## Efficient Representation of Quantum Many-body States with Deep Neural Networks





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# Outline

Neural Network Representation for Many\_body State

Deep vs. Shallow Architecture Limitation of Restricted Bolzmann Machine Power of Deep Bolzmann Machine

How to train Deep Boltzmann Machine

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- **1 Particle:**  $c_0|0\rangle + c_1|1\rangle \in \mathcal{H}$  Dimension is 2
- 2 Particles:  $c_{00}|00\rangle + c_{01}|01\rangle + c_{10}|10\rangle + c_{11}|11\rangle$   $\in \mathcal{H} \otimes \mathcal{H}$  Dimension is 4  $|00\rangle + |11\rangle$  Entanglement

n **Particles**:

 $\begin{array}{l} c_{0\cdots0}|0\cdots0\rangle+\cdots+c_{1\cdots1}|1\cdots1\rangle\\ \in \mathcal{H}^{\otimes n} \quad \text{Dimension is} \quad 2^{\boldsymbol{\mathcal{N}}} \end{array}$ 

Hilbert space is too large !! How to describe it??



#### **The Physical Relevant Corner**

## Many-Body Problem



## low description complexity?

## **Many-Body** Prohlem Hamiltonian $H = \sum_{i}^{n} H_{i} \qquad \begin{array}{c} H_{i} = I_{1} \otimes \cdots \otimes H_{i} \otimes \cdots \otimes I_{n} \\ \swarrow \\ \text{dimension is bounded by c} \end{array}$

dimension is bounded by constant

**low temperature property**  $\begin{array}{c} \text{ground state} \\ |\psi\rangle \end{array} \quad \begin{array}{c} \min \\ |\psi\rangle \end{array} \quad \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \end{array}$ thermal equilibrium property thermal state  $\frac{e^{-\beta H}}{4\pi e^{-\beta H}}$ time evolution  $e^{-iHt}|\psi_0\rangle$  — some simple state

succinct description but not quite useful hard to extract information





# finding ground state is particularly interesting $\min_{|\psi\rangle} \langle \psi | H | \psi \rangle - \lambda \langle \psi | \psi \rangle \text{ with } \langle \psi | \psi \rangle = 1$ QMA-hard for gapless system

intrinsic difficulty

## NP-hard for gapped system if Quantum PCP conjecture is true

## **QMA-hard**

gap between theory and practice:
heuristic algorithm (intuition
& extra information for special instance)
worst-case vs. typical case



## Many-Body Problem

**Mean-field** assumption  $\rho_k \sim_{n \to \infty} |\psi_1\rangle \langle \psi_1 | \otimes \cdots \otimes |\psi_k\rangle \langle \psi_k |$ success for bosonic system (Quantum de Finetti) fail for other strong correlated system Quantum Monte Carlo  $\frac{\langle \psi(\alpha) | H | \psi(\alpha) \rangle}{\langle \psi(\alpha) | \psi(\alpha) \rangle} \sim \sum_{c} p_{c}(\alpha) f_{c}(\alpha) \text{ sign problem}$ **Tensor Network** 



## Many-Body Problem

**Very successful for 1D system:** MPS faithfully represent at least gapped system polynomial time to extract information heuristic algorithm: DMRG, TEBD, etc polynomial time algorithm to find ground state

Landau, Zeph, Umesh Vazirani, and Thomas Vidick. "A polynomial time algorithm for the ground state of one-dimensional gapped local Hamiltonians." *Nature Physics* 11.7 (2015): 566-569.

Arad, I., Landau, Z., Vazirani, U., & Vidick, T. (2016). Rigorous RG algorithms and area laws for low energy eigenstates in 1D. *arXiv preprint arXiv:1602.08828*.

### 2 Some Previous Approach to tackle many-body problem

## Many-Body Problem

Schuch, N., Wolf, M. M., Verstraete, F., & Cirac, J. I. (2007). Computational complexity of projected entangled pair states. *Physical review letters*, 98(14), 140506.

Anshu, A., Arad, I., & Jain, A. (2016). How local is the information in tensor networks of matrix product states or projected entangled pairs states. *Physical Review B*, *94*(19), 195143.

Schwarz, M., Buerschaper, O., & Eisert, J. (2016). Approximating local observables on projected entangled pair states. *arXiv preprint arXiv:1606.06301*.

## fail for 2D system:

- 1. unknown whether PEPS is enough
- 2. extract information is hard
  - i. #P-hard in general case
  - ii. best known approximation algorithm:superpolynomial time under assumptions



#### **Neural Network Zoo**

## Neural Network



http://www.asimovinstitute.org/neural-network-zoo/

# natural to use generative model to represent quantum state



#### Restricted Boltzmann Machine (RBM)



## represent quantum state by neural network itself h



$$\Psi(\mathbf{v}) = \sum_{\mathbf{h}} e^{\mathcal{W}(\mathbf{v},\mathbf{h})}$$

# $\mathcal{W}(\mathbf{v}, \mathbf{h}) = \mathbf{v}^T \mathbf{W} \mathbf{h} + \mathbf{b}^T \mathbf{v} + \mathbf{c}^T \mathbf{h}$ Weight Function

## no intra-layer interactions



 $\mathbf{V}$ 

### 2 Restricted Boltzmann Machine (RBM)



Carleo, Giuseppe, and Matthias Troyer. "Solving the quantum many-body problem with artificial neural networks." *Science 2017.* 

numerical methods (combined with Monte Carlo) transverse field Ising:  $-h\sum_{i}\sigma_{i}^{x}-\sum_{i}\sigma_{i}^{z}\sigma_{j}^{z}$ anti-ferromagnetic Heisenberg:  $\sum_{i,j}\sigma_{i}^{x}\sigma_{j}^{x}+\sigma_{i}^{y}\sigma_{j}^{y}+\sigma_{i}^{z}\sigma_{j}^{z}$ , better than MPS  $\langle i,j \rangle$ 

Deng, Dong-Ling, Xiaopeng Li, and S. Das Sarma. "Exact machine learning topological states." *arXiv preprint arXiv:1609.09060* (2016).

exact representation for topological states SPT: 1D cluster state topological order: toric code

## 2 Restricted Boltzmann Machine (RBM)



Chen, J., Cheng, S., Xie, H., Wang, L., & Xiang, T. (2017). On the Equivalence of Restricted Boltzmann Machines and Tensor Network States. *arXiv preprint arXiv:1701.04831*. restricted Boltzmann Machine —> Tensor Network Tensor Network —> RBM with given architecture

Deng, Dong-Ling, Xiaopeng Li, and S. Das Sarma. "Quantum Entanglement in Neural Network States." *PhysRevX.7.021021*. area law (local connection) random RBM: entanglement spectrum, not thematize

Huang, Yichen, and Joel E. Moore. "Neural network representation of tensor network and chiral states." *arXiv preprint arXiv:1701.06246* (2017). using Grassmann number quasi-local  $H = \sum_{\vec{x} \in \mathbf{Z}^2} \left( c^{\dagger}_{\vec{x}+\vec{i}}c_{\vec{x}} + c^{\dagger}_{\vec{x}+\vec{j}}c_{\vec{x}} + ic^{\dagger}_{\vec{x}+\vec{j}}c^{\dagger}_{\vec{x}} + h.c. \right) - 2\mu \sum_{\vec{x} \in \mathbf{Z}^2} c^{\dagger}_{\vec{x}}c_{\vec{x}}$ 



#### Limitation of RBM

## Neural Network



limitation to represent Tensor Network (given architecture)

#### **Copy from**

Chen, J., Cheng, S., Xie, H., Wang, L., & Xiang, T. (2017). On the Equivalence of Restricted Boltzmann Machines and Tensor Network States. *arXiv preprint arXiv:1701.04831*.

## Neural Network

#### **Universal Approximation Theorem**

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In order to approximate probability distribution with **k** boolean variables, the number of hidden neurons is at most **2^k** *without constraints on architecture* (analog to bond dimension in MPS)

Le Roux, Nicolas, and Yoshua Bengio. "Representational power of restricted Boltzmann machines and deep belief networks." *Neural computation* 20.6 (2008): 1631-1649.

reasonable to consider: Efficient Representation: poly(n) parameters without constraints on architecture





under reasonable complexity assumptions RBM can not represent the following state efficiently: state generated by efficient quantum computer PEPS and other tensor network state ground state of local gapped Hamiltonian

## not closed in a quantum phase and dynamics !

## otherwise polynomial hierarchy collapse (a generalization of P=NP)

Gao, Xun, and Lu-Ming Duan. "Efficient Representation of Quantum Many-body States with Deep Neural Networks." *arXiv preprint arXiv:1701.05039* (2017).



hidden neurons are P/poly: can be solved conditionally independent in polynomial size circuit the circuit not easy to construct

P/poly:  $\Psi(\mathbf{v})$ polynomial-size boolean circuit  $\Psi(\cdot)$ input of function  $\mathbf{v}$ 

5

no need to know how to construct RBM (circuit) !

### Proof of Limitation for RBM



cluster state: ignore the colors



i. state generated from simple dynamics on simple initial state

ii. ground state of gapped, local, commuting Hamiltonian

iii. PEPS

time evolution on  $|+\rangle^{\otimes N}$ 

5

$$\mathcal{H} = -\sum_{\langle i,j \rangle} JZ_i Z_j + \sum_i B_i Z_i \qquad J = \frac{\pi}{4}$$

#### Proof of Limitation for RBM

## Neural Network

## **GWD** state:

5



#### cluster state with single qubit unitary and L blue: z-axis with angles, then Hadamard Supre green: Hadamard red

Gao, Xun, Sheng-Tao Wang, and Lu-Ming Duan. "Quantum Supremacy for Simulating A Translation-Invariant Ising Spin Model." *Phys.Rev.Lett*. 2017



5



#### translational invariant ! only depends on size

## cluster state is resource state in measurement-based quantum computing



5



#### translational invariant ! only depends on size

if  $\Psi(v)$  of GWD state can be represented by RBM efficiently

 $P^{\#P} \subseteq P^{P/poly}$  widely believed not true: polynomial hierarchy collapse ! generalization P=NP, very unlikely

extend to approximate case in terms of trace distance

with another reasonable complexity assumptions

related to quantum supremacy for sampling random quantum circuit of supported by quantum chaos theory

**Google's quantum supremacy plan** Boixo, Sergio, et al. "Characterizing quantum supremacy in near-term devices." *arXiv preprint arXiv:* 1608.00263 (2016).







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**Characterize RBM state** 



## Although limitation of RBM still useful in practice

how to characterize the state represented by RBM? string-bond state?





Deep Boltzmann Machine (DBM)

## Deep Neural Networks

## Deep vs. Shallow

DBM can represent the following state efficiently:

**quantum computing or dynamics:** O(nT) **ground state:**  $O\left(\frac{1}{\Delta}\left(n + \log \frac{1}{\epsilon}\right)m^2\right)^{-1}$  depth or evolution time energy gap in Hamiltonian

## gapless, even non-local

**tensor network state:**  $O(D^{2d}n)$  D: bond dimension d: coordination number

## closed in a quantum phase and dynamics



#### **Deep Neural Proof of Power of** 2 Networks **Deep Boltzmann Machine Quantum Computing or Dynamics** rule I simulates matrix multiplication Hadamard Gate Phase phase gadget is not necessary $Z(\theta)$ Gate absorbed into bias term $\phi_{t+1}(\cdots x_i, x_{i+1}\cdots) = (-1)^{x_i x_{i+1}} \phi_t(\cdots x_i, x_{i+1}, \cdots)$ control-Z rule II simply multiplication (d) Construction of $\phi_{t+1}$ $H, Z(\theta), \text{control-}Z \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$ (d) Construction of $\phi_{t+1}$ **Identity Gadget** H $Z(\theta)$ $Z(\theta)$ universal set **DBM** representation HΙ for $Z(\theta)$ $Z(\theta)$ Ι quantum computing $\phi_0$

Øт

## Proof of Power of Deep Boltzmann Machine Quantum Computing or Dynamics

### dynamics: quantum simulation

2

Lloyd, S. Universal quantum simulators. Science 273, 1073 (1996).
Poulin, D., Qarry, A., Somma, R. & Verstraete, F. Quantum simulation of time-dependent hamiltonians and the convenient illusion of hilbert space. Physical review letters 106, 170501 (2011).
Berry, D. W., Childs, A. M., Cleve, R., Kothari, R. & Somma, R. D. Simulating hamiltonian dynamics with a truncated taylor series. Physical review letters 114, 090502 (2015).

# closed under a quantum phase: simulating adiabatic evolution



#### rule I simulates contraction of bond index



### Proof of Power of Deep Boltzmann Machine Tensor Network

## Deep Neural Networks

#### see also

Chen, J., Cheng, S., Xie, H., Wang, L., & Xiang, T. (2017). On the Equivalence of Restricted Boltzmann Machines and Tensor Network States. *arXiv preprint arXiv:* 1701.04831.

Huang, Yichen, and Joel E. Moore. "Neural network representation of tensor network and chiral states." *arXiv preprint arXiv:1701.06246* (2017).

why not use universal approximate theorem

representing  $A_{x_1 \cdots x_c}$  directly ?

## physical property is very sensitive to local tensor !

#### toric code



$$\begin{split} T_{ijkl} &= 1 \rightarrow T_{ijkl} = 1, \text{ if } i+j+k+l = 0 \mod 2 \\ T_{ijkl} &= 0 \rightarrow T_{ijkl} = \epsilon, \text{ if } i+j+k+l = 1 \mod 2 \end{split}$$

so long as  $\ \epsilon \neq 0$ topological entanglement entropy disappears ! Chen, X., Zeng, B., Gu, Z. C., Chuang, I L., & Wen, X. G. (2010). Tensor product representation of a topological ordered phase: Necessary symmetry conditions. *Physical Review B*, *82*(16), 165119.

#### Proof of Power of Deep Boltzmann Machine Ground State

## Deep Neural Networks

#### inspired by

2

Berry, D. W., Childs, A. M., Cleve, R., Kothari, R., & Somma, R. D. (2015). Simulating Hamiltonian dynamics with a truncated Taylor series. *Physical review letters*, *114*(9), 090502. using Taylor series instead of Trotter decomposition:

exponential improvement on precision







3

## Deep Neural Networks

## non-local even non-sparse property (compared to PEPS):

proved result for representing ground state beyond 1D

#### appropriate for highly entanglement state like time-evolution

expected to have less parameters

#### deeper compared to RBM

harder to extract information or training ?

both #P-complete (computing local observable)



**Training DBM** 

## Deep Neural Networks

## Extract Information prototype Monte-Carlo algorithm

$$\rangle = \frac{\sum_{h} p_{h} f_{h}}{\sum_{h} p_{h} g_{h}}$$

 $\frac{p_{h'}}{p_h}$  easy to compute Metroplis algorithm  $Pr(h \to h') = \min\left(1, \frac{p_{h'}}{p_h}\right)$ 

 $\langle \psi | O | \psi$ 

f and g are easy to compute

fluctuation is too large?the same problems alsolocal minimum?occur in RBM ?

## Deep Neural Networks

## seems hard to train (sign problem)

inevitable! (intrinsic) trade off between representational power & computational difficulty

## even though our work shows we can benefit a lot from depth like in deep learning

quasi DBM/RBM fewer neurons in the second hidden layer

or other deep architecture



#### **Restricted Boltzmann Machine:**

limitation for dynamics (quantum phase), PEPS, GS tool of proof: complexity theory conjecture, string-bond state

#### **Deep Boltzmann Machine:**

most of physical states (physical relevant corner?) dynamics, tensor network ground state (even gapless & non-local)

#### **Training Deep Boltzmann Machine:**

Prototype Monte Carlo connection with other fields