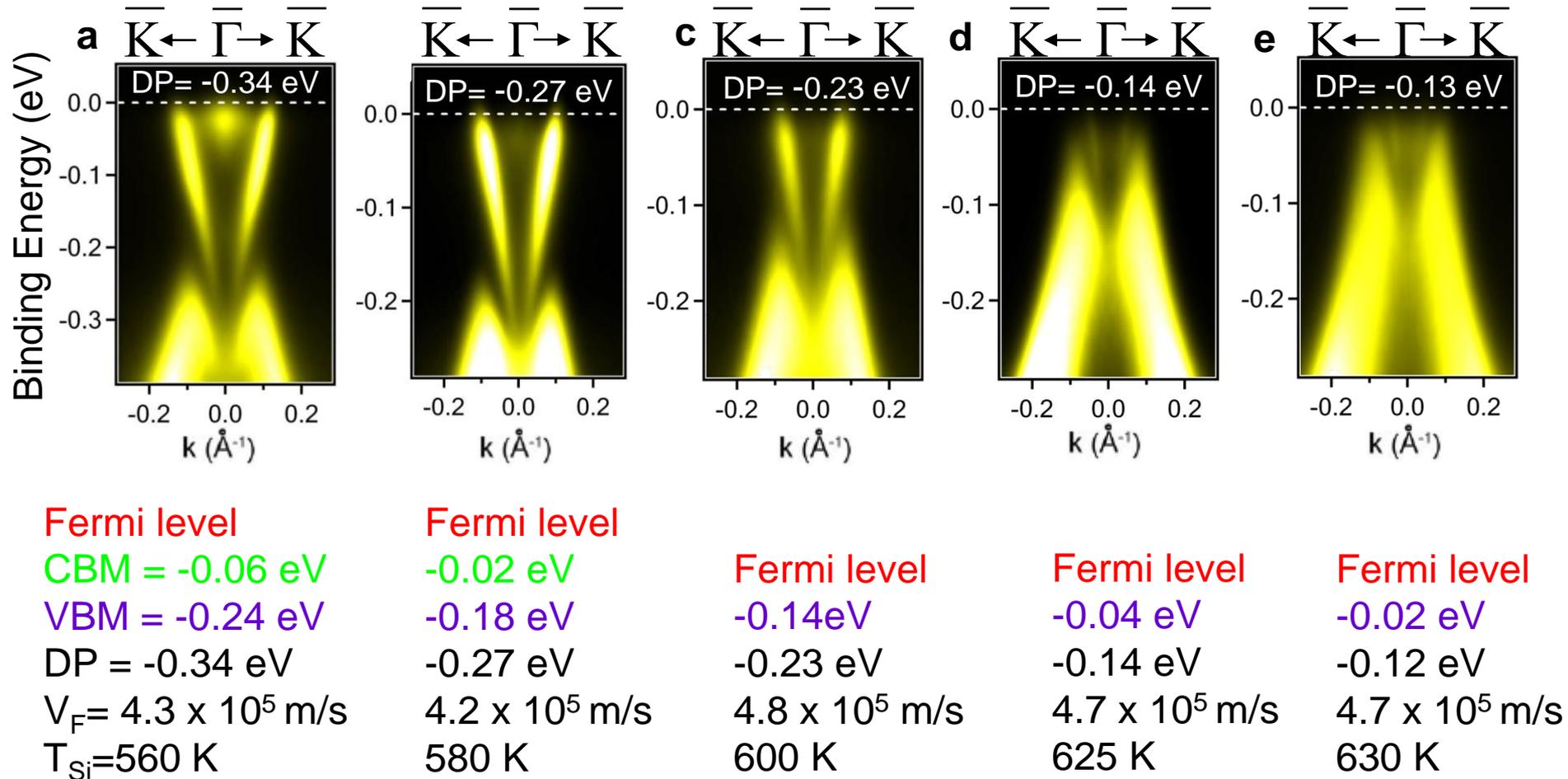
6 APRIL 2012 VOL 336 SCIENCE

Science 336, 52-55 (2012)

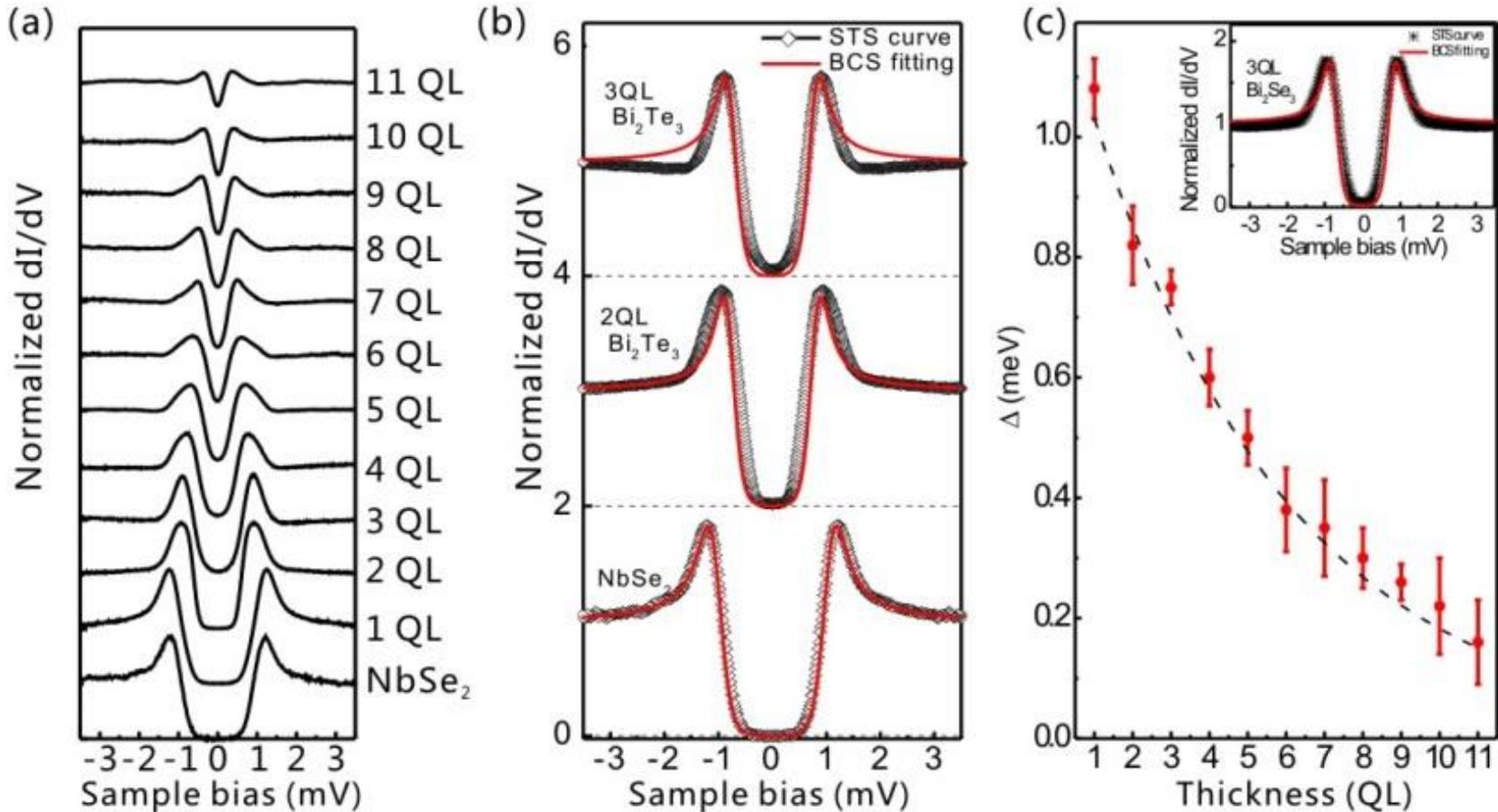
Bi₂Te₃ on NbSe₂



Fermi levels of Bi_2Te_3 films grown at different temperatures



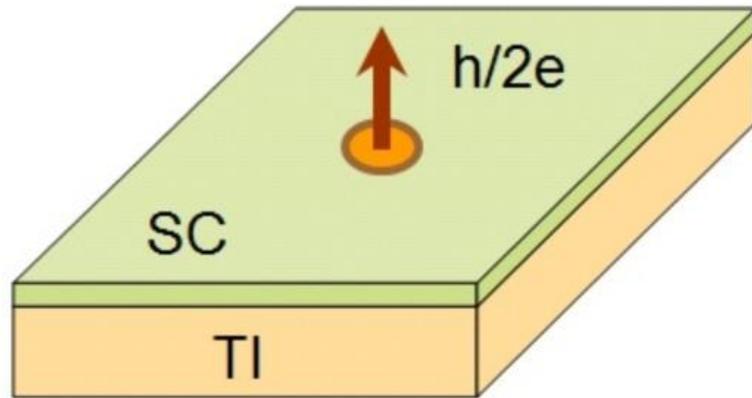
Topological superconductor



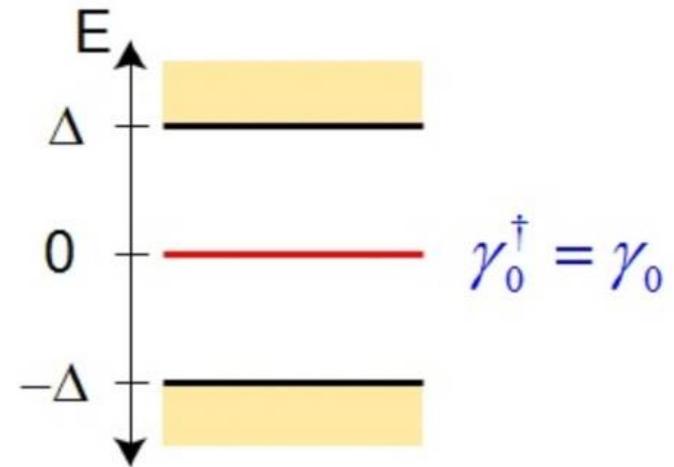
Full gap topological superconductor!

Content

- ① Introduction
- ① Topological superconductor by proximity effects
- ① **Majorana fermion in the vortex**
- ① Summary



Quasiparticle Bound state at $E=0$

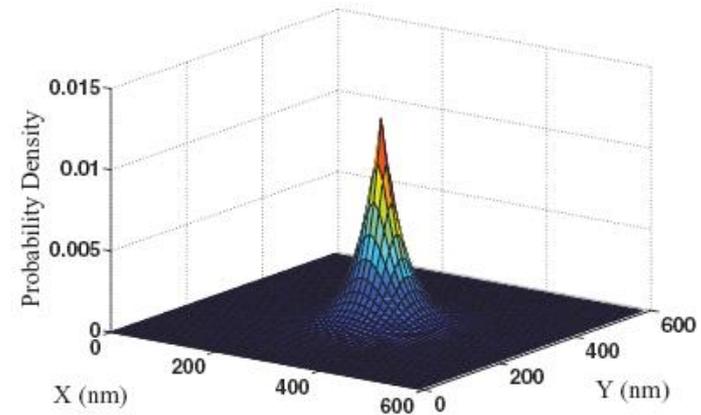
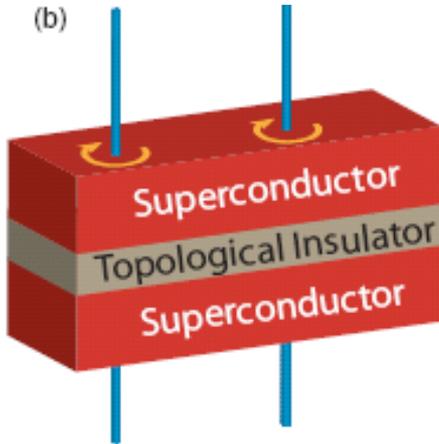


Majorana Fermion γ_0

L. Fu and C.L. Kane PRL 100, 096407 (2008)

C.W.J. Beenakker, *Ann. Rev. Condens. Matter Phys.*, 4:113-136 (2013)

Vortex lines in topological insulator-superconductor heterostructures

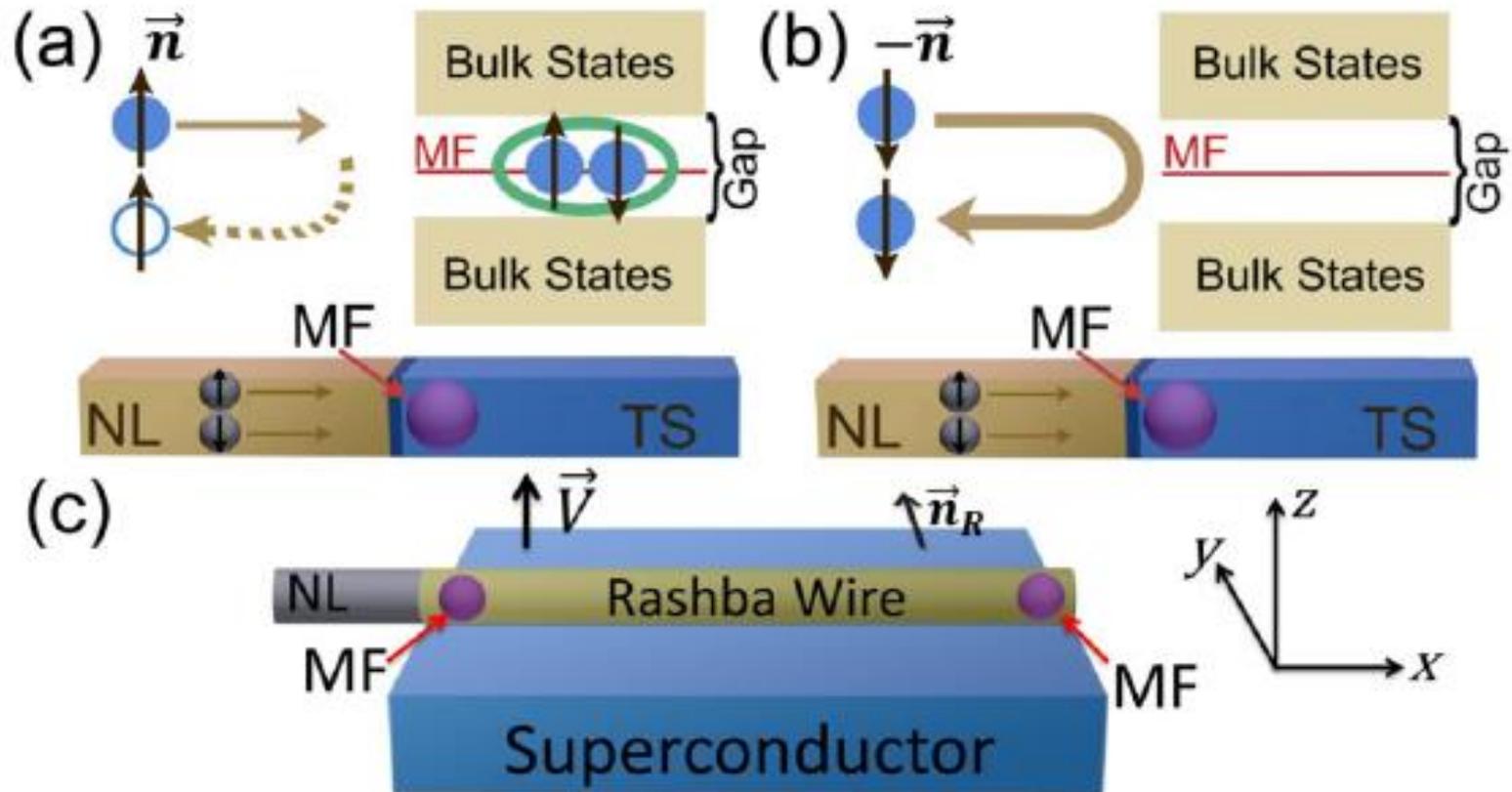


- ⊗ In most case, the bulk effects can be ignored
- ⊗ A Majorana fermion is stable with a spatial extent~40nm
- ⊗ Chemical potential & Majorana states
- ⊗ Majorana fermions can survive for thick samples

Hughes group, **PRB 84**, 144507 (2011)
PRB 87, 035401 (2013)

Selective Equal-Spin Andreev Reflections Induced by Majorana Fermions

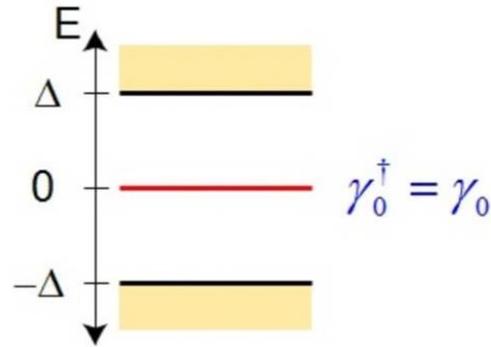
P. A. Lee, K. T. Law, PRL 112, 037001 (2014)



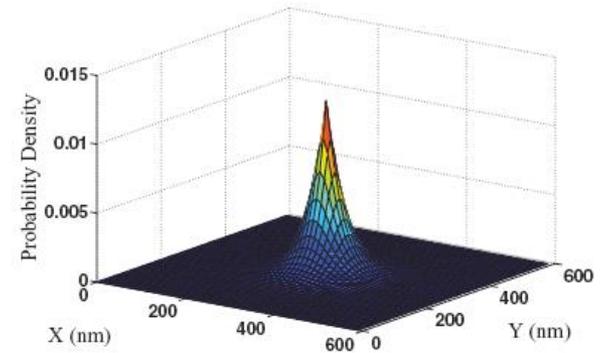
The SESARs can also be used to detect MFs if spin-polarized leads are used.

Features of MF

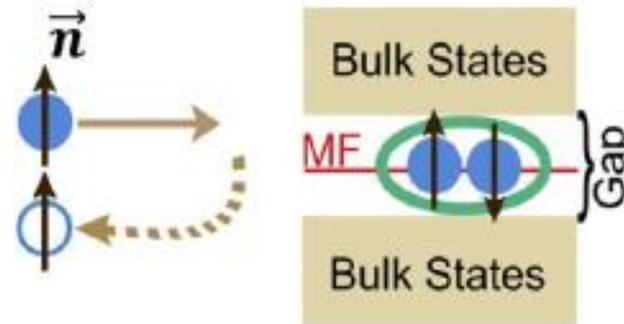
Zero energy



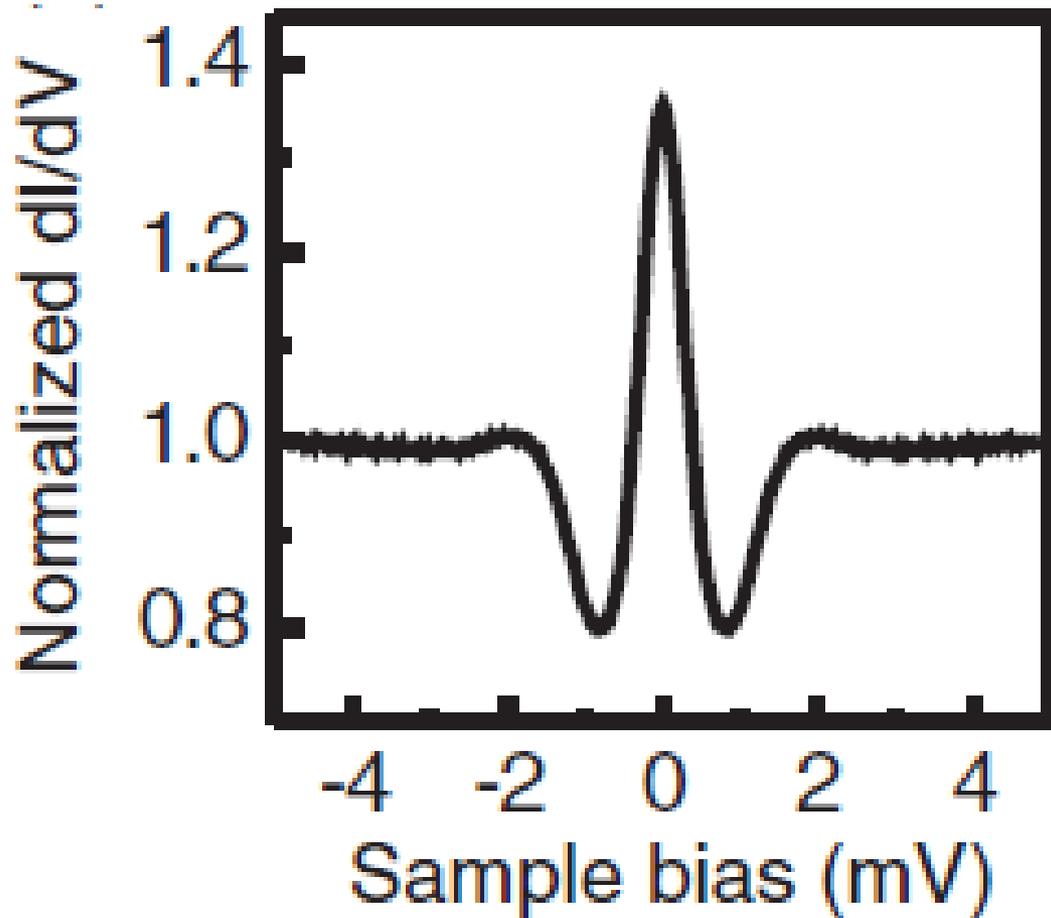
Cone shape distribution



SSAR or SESAR



Detect Majorana fermion by zero-energy peak



...ning Fermi energy

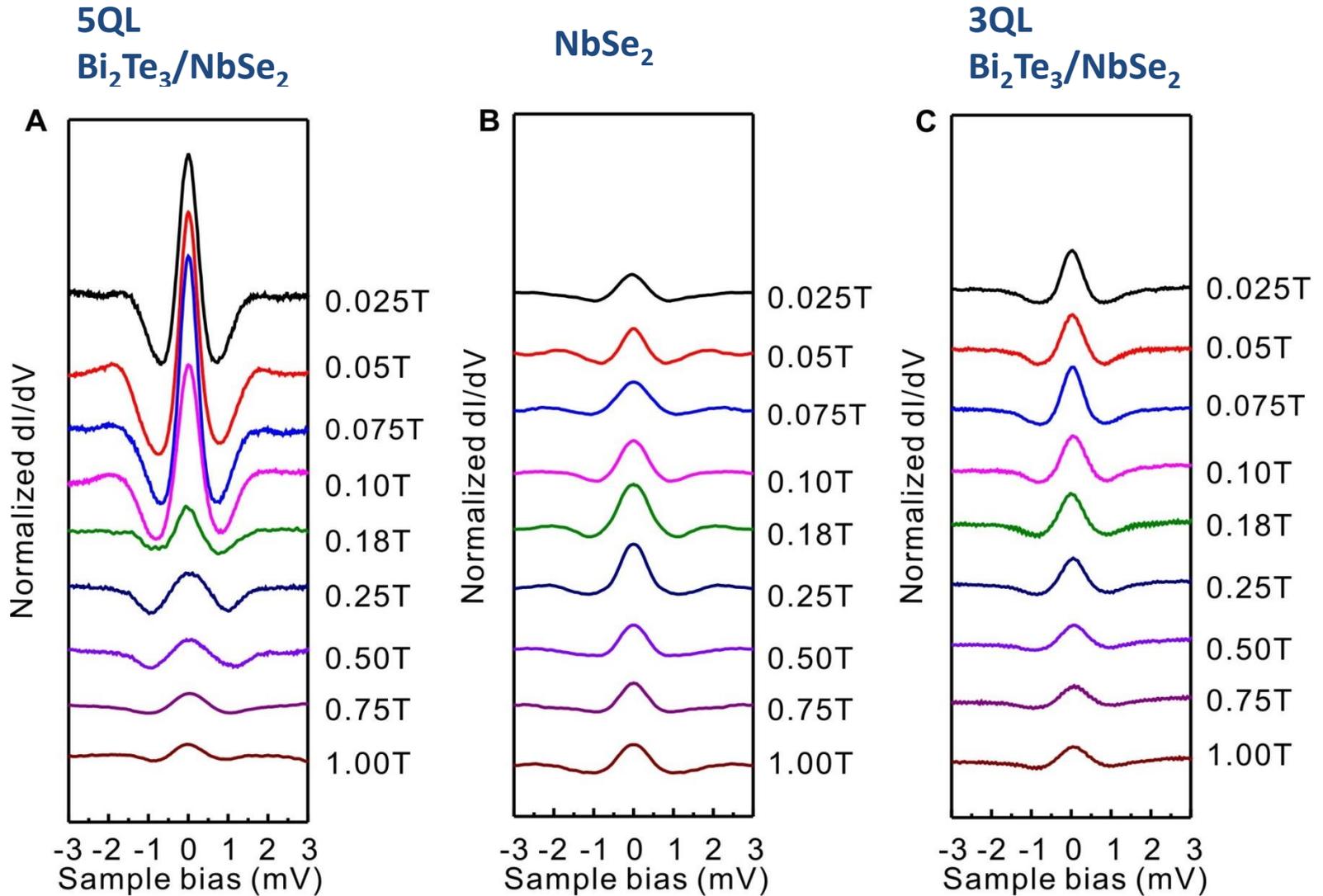
$$\sqrt{\Delta^2 + E_D^2}$$

...or 5QL $\text{Bi}_2\text{Te}_3/\text{NbSe}_2$

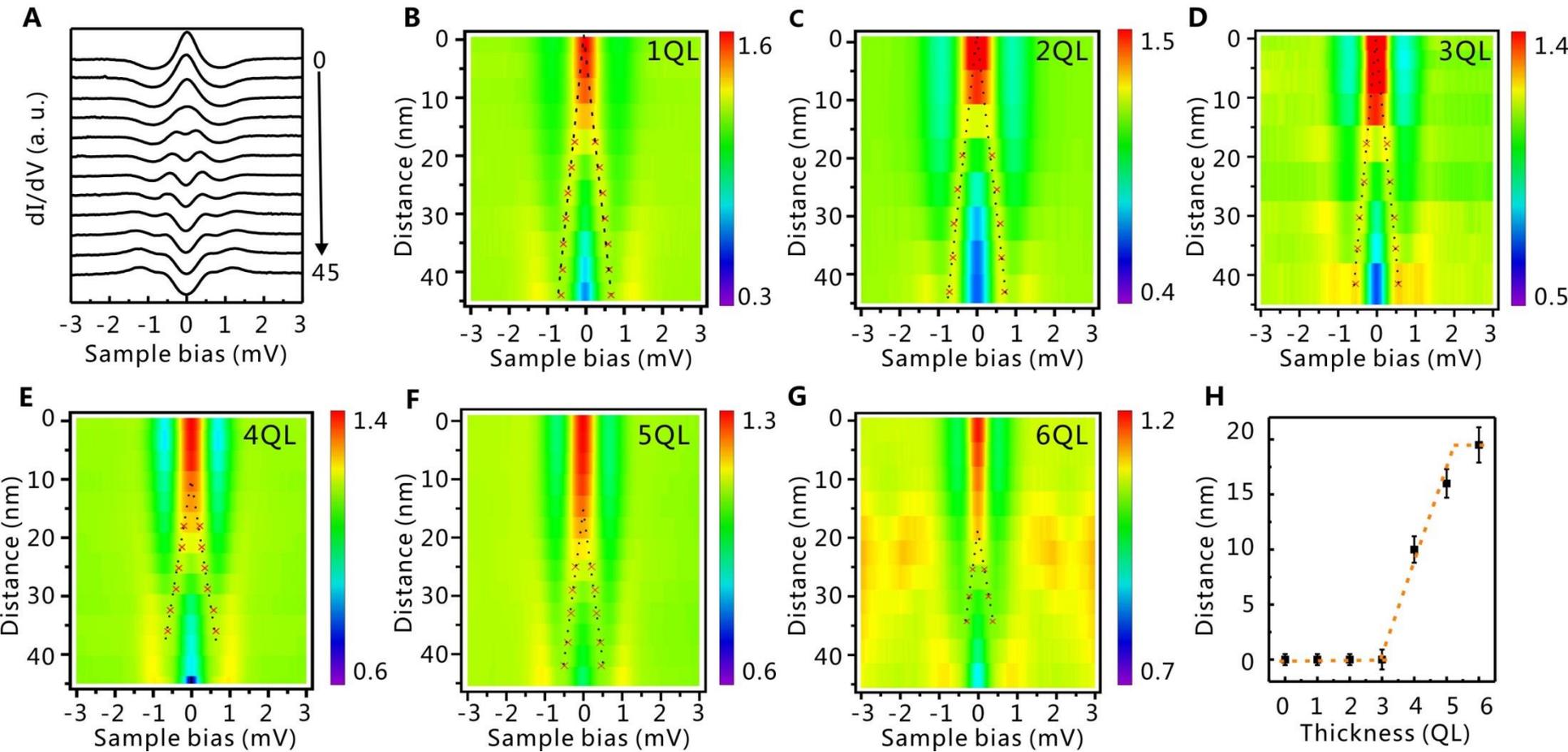
ow:
 0.8meV $E_D \sim 200\text{meV}$
 $E_g \approx 0.028\text{meV} \sim 30\text{mK}$

$E_D \sim 5\text{meV} \rightarrow E_g \approx 0.1\text{meV}$
 $T \sim 300\text{mK} \rightarrow \text{Thermo energy} \sim 0.026 \text{ mV}$

Magnetic field dependence of ZBP



Spatial distribution of Majorana fermion

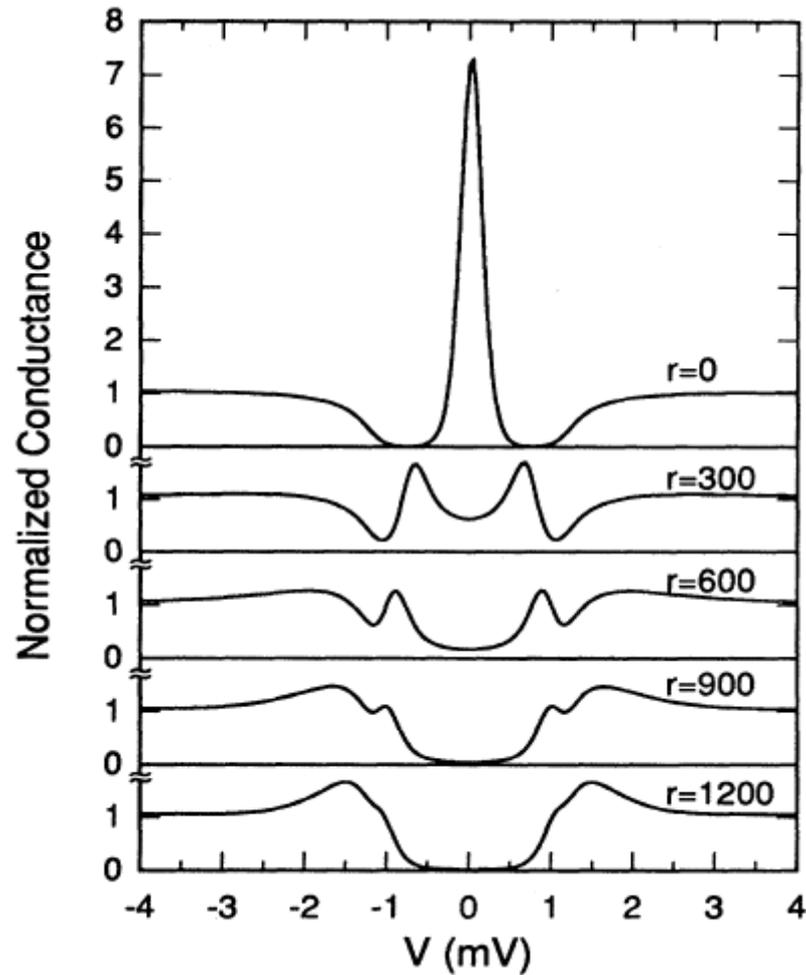


Splitting of Zero-bias peaks

Non-zero splitting

Core states splitting in CSC

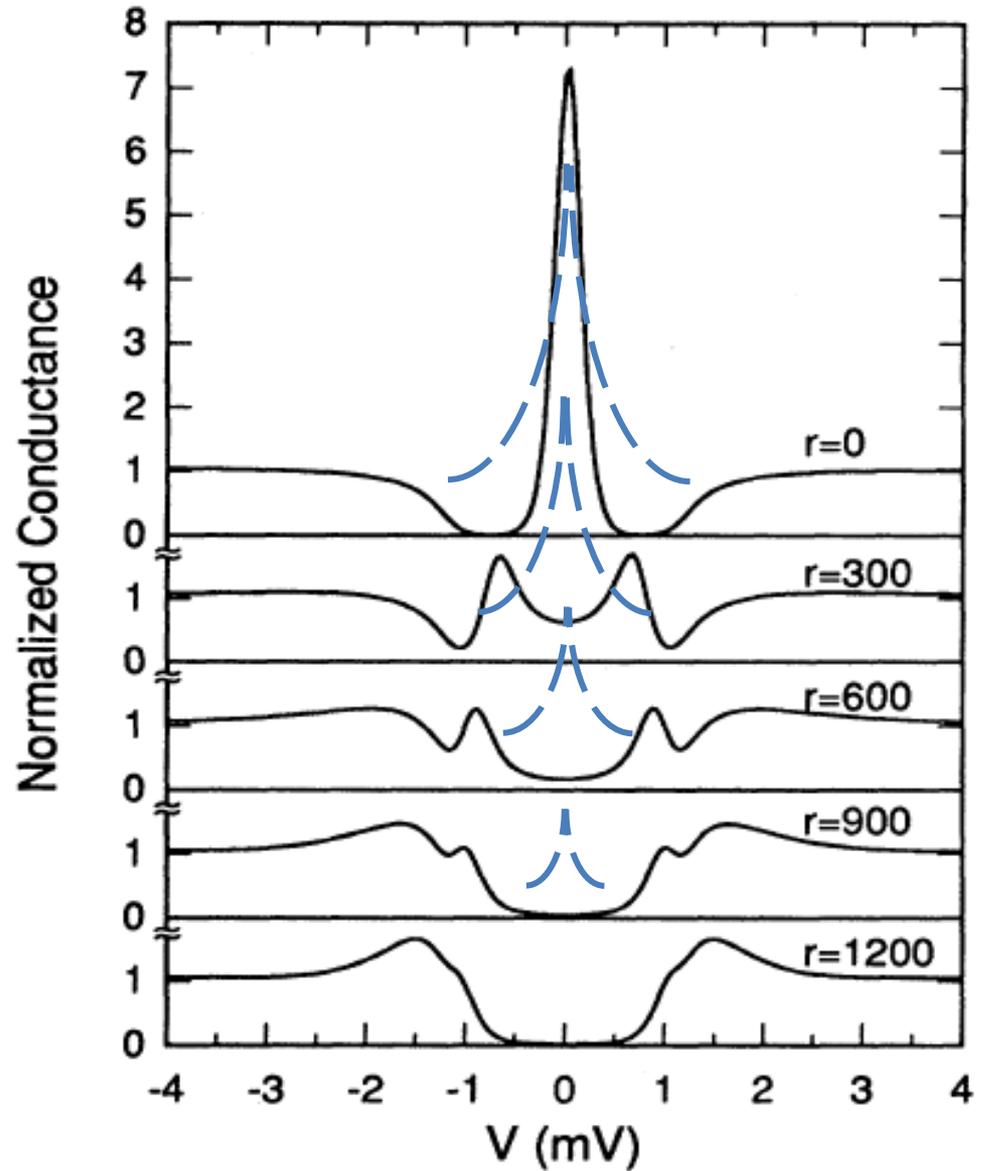
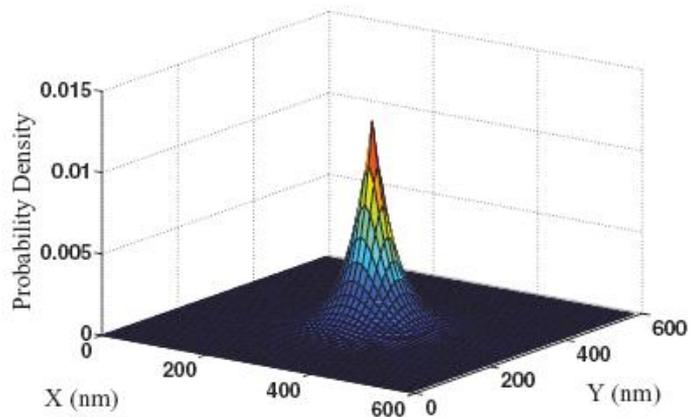
NbSe₂



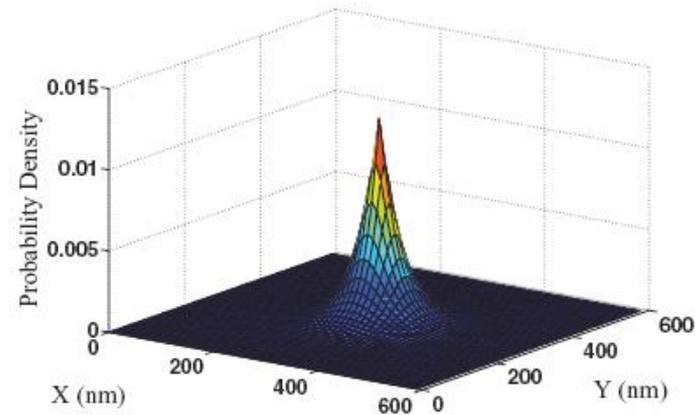
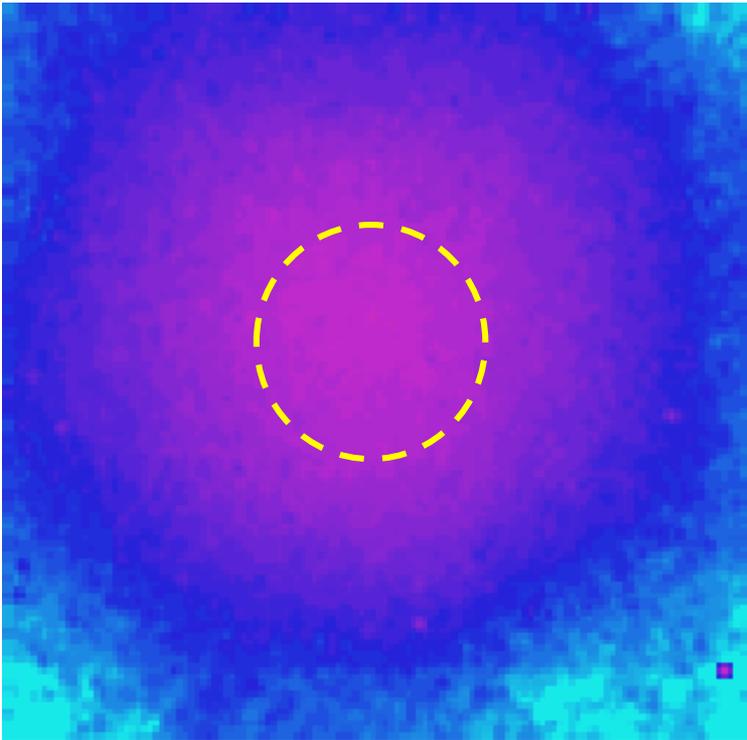
Hess et al., Phys. Rev. Lett. 62, 214 (1989)

F. Gygi, M. Schluter, PRB 43, 7609 (1991)

Core states splitting in TSC



Spatial extent of Majorana fermion



- A Majorana fermion in a spatial extent $\sim 40\text{nm}$

Hughes group,
PRB 84, 144507 (2011)
PRB 87, 035401 (2013)

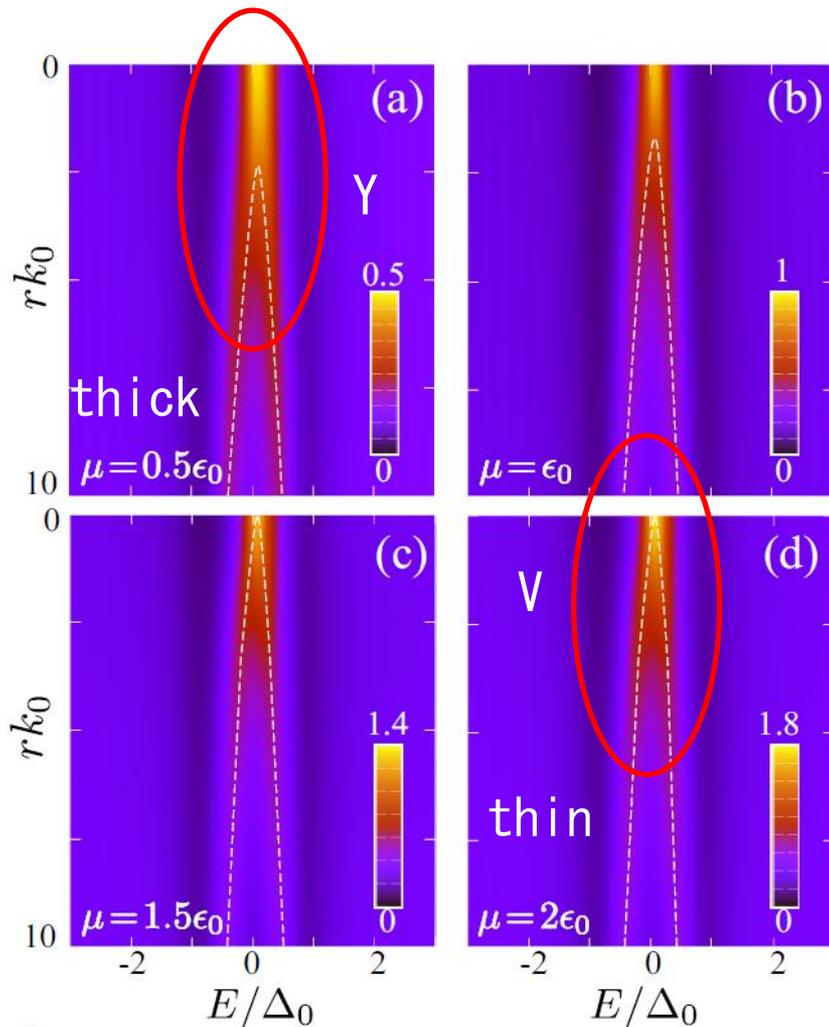


Strong evidence for existence of Majorana mode

Evolution of DOS with thickness

PRL 115, 177001, 2015

□ Energy-space distribution of DOS of quasiparticles: dI/dV in experiments



smearing factor
in energy

$$\eta = 0.2\Delta_0 \sim 4\text{K}$$

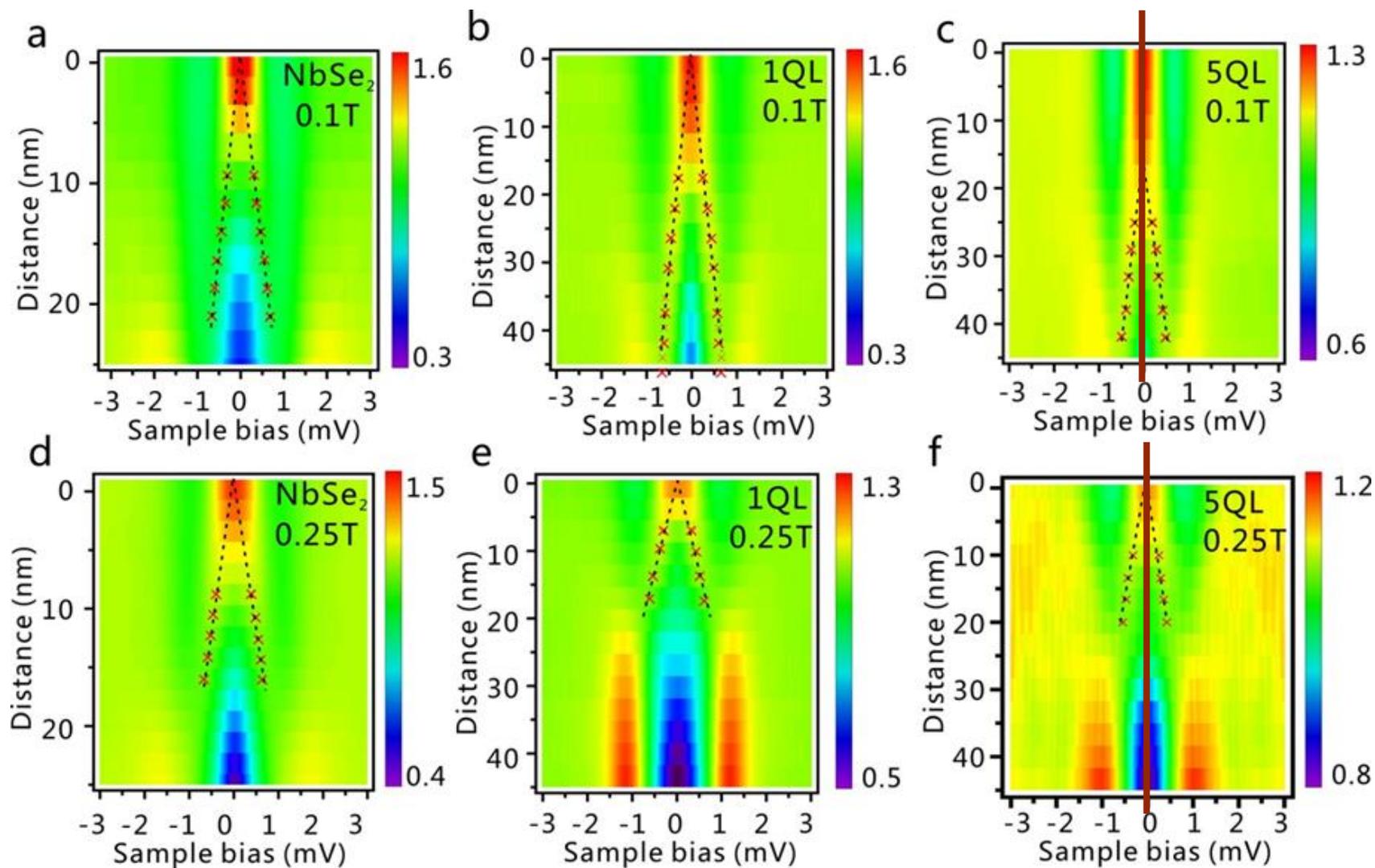
Y shape \Leftrightarrow w MBS
V shape \Leftrightarrow w/o MBS

full agreement with experiments !

□ Thickness vs. chemical potential

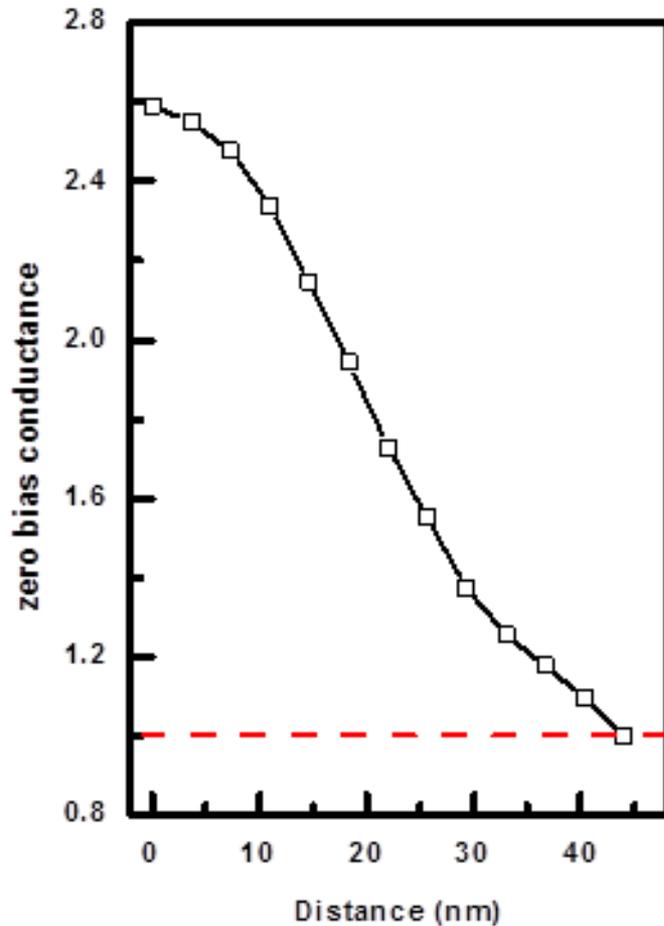
theoretically thickness only cannot induce phase transition, but μ can.

More evidence

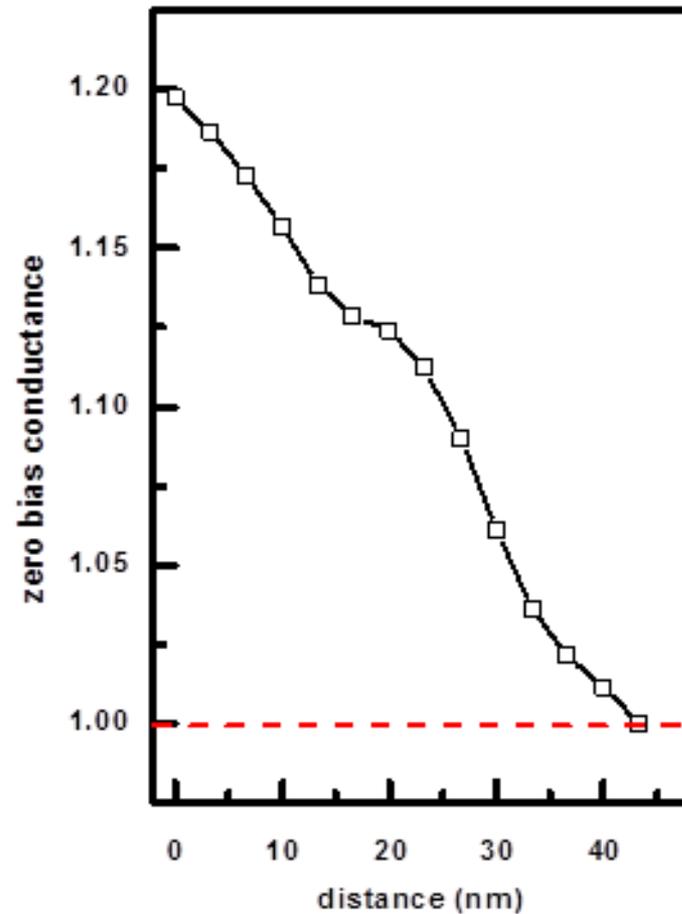


Spatial distribution at zero energy

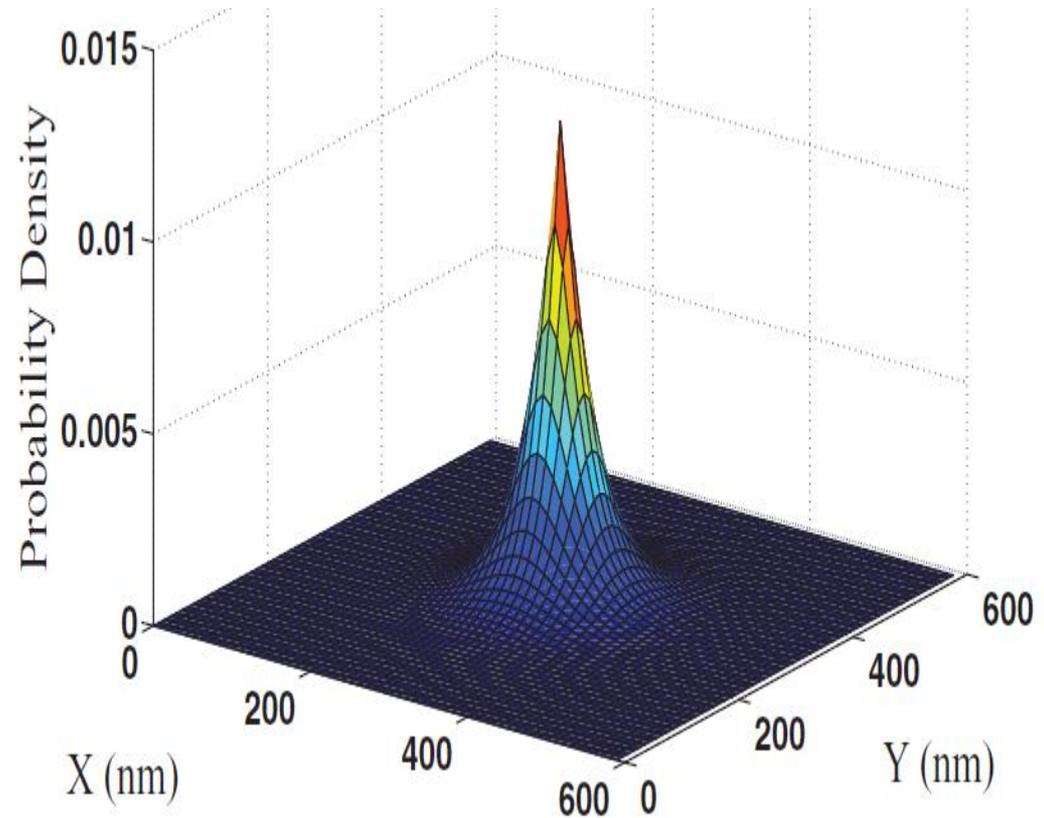
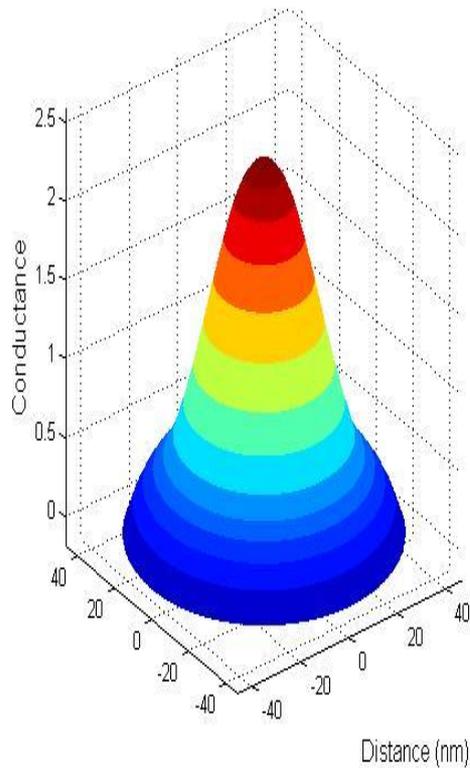
5QL @ 0.1T



5QL @ 0.18T



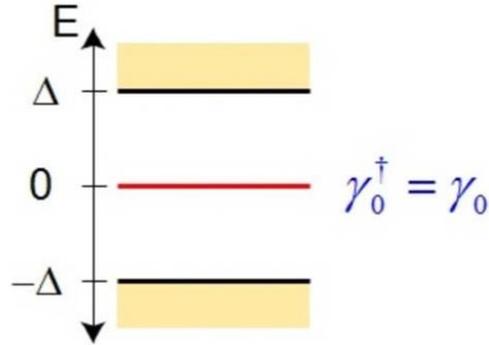
Experimental spatial distribution of Majorana fermion



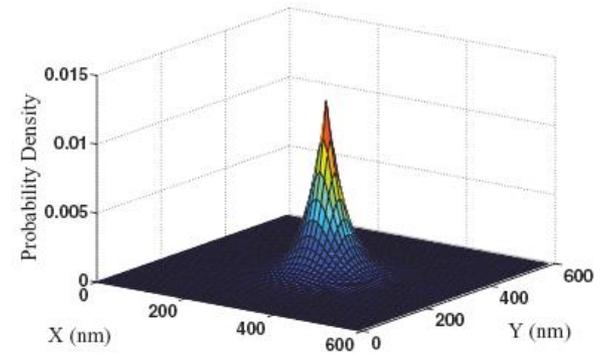
Substrate the data with MF (5QL at 0.1T) by the data without MF (5QL at 0.18T), one should get the contribution of MF (left). The right one is the calculated probability distribution of a Majorana bound state in vortex(PRB84,144507).

Features of MF

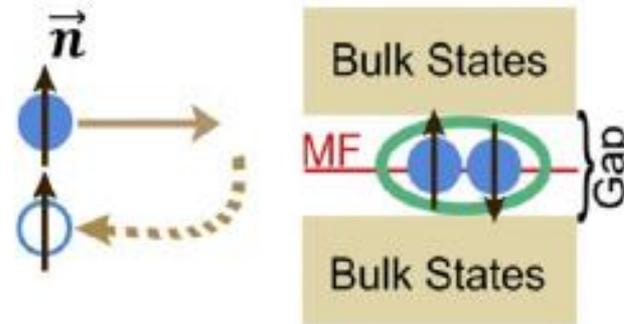

Zero energy



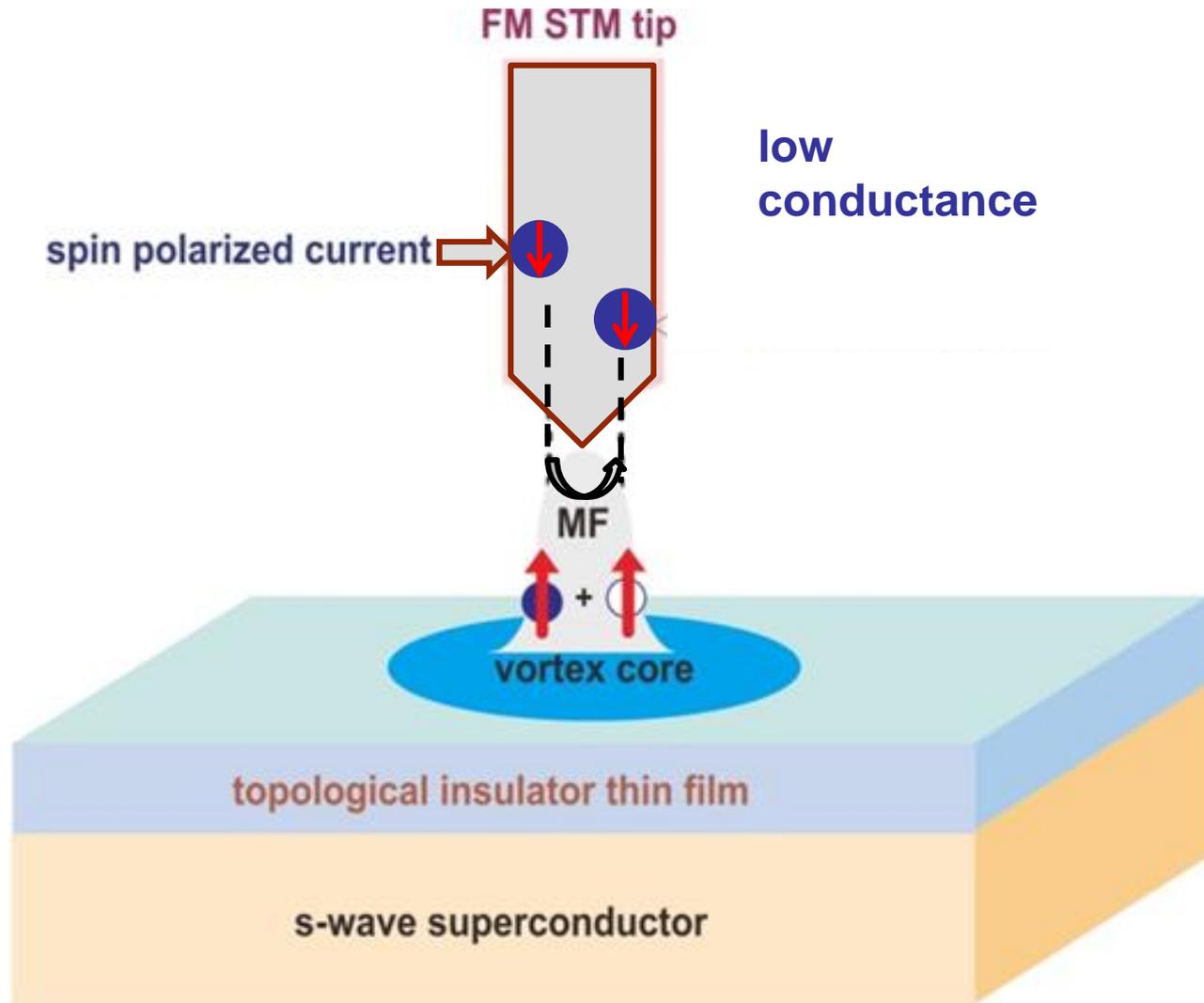

Cone shape distribution



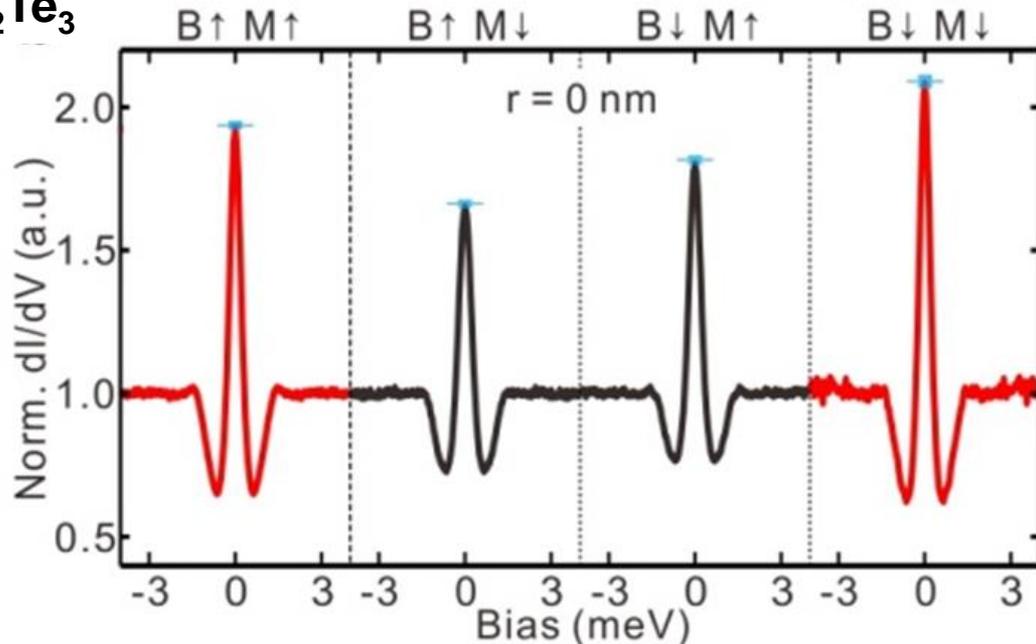
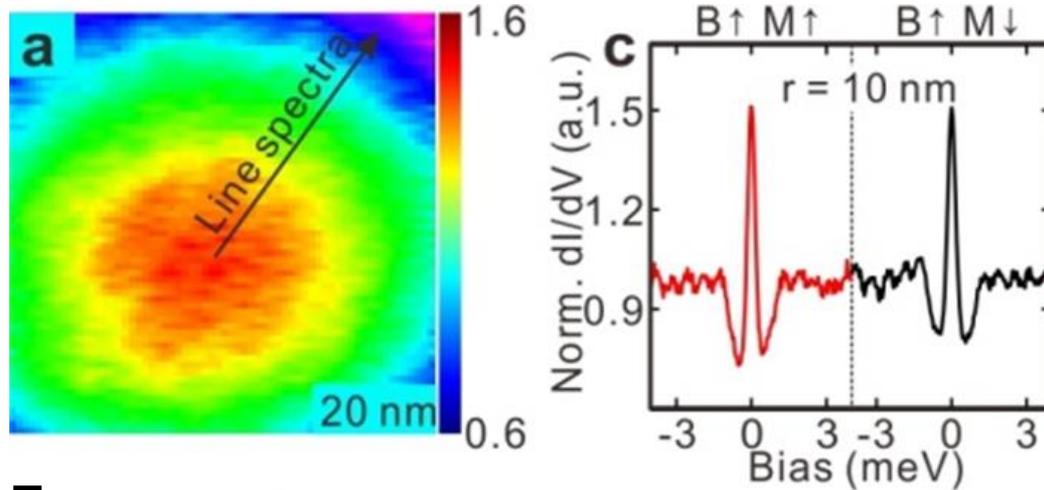

SSAR or SESAR



Spin selective Andreev reflection



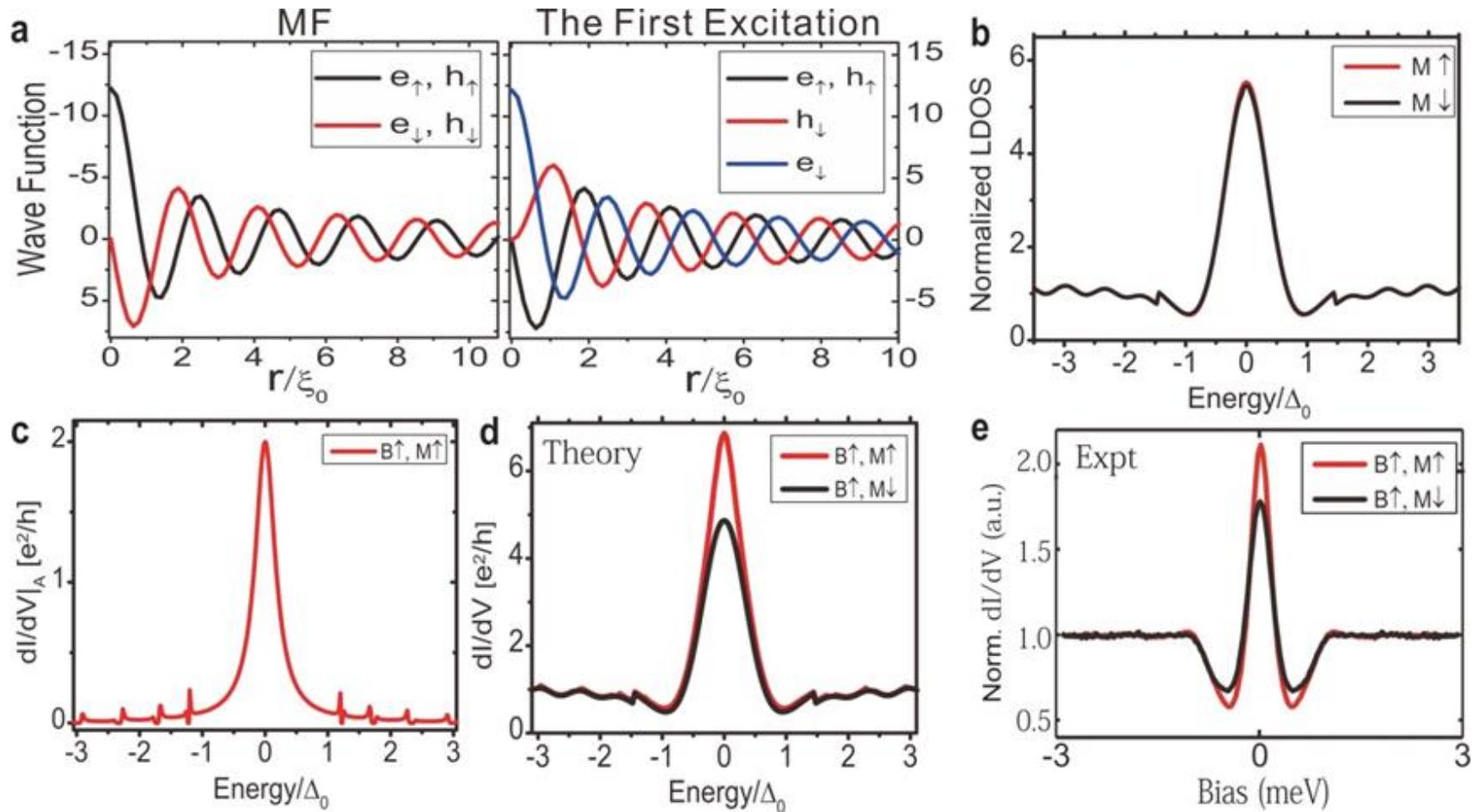
SSAR observed by SPSTM



5QL Bi_2Te_3

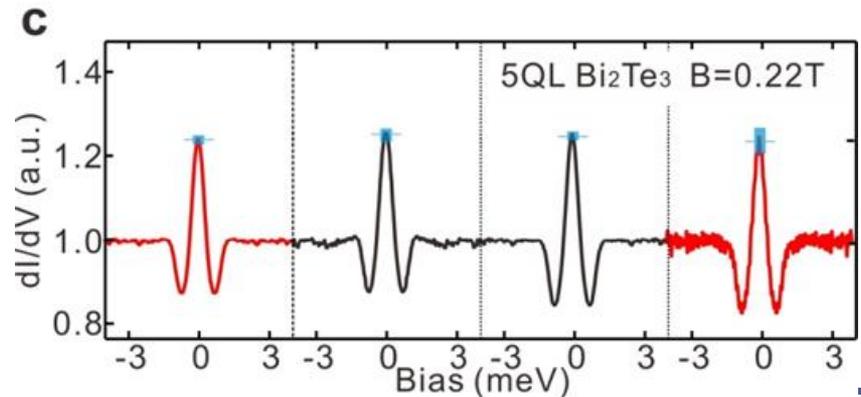
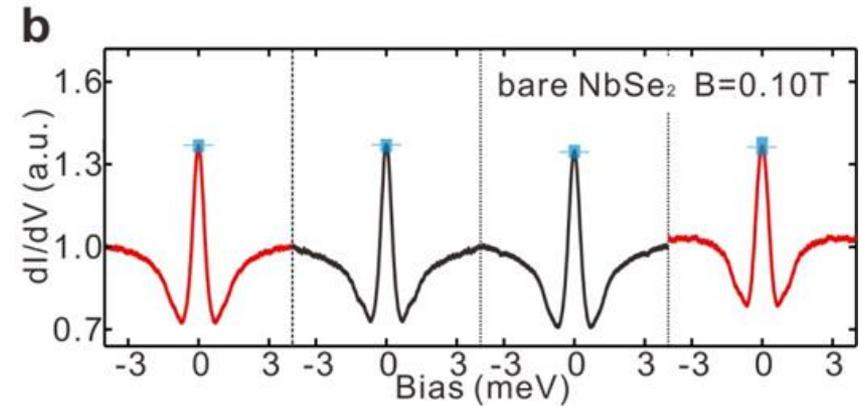
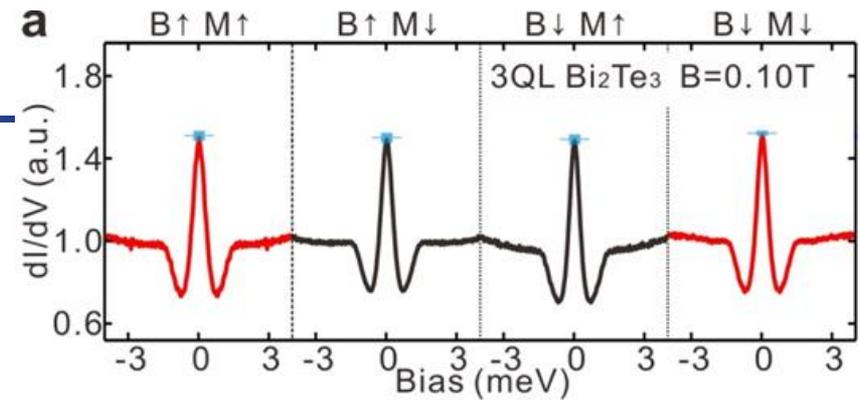
- ⊙ Non-magnetic
- ⊙ Low applied field (0.1T)
- ⊙ SSAR can only induced by MF

Comparison with model calculations

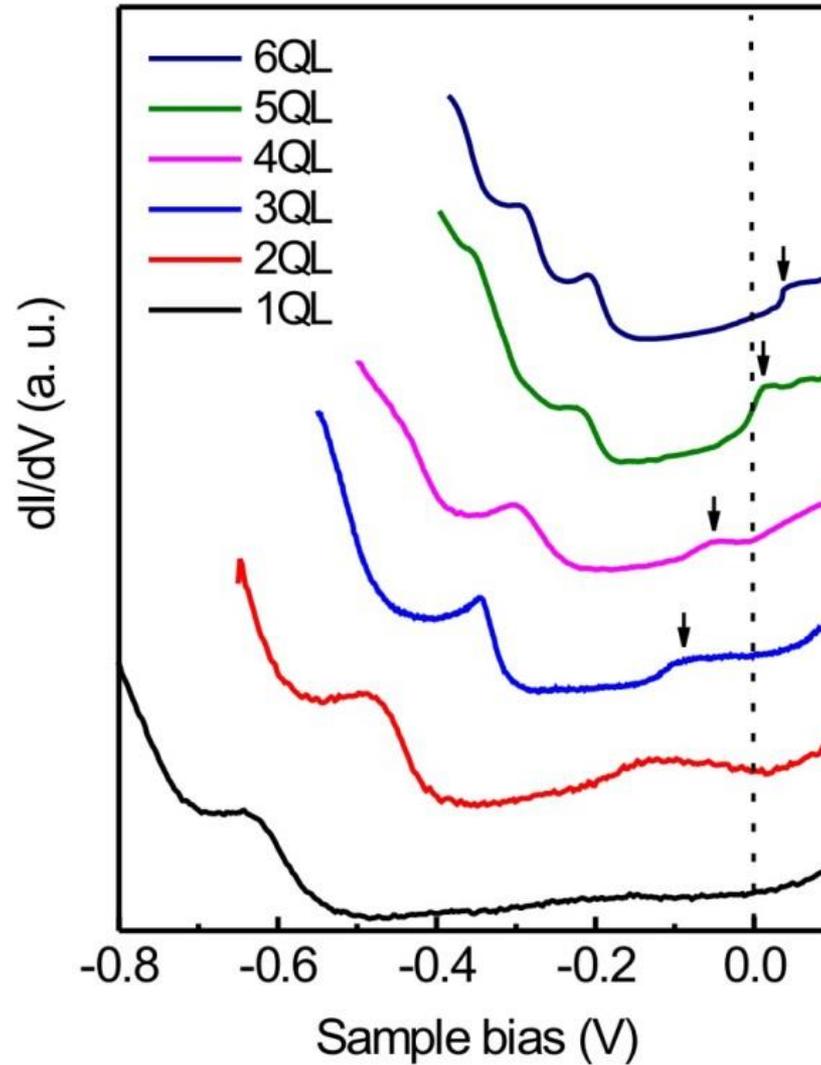


Compare with no-MF cases

arXiv:1603.02549
Phys.Rev.Lett. 116, 257003 (2016)

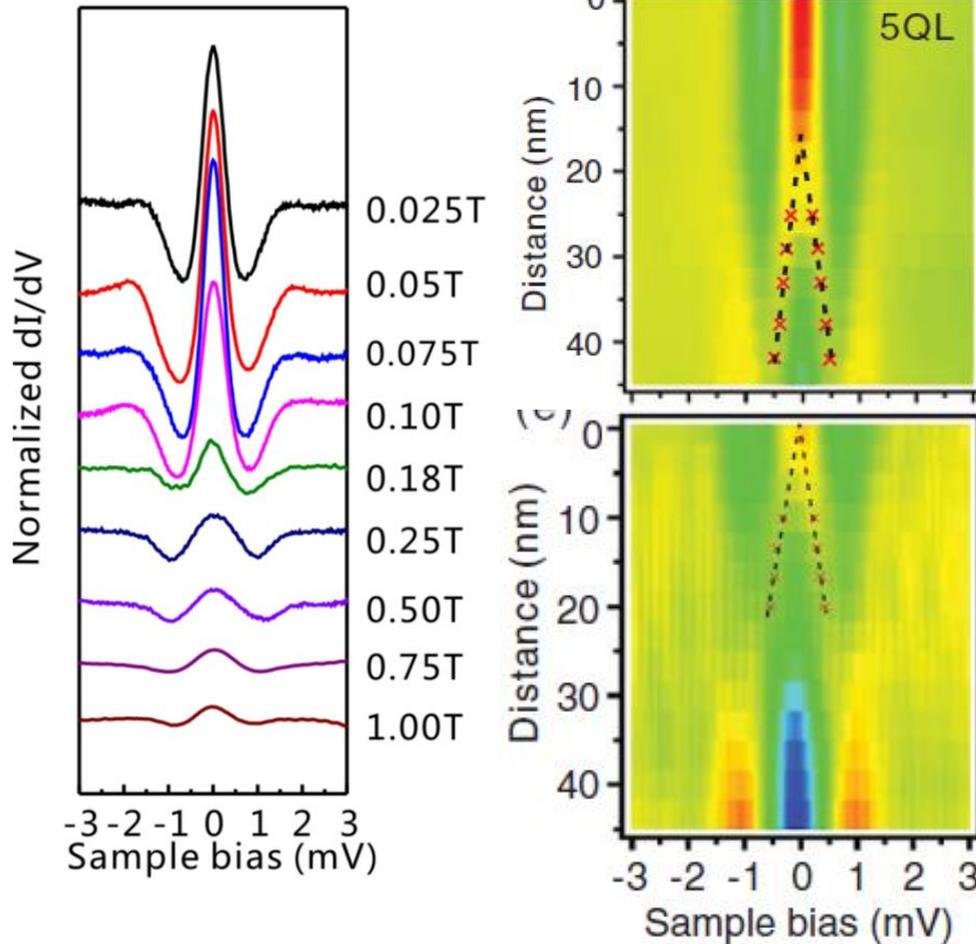


Why transition at 4QL?

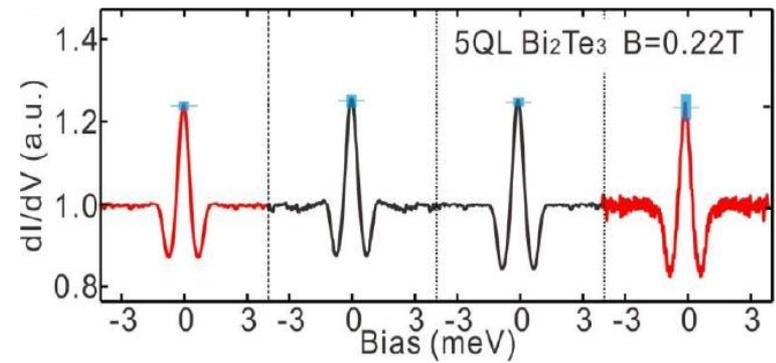
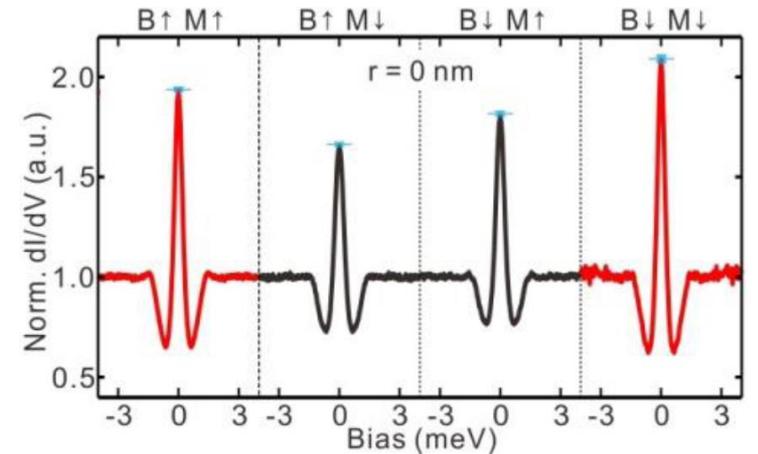


All evidences are consistent

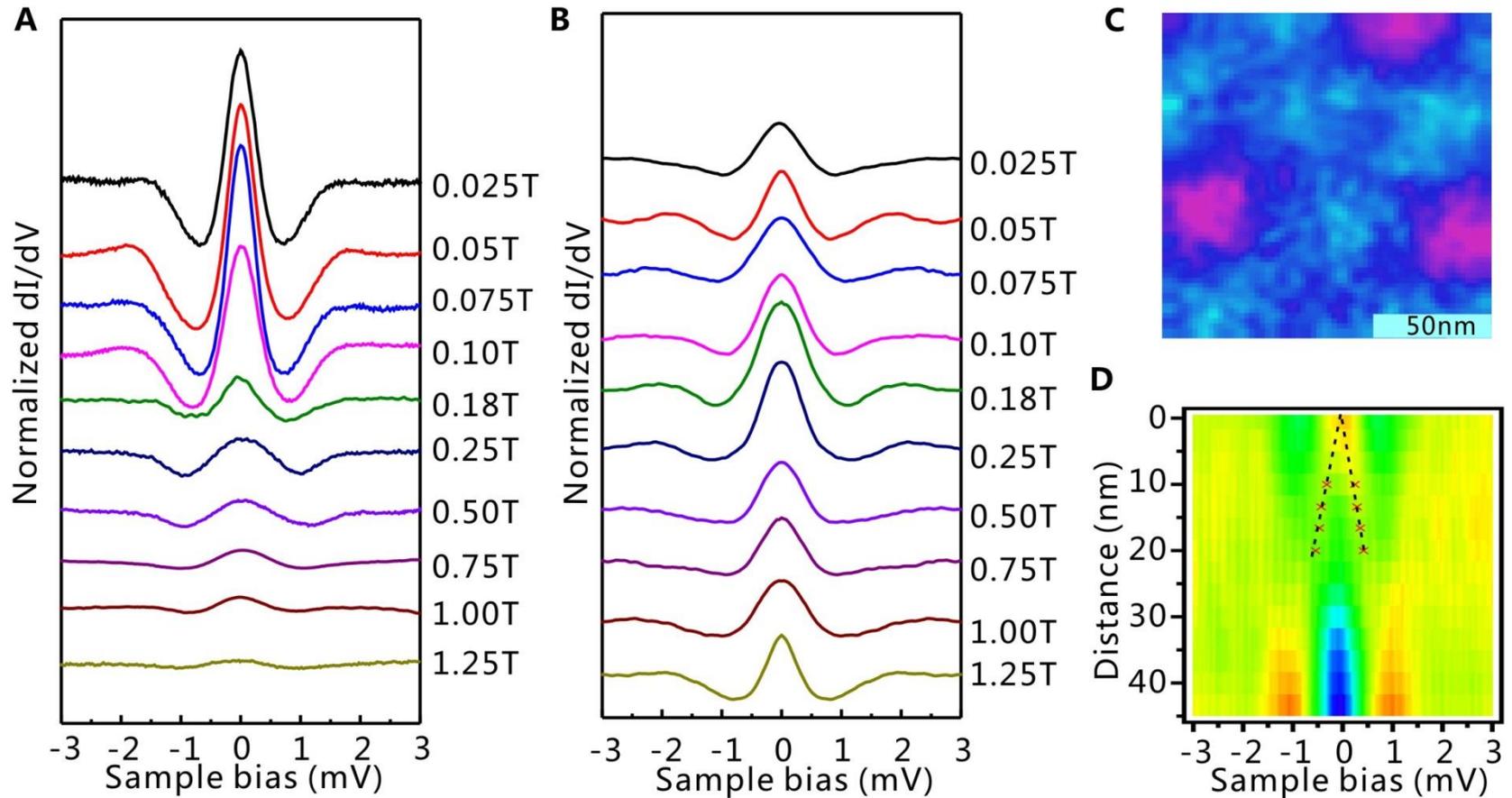
5QL Bi_2Te_3



5QL Bi_2Te_3 $B=0.1\text{T}$



Turn MF on/off



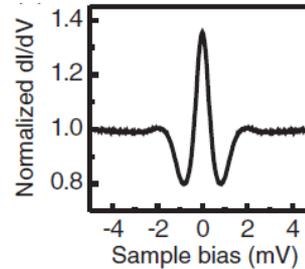
PRL 114, 017001 (2015)

npj Quantum Materials (2017) 2:34 ; doi:10.1038/s41535-017-0037-4

Summary

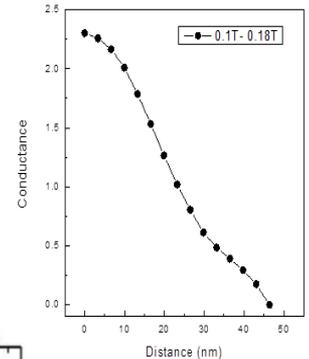

Artificial topological superconductor by proximity effect


Zero energy

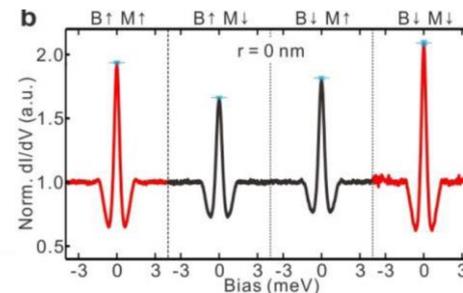


Need more work


Cone shape distribution



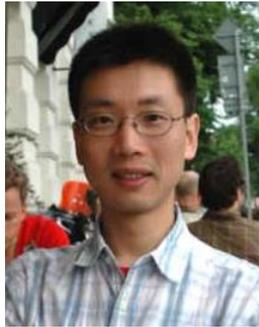

SSAR or SESAR



Acknowledgments



QK Xue



C.H. Liu



C.L. Gao



D. Qian



Zhuan Xu



S.C. Li



F.C. Zhang



Liang Fu

Theoreticians: F. Liu, S.B. Zhang, X.C. Xie, Z. Fang, X. Dai, S. Q. Shen, Xiao-Liang Qi, Shou-Cheng Zhang., Q.H. Wang, Y. Chen, Y. Zhou

M.X. Wang, J.P. Xu, H.H. Sun, G.Y. Wang, D. Xu, L. Miao, F. Yang, M.Y. Yao, Z. F. Wang, K.W. Zhang, Lun-Hui Hu, Chuang Li

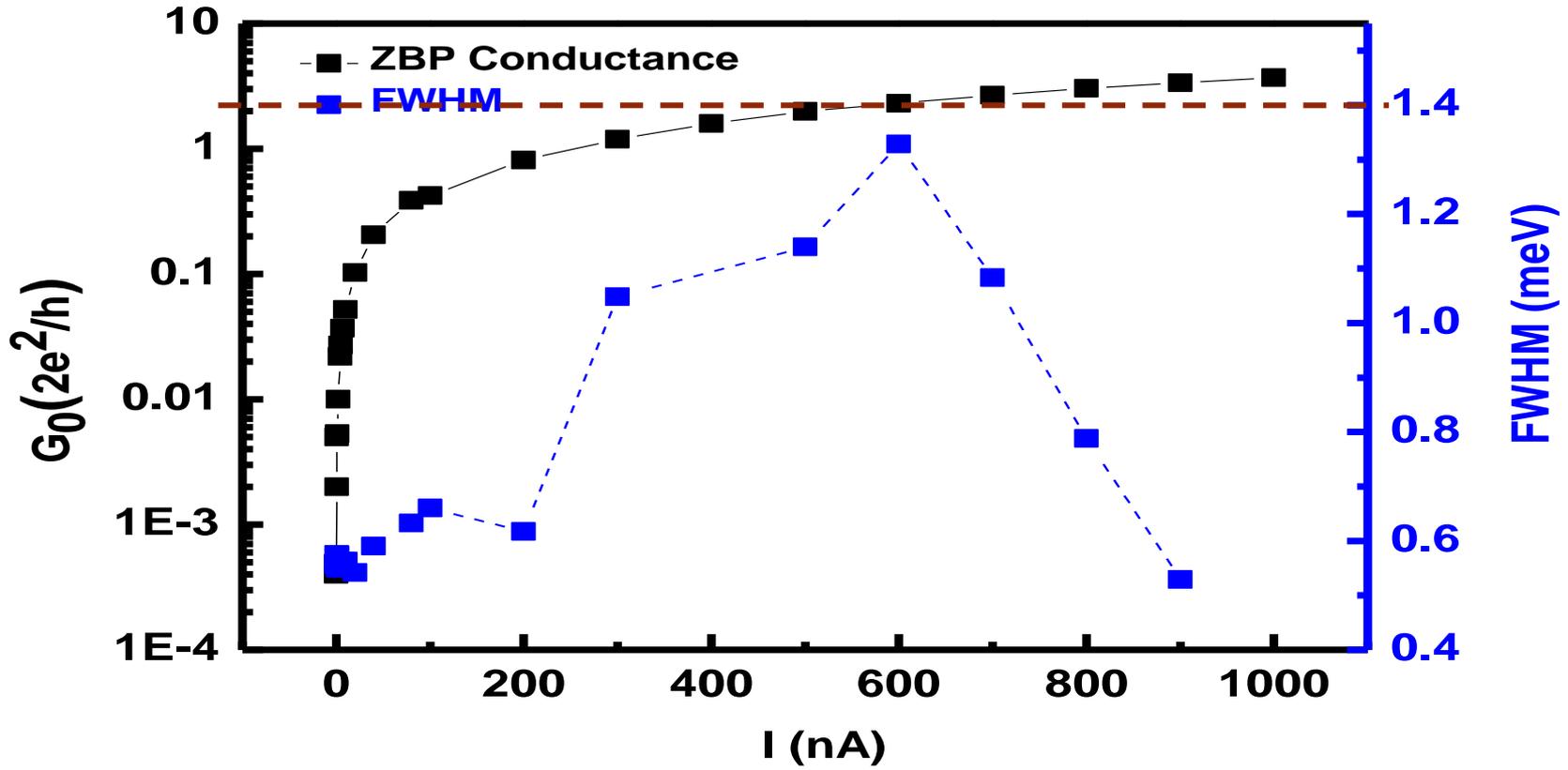
Supported by NSFC, MOST and MOE



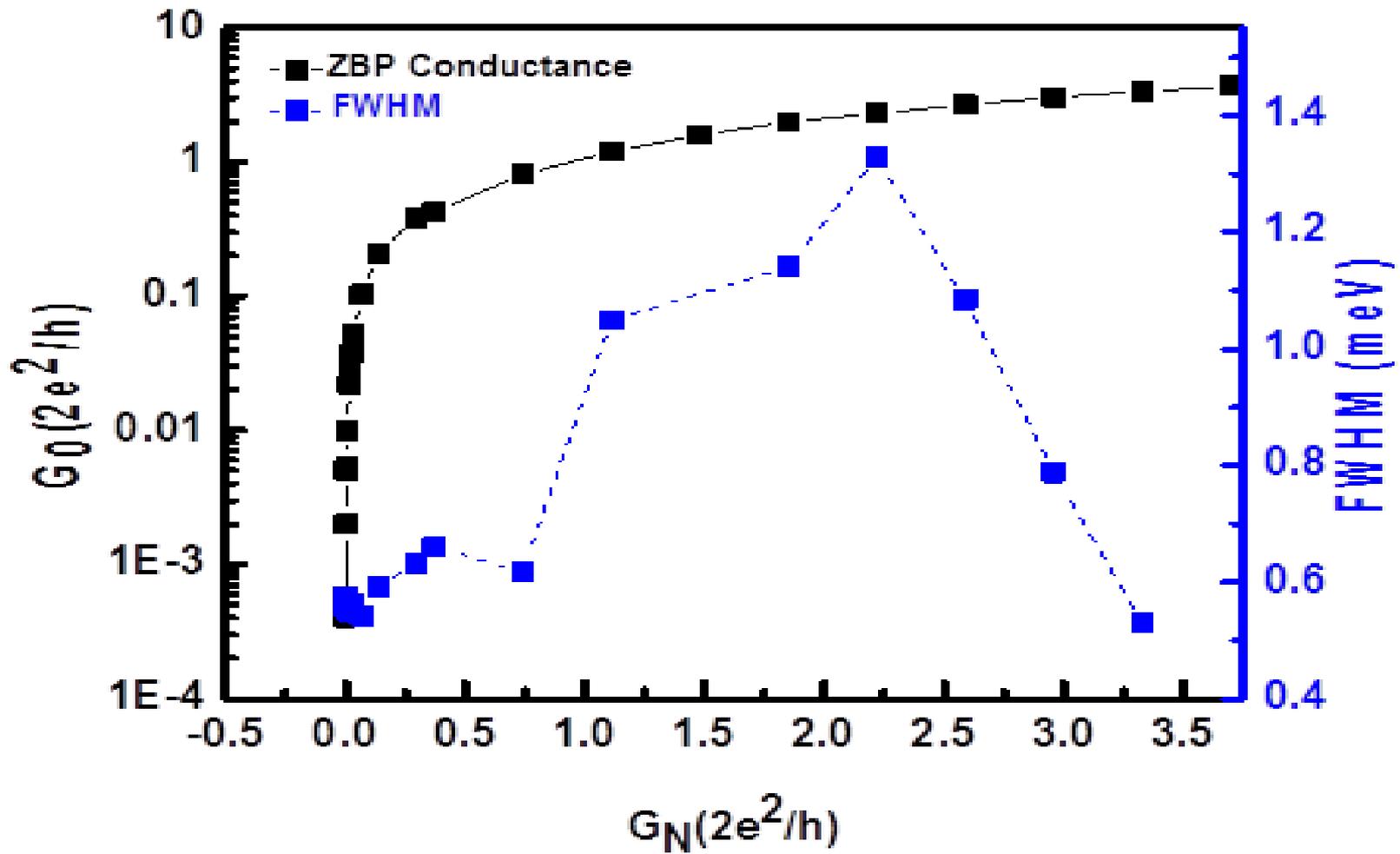
上海交通大學
SHANGHAI JIAO TONG UNIVERSITY

Thank you very much!

Conductance and width of ZBP VS tunneling current



Tip touches sample ~600nA
Conductance saturate: $2G_0$
MZM+2 normal modes

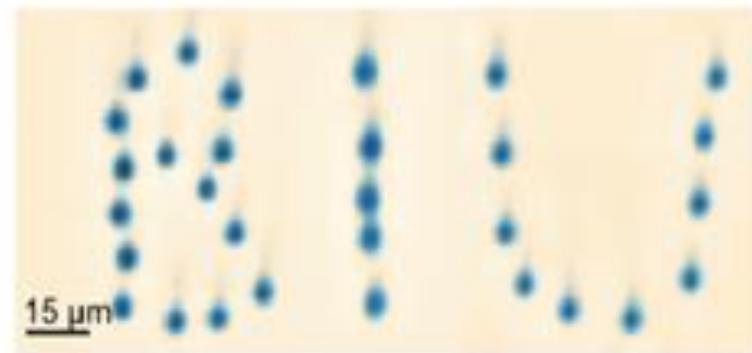
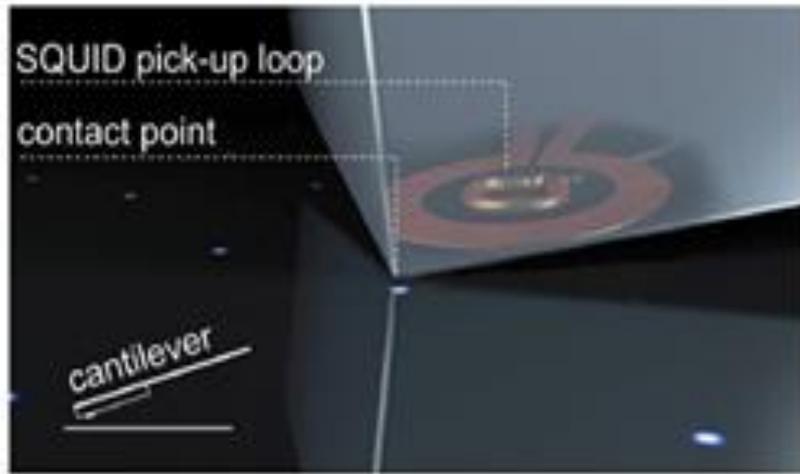


Possible application in topological quantum computing

Advantages

- ① Simple
 - ① Low magnetic field $\sim 0.1\text{T}$
 - ① E_f can be tuned, no gate is needed
 - ① Protected by superconducting gap
 - ① Easy to increase the temperature
 - ① No effect by impurities
 - ① 2D system, easy to manipulate
-

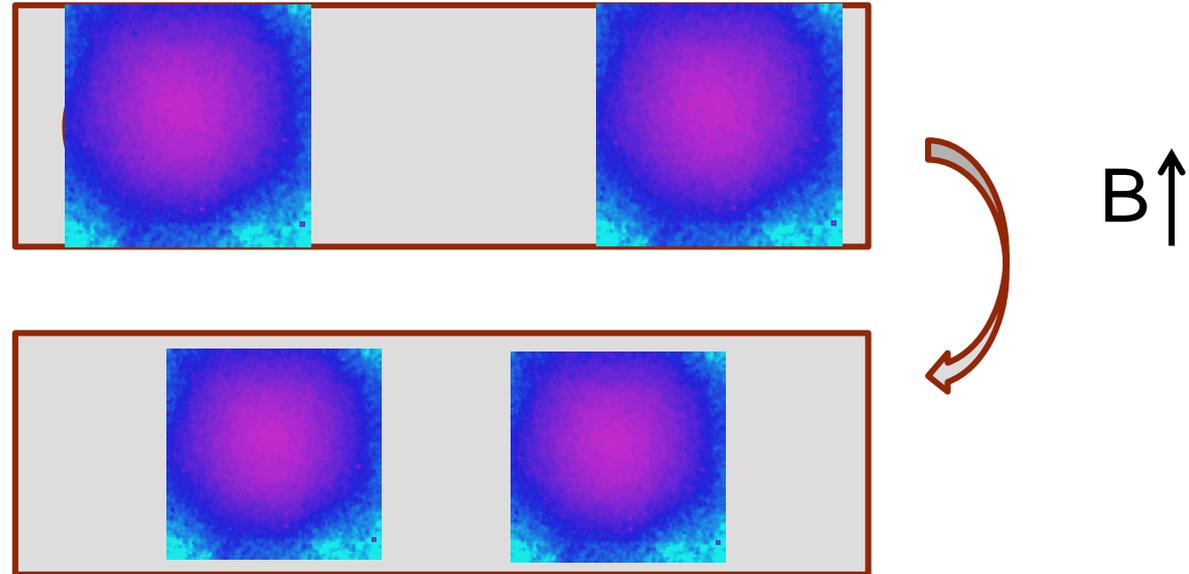
Braiding



- ⊙ Vortex can be easily moved and positioned by scanning SQUID

Fusion

⊙ Fusion = read out



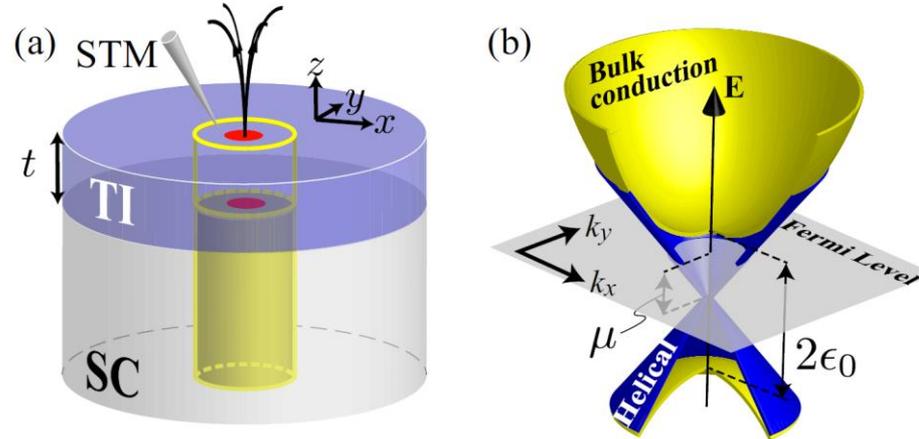
⊙ How to detect the result?

- Junction? Microwave? Spin?.....

Model system

T. Kawakami and XH: arXiv.1506.03194

□ Geometry



□ Hamiltonian

$$\mathcal{H}_{\text{BdG}} = \begin{pmatrix} \hat{\mathcal{H}}_{\text{TI}}(\mathbf{k}) & \hat{\Delta} \\ \hat{\Delta}^\dagger & -\hat{\mathcal{H}}_{\text{TI}}^*(-\mathbf{k}) \end{pmatrix}$$

$$\hat{\mathcal{H}}_{\text{TI}} = \epsilon \hat{\sigma}_z - iv_F \left[e^{-i\phi \hat{s}_z} \left(\partial_r \hat{s}_x + \frac{1}{r} \partial_\phi \hat{s}_y \right) + \partial_z \hat{s}_z \right] \hat{\sigma}_x - \mu$$

σ : orbital
s: spin

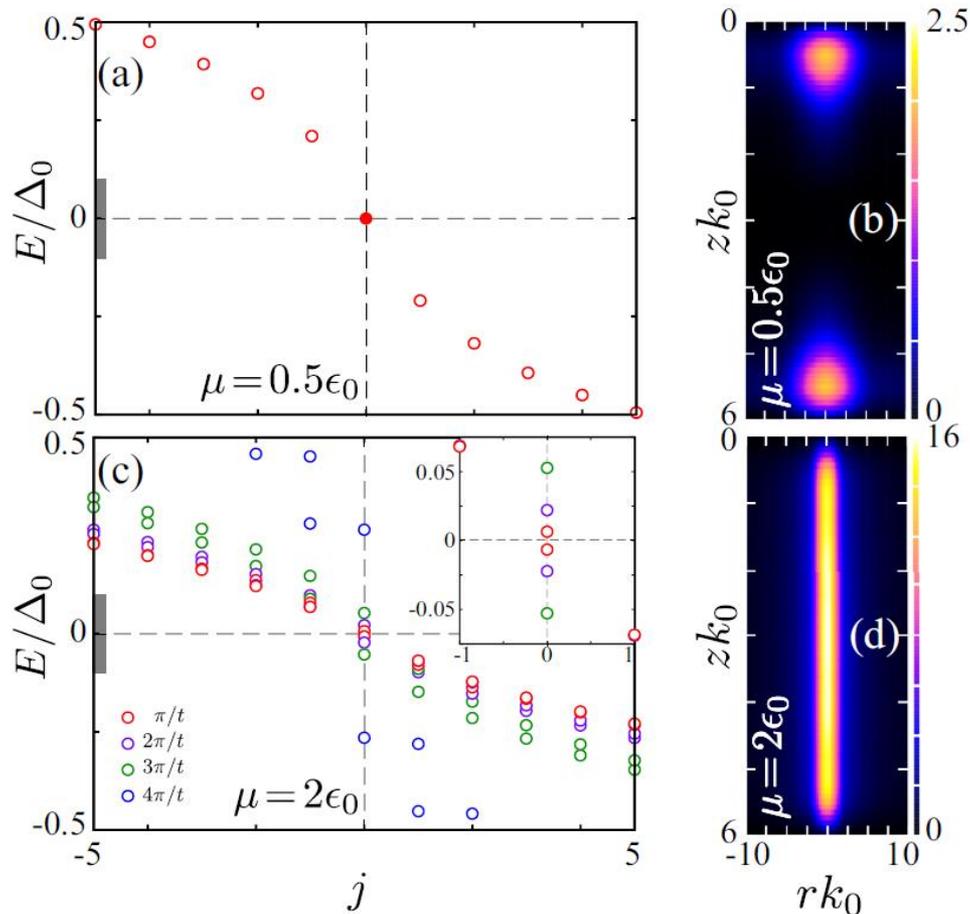
$$\epsilon = -\epsilon_0 - \frac{1}{2m} \left(\partial_r^2 + \frac{1}{r} \partial_r + \frac{1}{r^2} \partial_\phi^2 + \partial_z^2 \right) \quad \hat{\Delta} = e^{i\phi} \Delta_0 i \hat{s}_y$$

□ Parameters for calculation: * Band gap of Bi_2Te_3 : $2\epsilon_0 \sim 0.2\text{eV}$

* SC gap: $\Delta = 0.02\epsilon_0 \sim 2\text{meV}$ * chemical potential: $\mu = 0.5\epsilon_0 \sim 2\epsilon_0$

Majorana bound states in thick TI film

- Energy dispersion and distribution of DOS of quasiparticles



* typical length

$$k_0^{-1} = v_F/\epsilon_0 \simeq 2\text{nm}$$

* coherence length

$$\xi = v_F/\Delta_0 = 50k_0^{-1}$$

MBS at small Fermi level \Leftrightarrow thick TI film

suppression of MBS \leftarrow
bulk conduction bands
induce interactions in
thin TI film

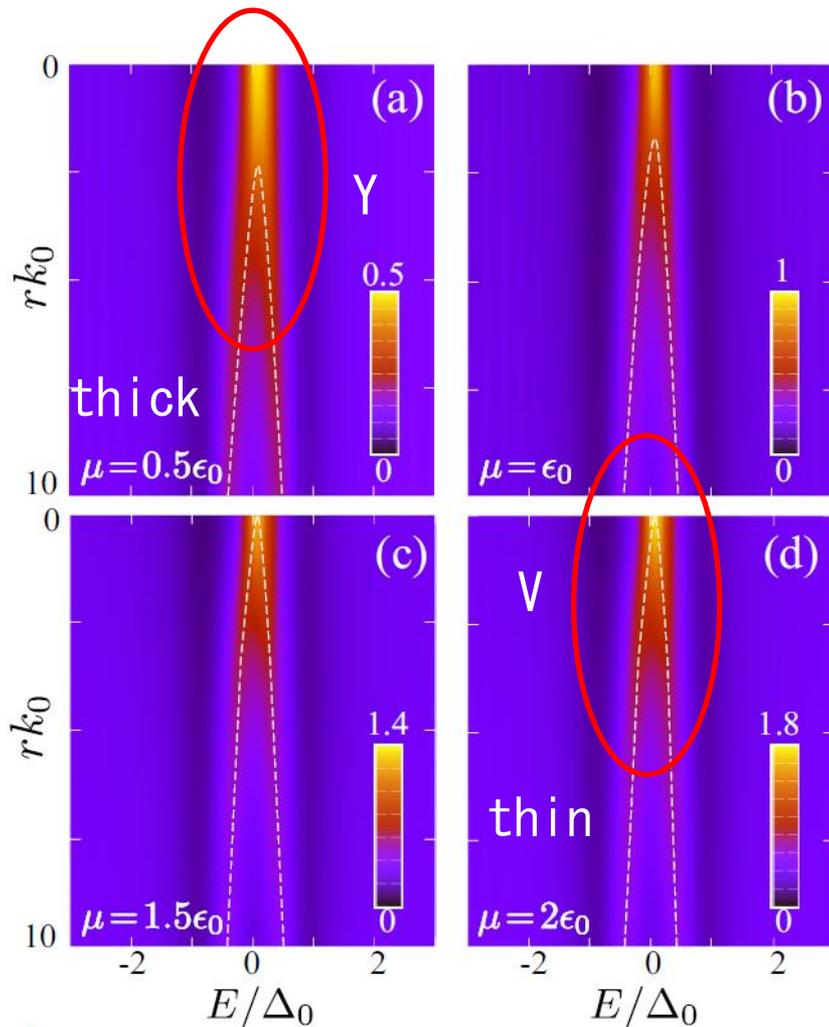
Fu and Kane, PRL **100**, 96407 (2008). Hosur, *et al.*, PRL **107**, 097001 (2011).

Z.-Z. Li, F.-C. Zhang, and Q.-H. Wang, Sci. Rep. **4**, 6363 (2014).

Evolution of DOS with thickness

PRL 115, 177001, 2015

□ Energy-space distribution of DOS of quasiparticles: dI/dV in experiments



smearing factor
in energy

$$\eta = 0.2\Delta_0 \sim 4\text{K}$$

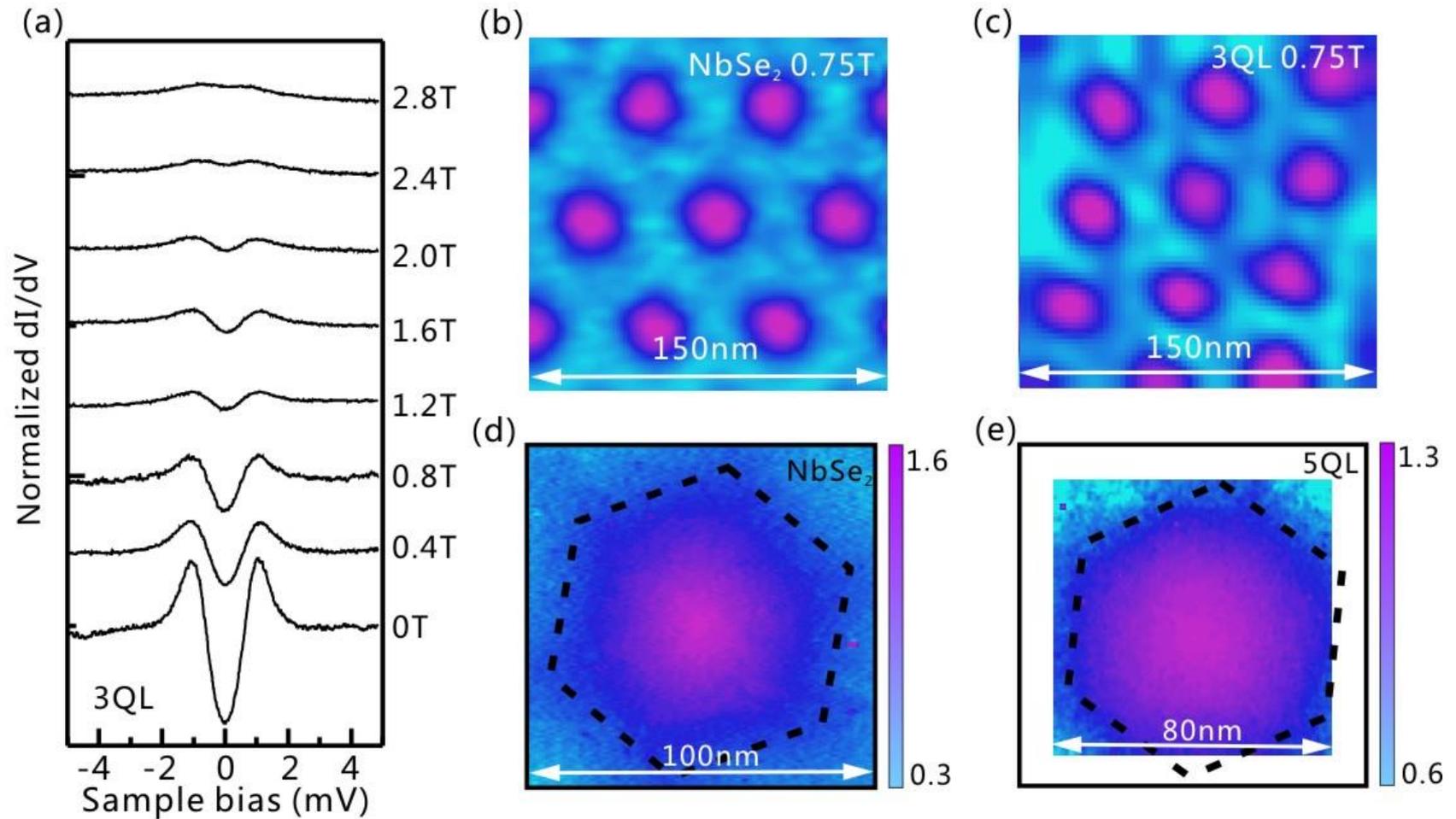
Y shape \Leftrightarrow w MBS
V shape \Leftrightarrow w/o MBS

full agreement with experiments !

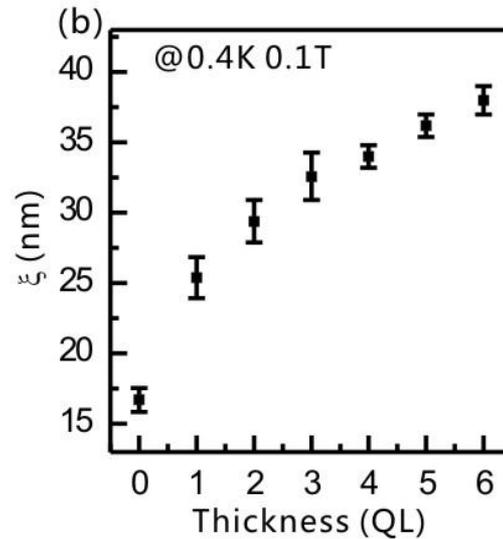
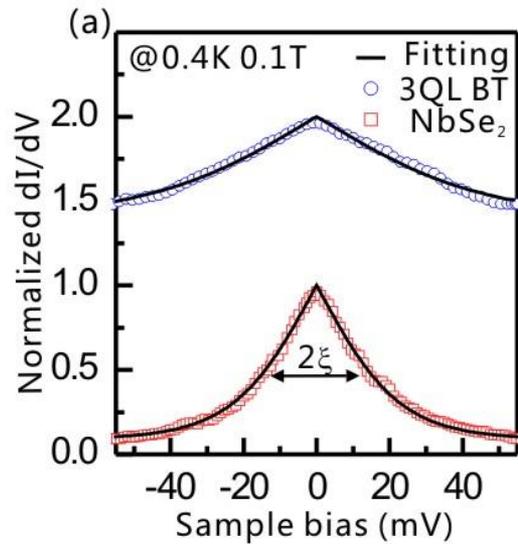
□ Thickness vs. chemical potential

theoretically thickness only cannot induce phase transition, but μ can.

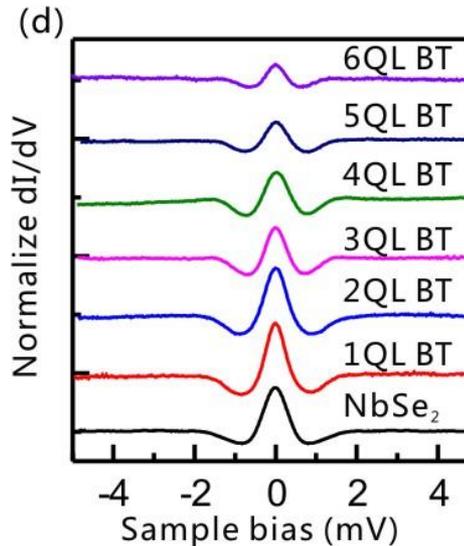
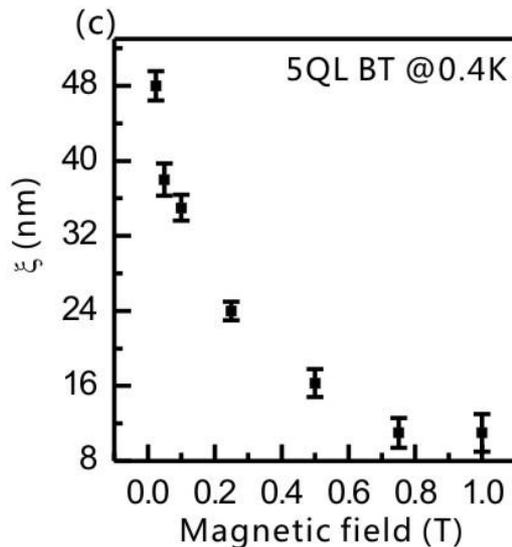
Vortex on topological superconductor



Coherence length and core states

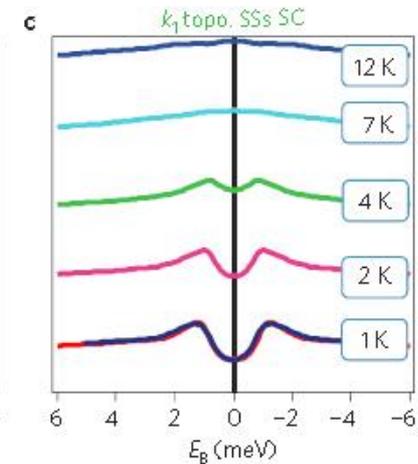
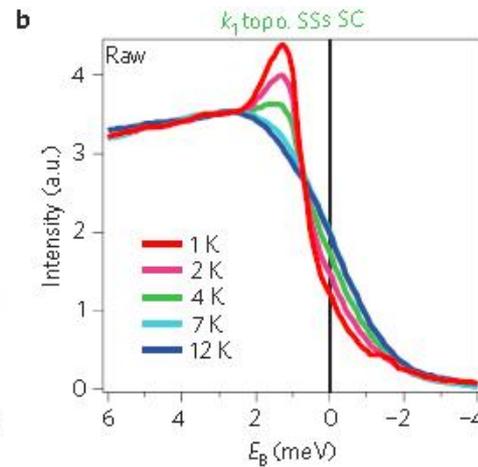
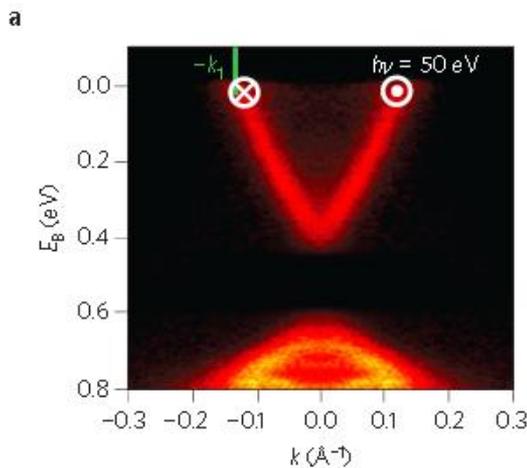


- Coherence length deduced from Vortex
- Much larger than that in $NbSe_2$
- Saturate at 3QL
- Change with magnetic field
- Saturate at $\sim 0.7T$
- Core states observed



Momentum-space imaging of Cooper pairing

3QL



7QL

